

The 18th CNS International Summer School (CNSSS19)

Aug 21-27, 2019, Tokyo



Theoretical Study of Superheavy Nuclei

--- Structure Properties & Synthesis Mechanism ---

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NSFC, CAS & MOST

HPC Cluster of KFTP/ITP-CAS

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Atomic nucleus & nuclide

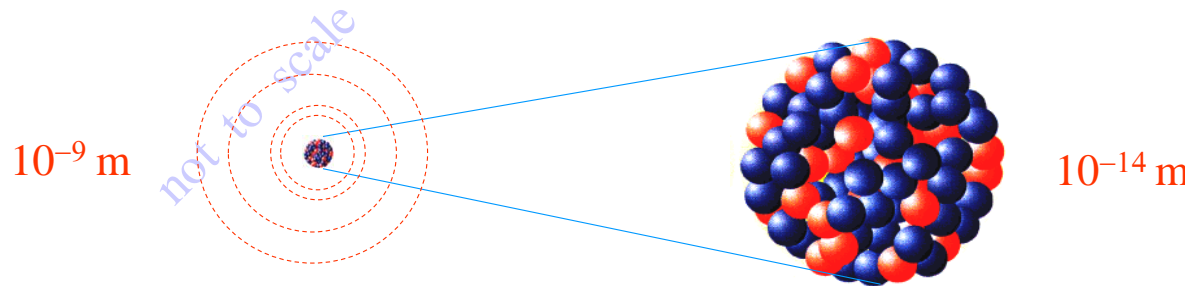
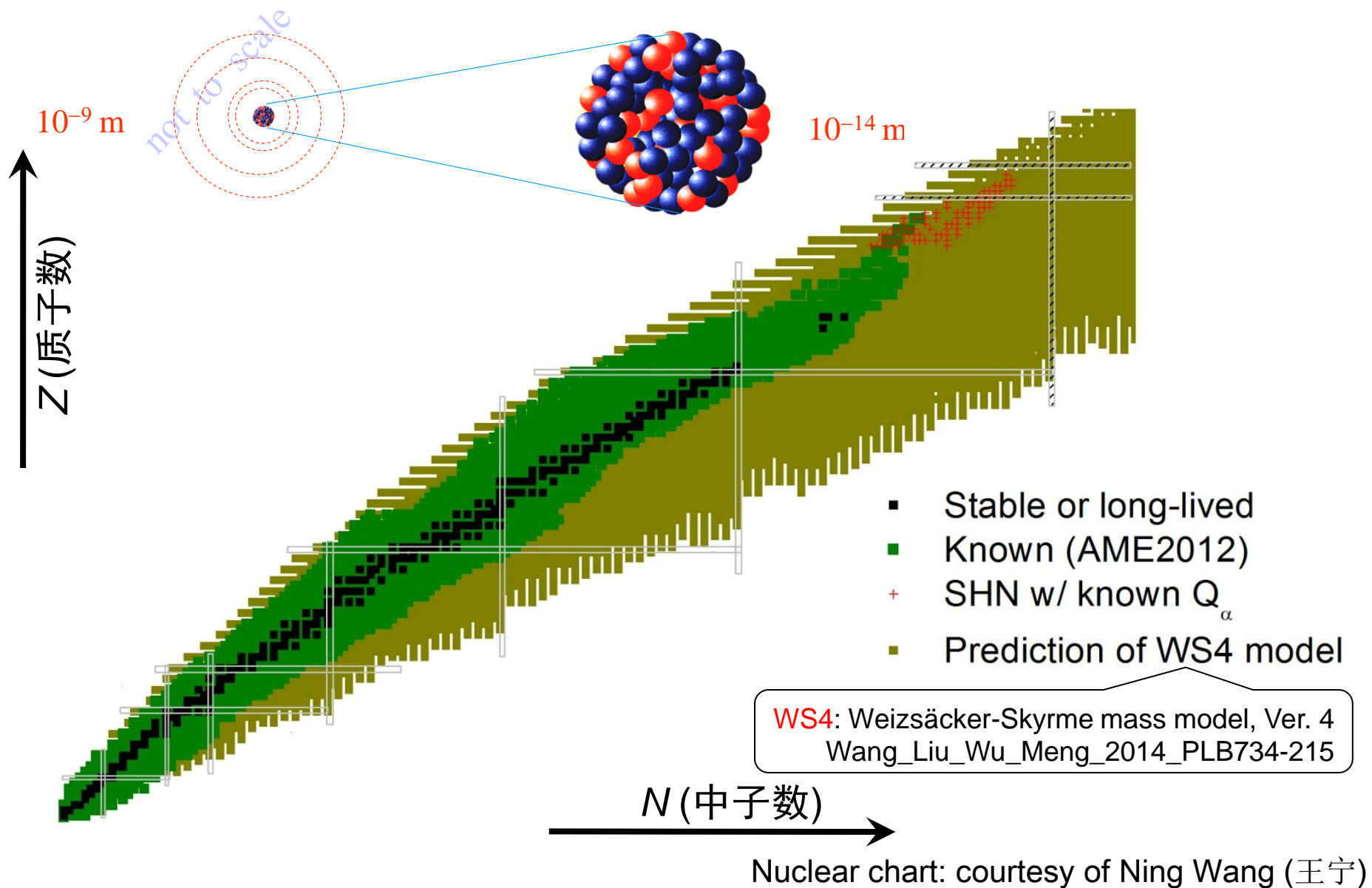
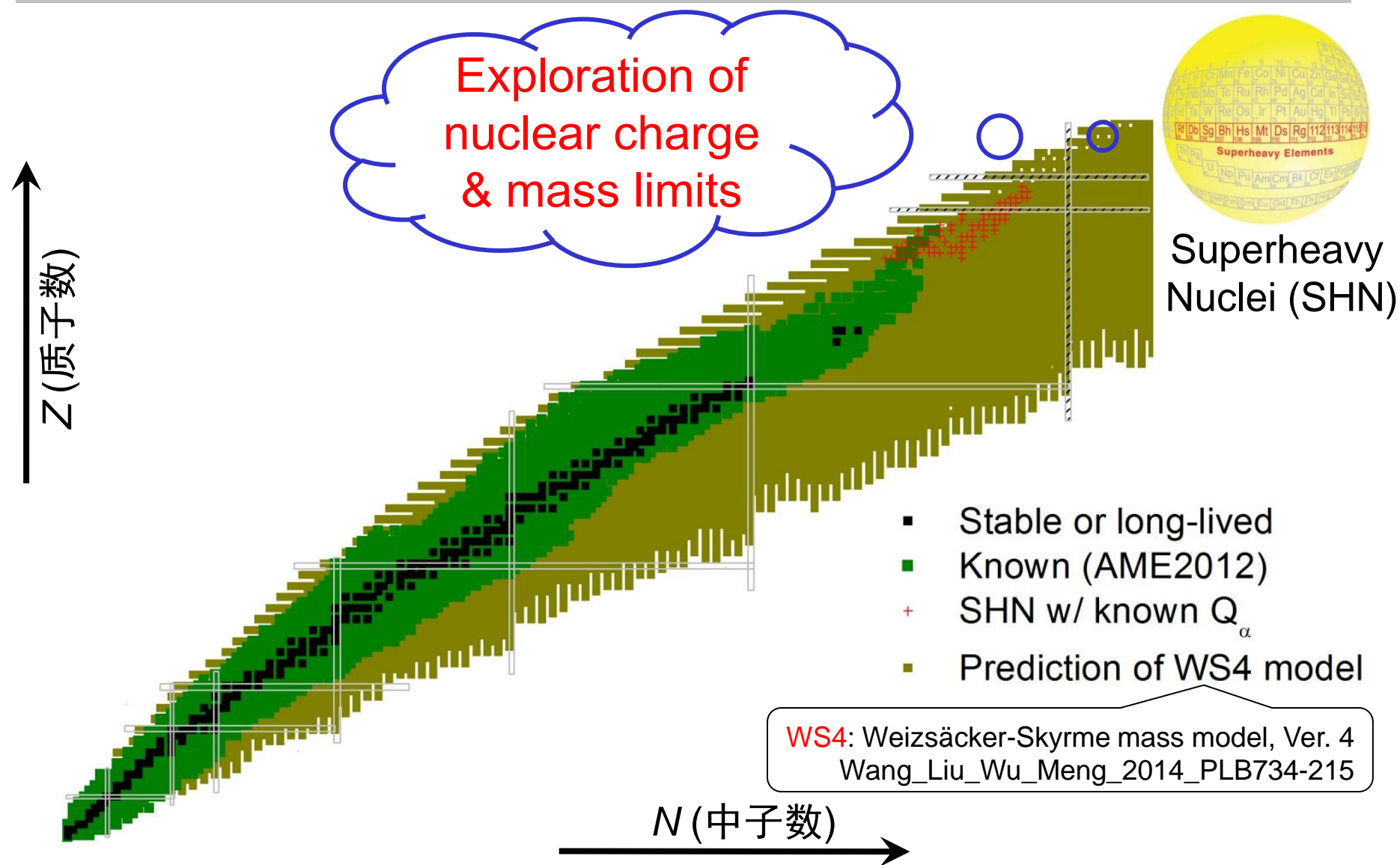


Chart of nuclides



Extension of chart of nuclides: Superheavy nuclei



Nuclear chart: courtesy of Ning Wang (王宁)

Main content of lectures

- ❑ Predictions, experimental progress & challenges
- ❑ Structure, decay & fission properties
- ❑ Synthesis mechanism: Heavy-ion fusion & multi-nucleon transfer reactions

Further readings:

- *Special Issue on Superheavy Elements*, Nucl. Phys. A 944 (2015)
- Lu, Zhao & SGZ, Chapter 5 in *Relativistic Density Functional for Nuclear Structure* (World Scientific, 2016, Editor: Jie Meng)
- Ackermann & Theisen, Phys. Scr. 92 (2017) 083002
- Giuliani et al., Rev. Mod. Phys. 91 (2019) 011001

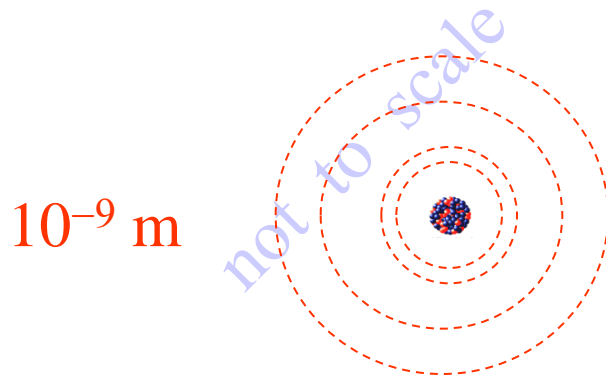
Lecture 1

- Predictions of the island of stability of SHN
- Experimental progress on synthesis of SHN
- Challenges in synthesizing SHN

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Nuclear charge limit on atomic level

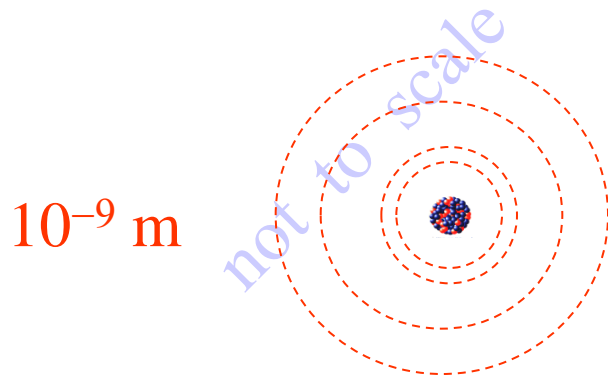


□ Bohr's model: Velocity of innermost electron

- Hydrogen: $v_1 = \alpha c$
- Heavy atoms: $v_1 \sim Z\alpha c$, $Z < 137$

Greiner & Reinhardt 1994, QED
Indelicato 2013, Nature 498, 40-41

Nuclear charge limit on atomic level



□ Bohr's model: Velocity of innermost electron

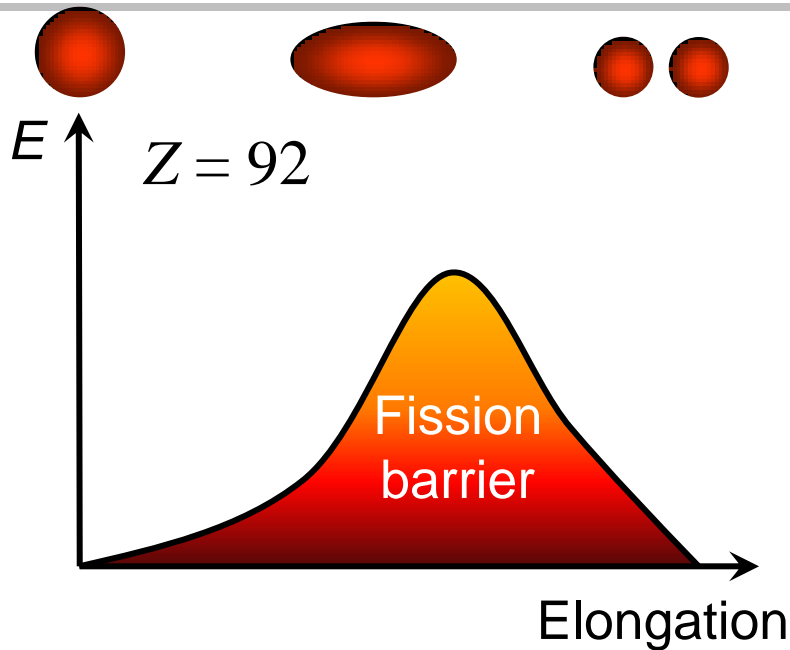
- Hydrogen: $v_1 = \alpha c$
- Heavy atoms: $v_1 \sim Z\alpha c$, $Z < 137$

□ QED: Energy of innermost electron

- Point-like nucleus:
 $E_1 \sim m_e c^2 [(1 - (Z\alpha)^2)^{1/2} - 1]$, $Z < 137$
- Finite-size nucleus: $Z < 173$

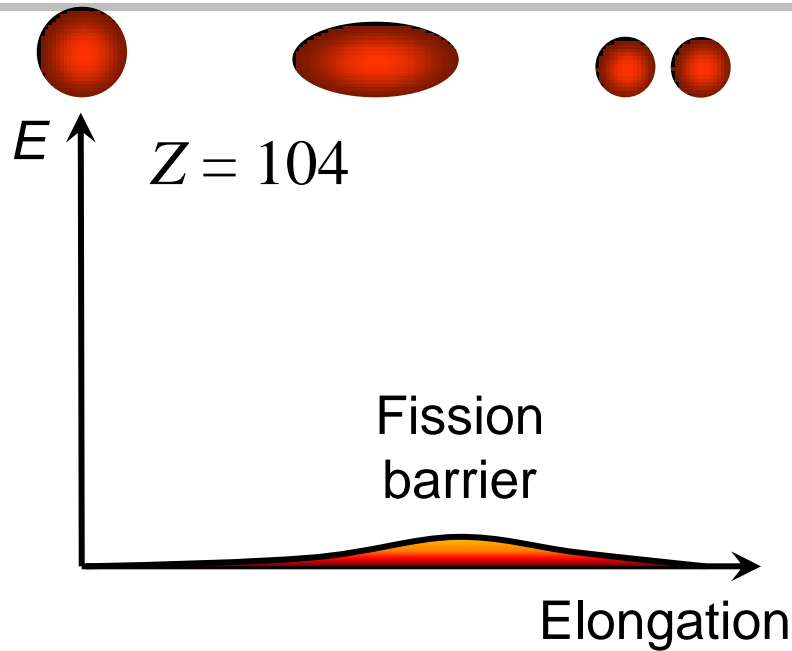
Greiner & Reinhardt 1994, QED
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Nuclear charge & mass limits on nuclear level



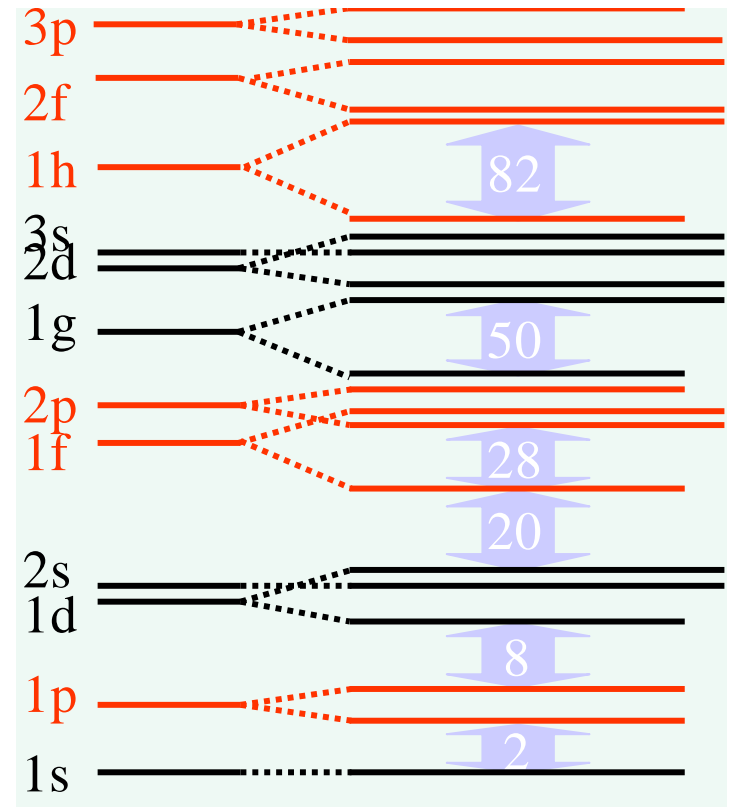
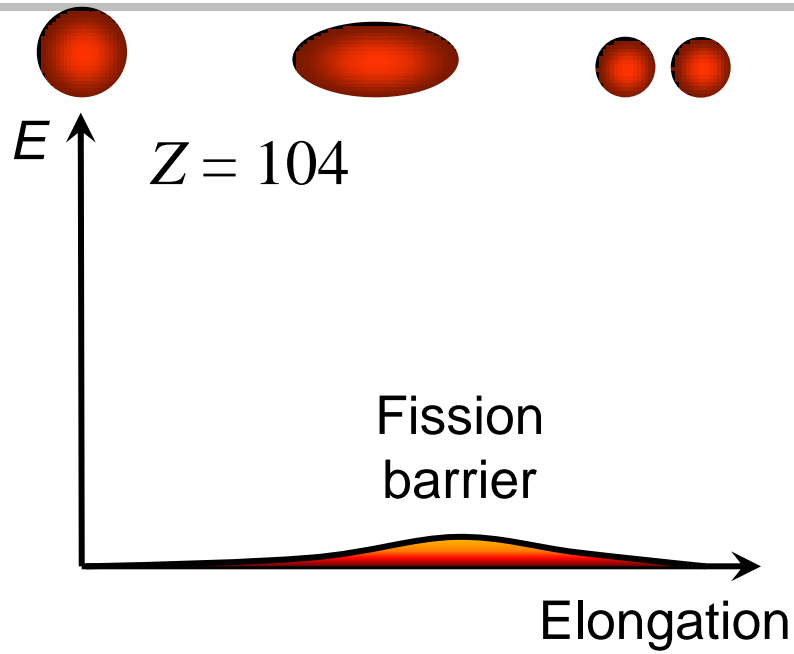
- Nuclear force: Short-ranged, attractive & saturated
- Coulomb force: Long-ranged & repulsive
- As charged liquid drops, nuclear charge & mass limited by competition of nuclear & Coul. forces

Nuclear charge & mass limits on nuclear level



- Nuclear force: Short-ranged, attractive & saturated
- Coulomb force: Long-ranged & repulsive
- As charged liquid drops, nuclear charge & mass limited by competition of nuclear & Coul. forces

Quantum shell effects



Single nucleon potential, spectra & magicities

- Harmonic oscillator & square well
- Woods-Saxon
- Self-consistent (self-bound)



Eugene Paul Wigner

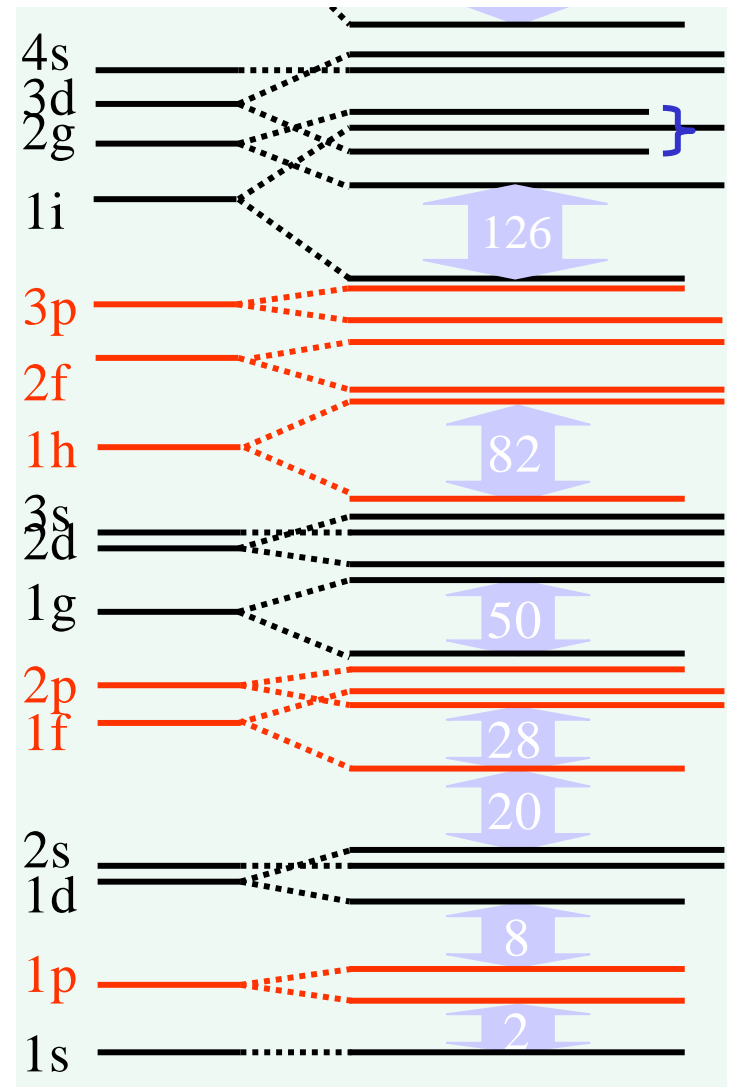


Maria Goeppert Mayer

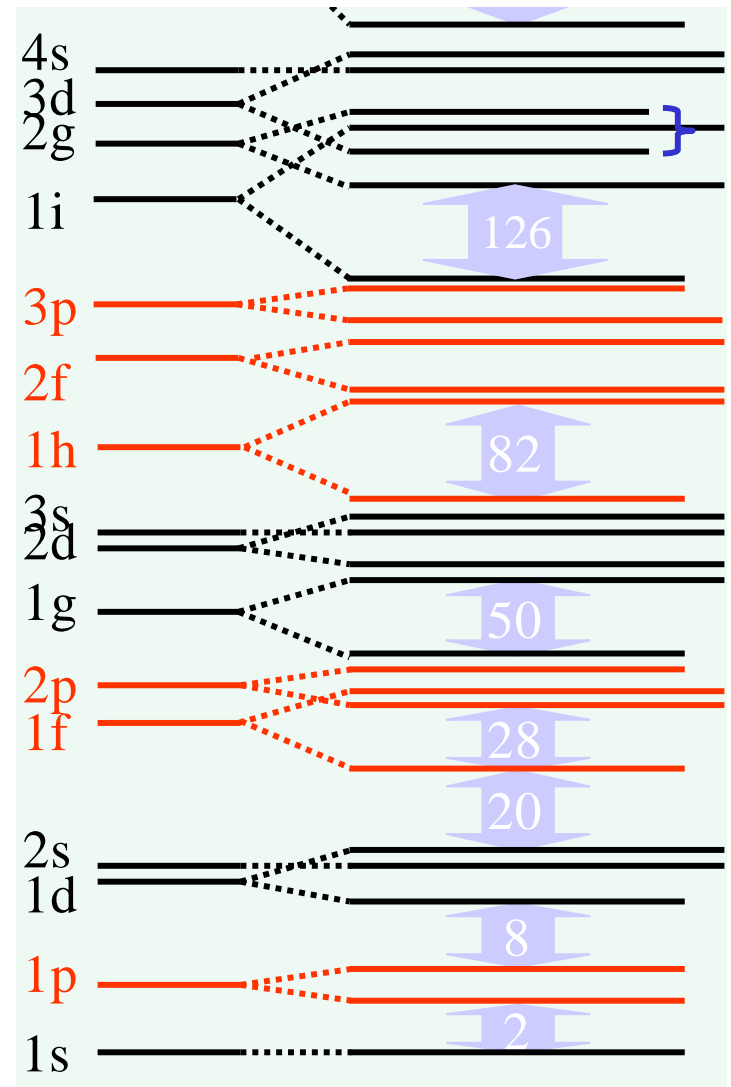
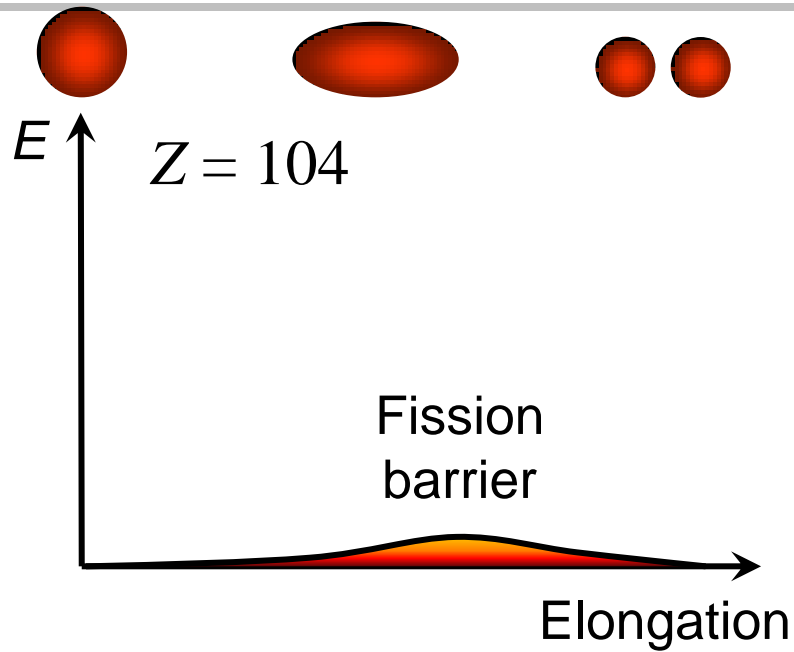


J. Hans D. Jensen

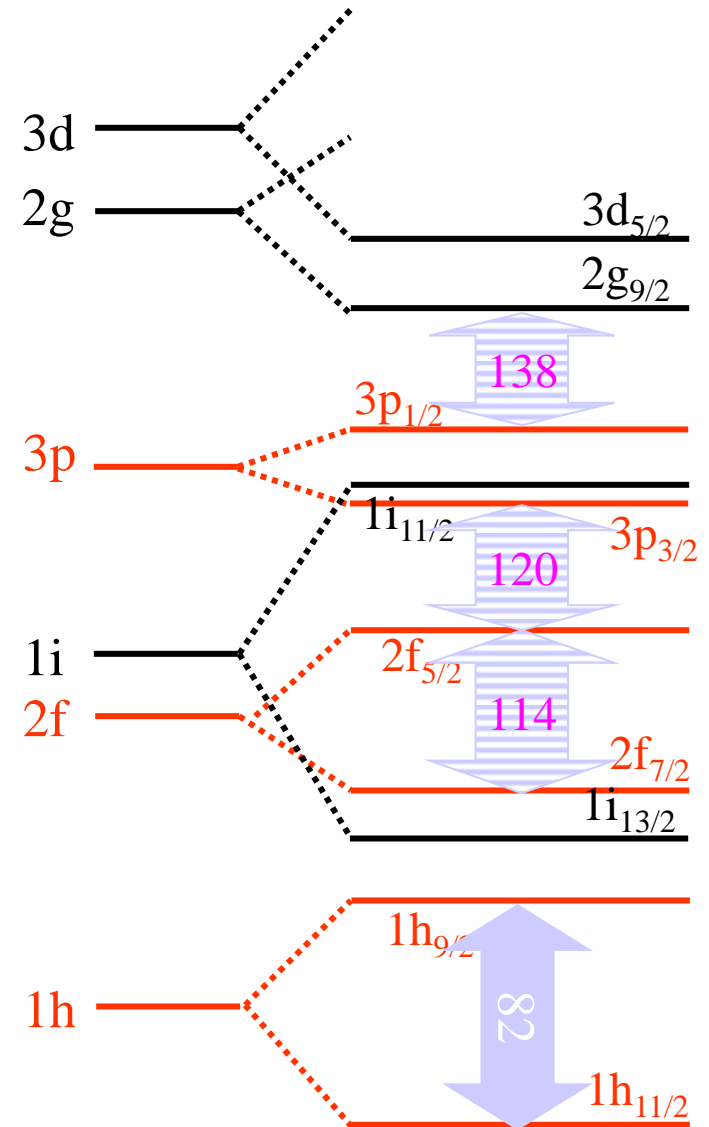
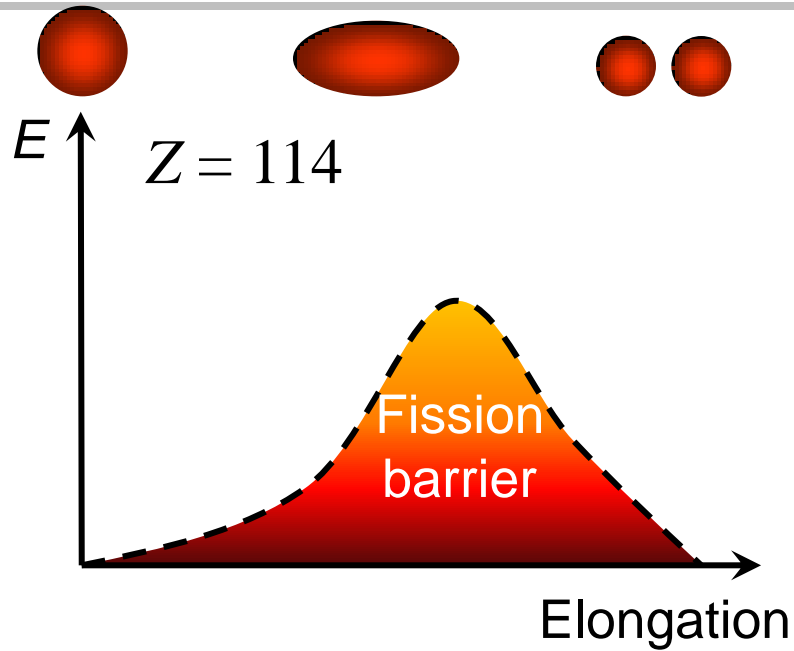
Nobel Prize of Physics (1963)



Quantum shell effects



Quantum shell effects \Rightarrow Superheavy nuclei



Before 1966: Extrapolation of shell structure

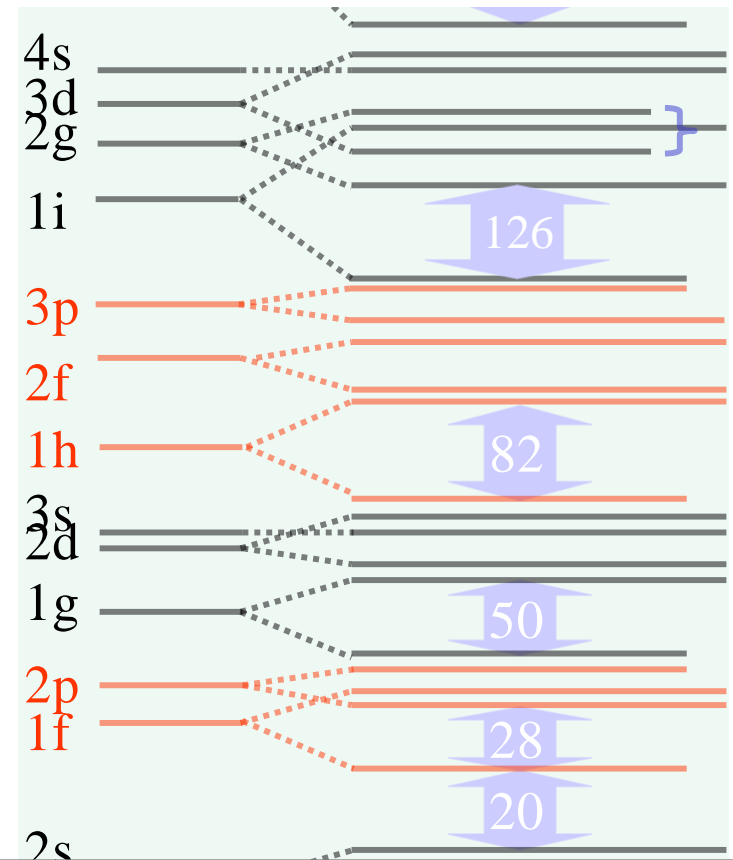
Wheeler 1955, in N. Bohr & Development of Phys.

bility analysis of very heavy nuclei. In so far as one can extrapolate the stability criteria for very heavy nuclei, one is led to expect the existence of nuclei twice as heavy as known nuclear species, that can be created by massive neutron irradiation, and that will live long enough to be studied in the laboratory.

Scharff-Goldhaber 1957_Nucleonics15-122

Relatively long-lived isotopes may well be found among the far-transuranic nuclides because of magic-number stability

There may be, for instance, another region of relative stability at the doubly magic nucleus ${}_{126}X^{310}$ (the closing of the j neutron shell).



Werner_Wheeler 1958_PR109-126

PHYSICAL REVIEW

VOLUME 109, NUMBER 1

JANUARY 1, 1958

Superheavy Nuclei

FREDERICK G. WERNER AND JOHN A. WHEELER

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

(Received August 19, 1957)

1966: Semi-quantitative shell effects

1.E.2:
1.D.4

Nuclear Physics **81** (1966) 1—60; © North-Holland Publishing Co., Amsterdam

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NUCLEAR MASSES AND DEFORMATIONS

WILLIAM D. MYERS and WLADYSLAW J. SWIATECKI

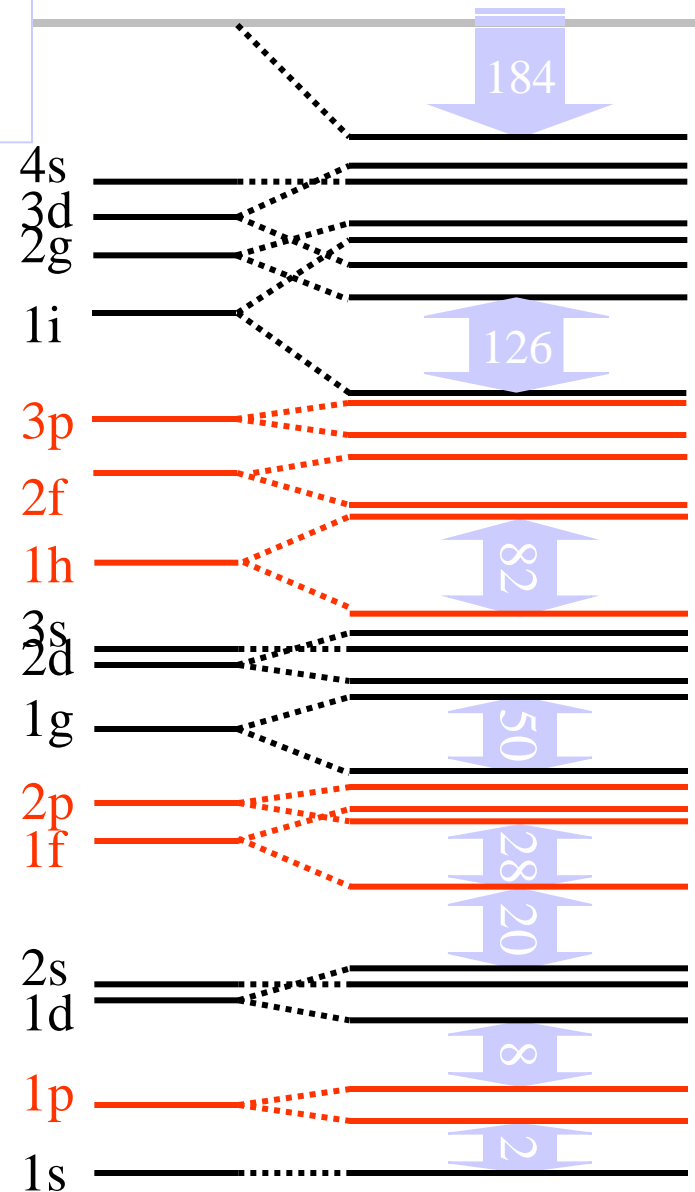
Lawrence Radiation Laboratory, University of California, Berkeley, California[†]

Received 7 September 1965

half-life. This is illustrated in fig. 2 where we have plotted the deformation energy predicted by our mass formula for the case $Z = 126$, $N = 184$. This nucleus has a fissility parameter $x = 1.05$; as a result, in the absence of shell effects, it would have a vanishing barrier against fission and a spontaneous fission half-life of the order of nuclear collective oscillations or 10^{-22} sec. Because of the assumed doubly magic number, however, the ground-state mass of this nucleus would be depressed (by about 10.2 MeV according to our formula). Since this depression is, according to our treatment, a rapidly decreasing function of deformation, there results a considerable barrier against fission, with a height of 9.0 MeV: the extra binding associated with the doubly magic number has stabilized the otherwise highly fissile nucleus. An

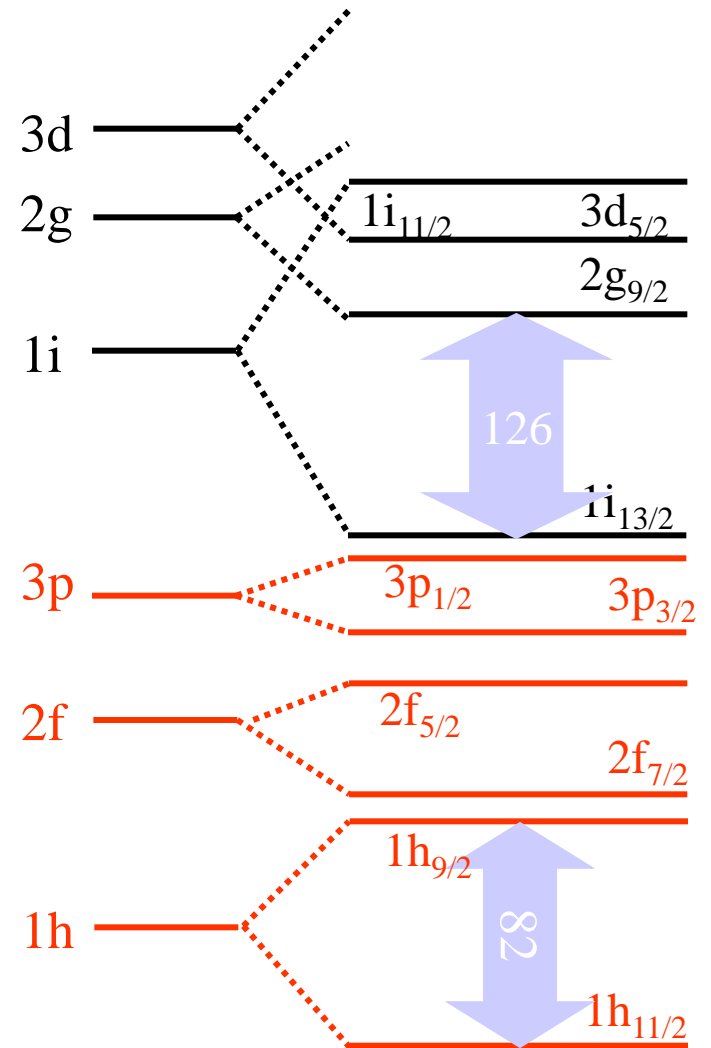
Strong Coulomb effects on proton shell structure

Large A
Large Z



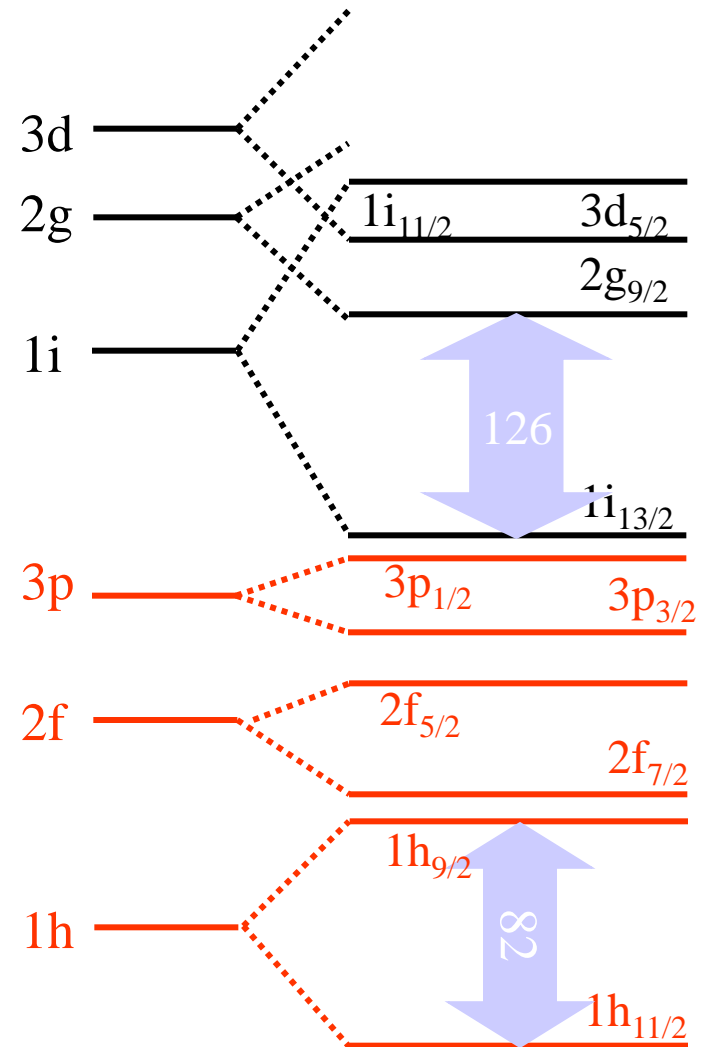
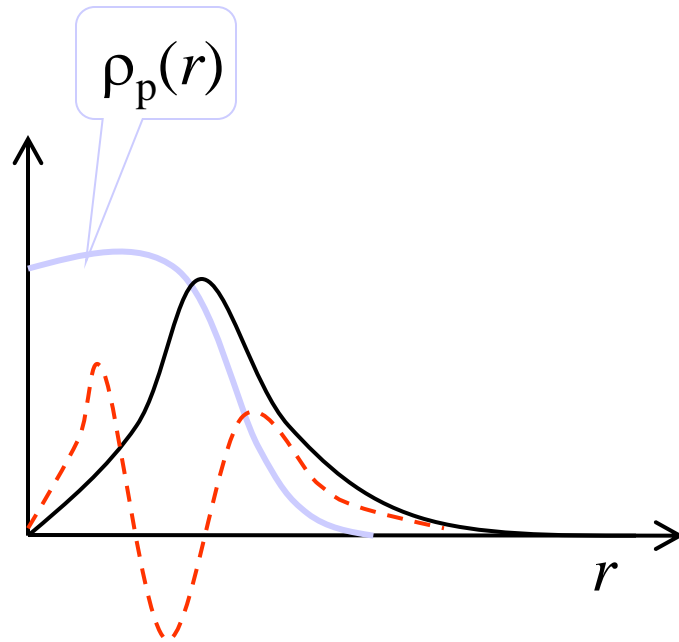
Strong Coulomb effects on proton shell structure

Coulomb Repulsion



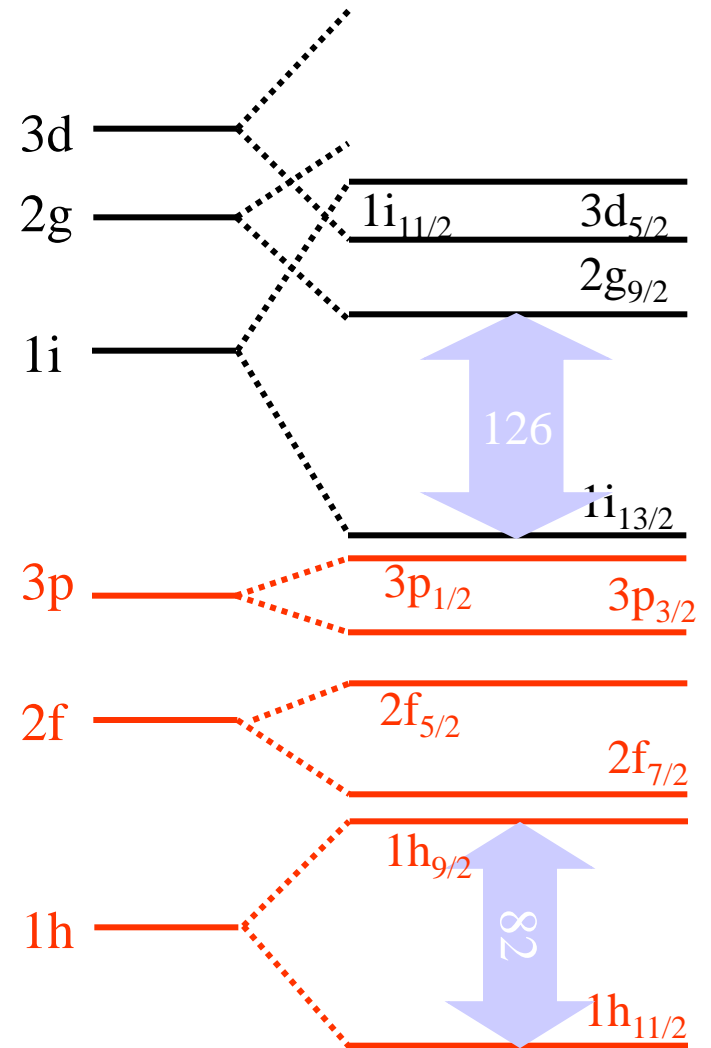
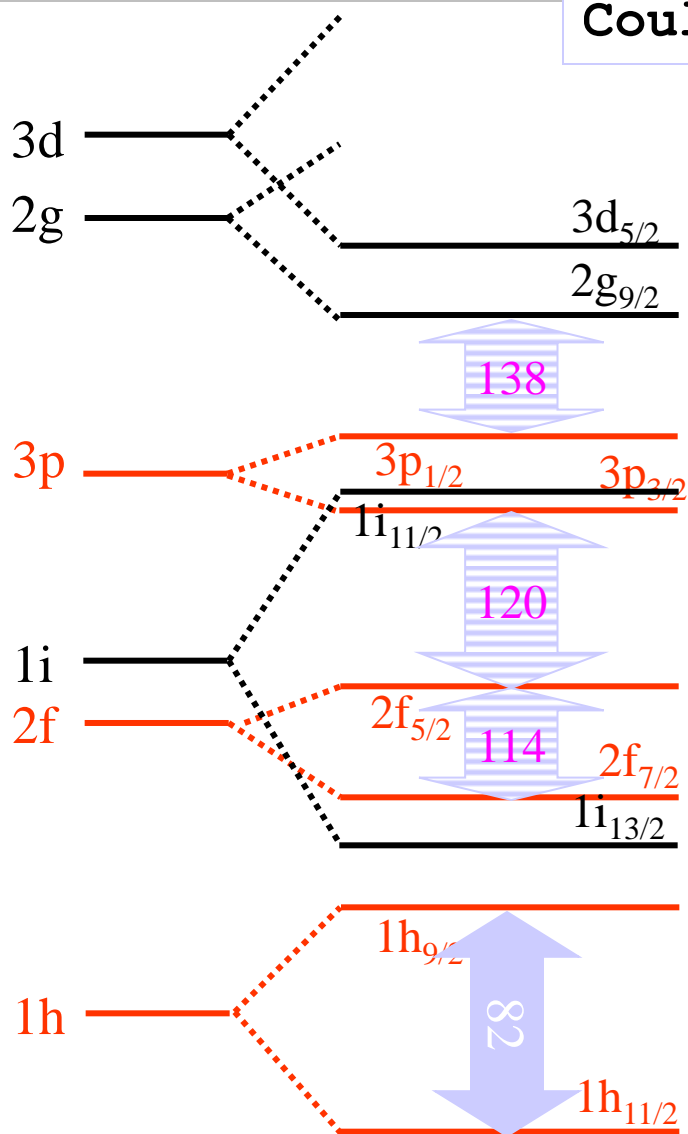
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Strong Coulomb effects on proton shell structure

Coulomb Repulsion



1966 & after: Quantitative shell effects

- Closed shells for $Z > 82$ and $N > 126$ in a diffuse potential well
 - A Woods-Saxon potential Sobiczewski...1966_PhysLett22-500
 - $Z = 114$ & $N = 184$; $B_f = 10$ MeV
- Predictions of new magic regions and masses for super-heavy nuclei from calculations with realistic shell model single particle Hamiltonian
 - A non-local potential Meldner1967_ArkivFysik36-593
 - $Z = 114$ & $N = 184$

1967: A different opinion

VOLUME 18, NUMBER 17

PHYSICAL REVIEW LETTERS

24 APRIL 1967

SHAPE OF HEAVY NUCLEI*

Philip J. Siemens† and H. A. Bethe

Laboratory of Nuclear Studies, Cornell University, Ithaca, New York

(Received 10 March 1967)

The energy reduction for this semimagic nucleus is likely to be considerably less than for Pb^{208} for which it is about 15 MeV. On the other hand, without shell structure, the elongated shape has an energy about 18 MeV less than the sphere for $Z = 114$. We therefore believe that shell effects are unlikely to make the nucleus $Z = 114$ stable.

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- More quantitative investigations needed !!!
 - Single particle shell structure
 - How much stability the closed shell brings

1966 & after: Quantitative shell effects

- Closed shells for $Z > 82$ and $N > 126$ in a diffuse potential well
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 - A non-local potential Meldner1967_ArkivFysik36-593
 - $Z = 114$ & $N = 184$
- “On the spontaneous fission of nuclei with Z near 114 and N near 184”
 - Nilsson potential; potential energy surface with shell corrections calculated by the Strutinsky method Nilsson...1968_NPA115-545
 - $B_f = 10$ MeV & $T_{1/2}(\text{s.f.}) \sim 10^{19}$ years
- “On the stability of superheavy nuclei against fission”
 - Nilsson potential; potential energy surface with rotational invariance
 - $T_{1/2}(\text{s.f.}) \sim 10^{0-17}$ years Mosel_Greiner1969_ZPA226-261

More investigations

Not the end yet ! For the location of the island of stability, more different predictions:

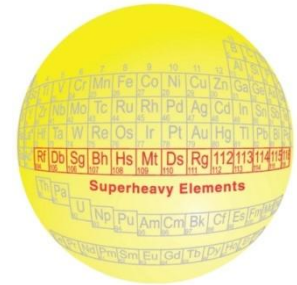
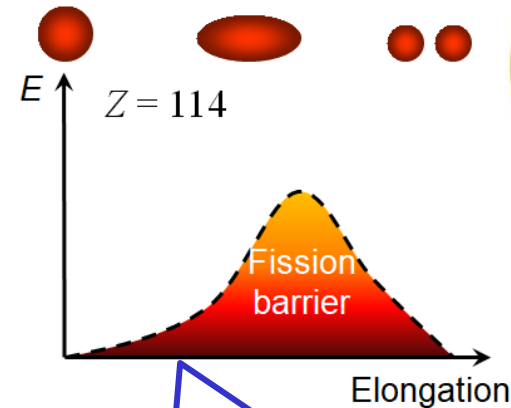
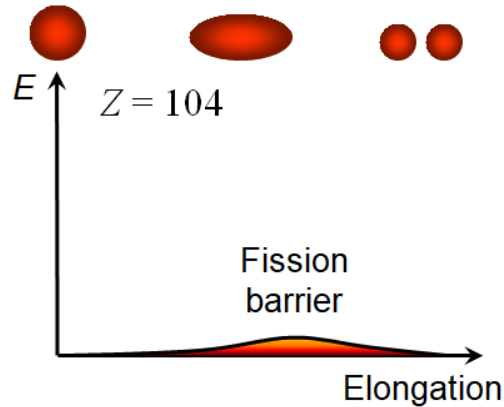
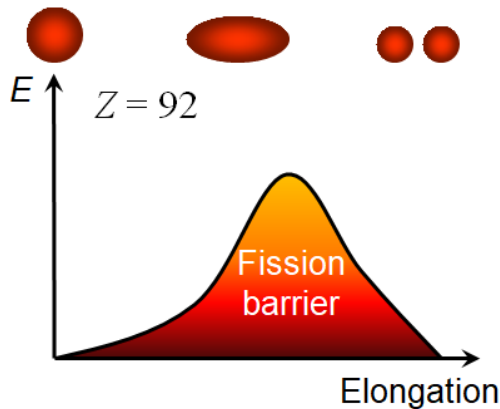
- ❑ Macroscopic-microscopic models
- ❑ Self-consistent approaches
 - Non-relativistic
 - Relativistic
- ❑ ...

What are “Superheavy Nuclei” ?

- ❑ Can not exist if described as a charged liquid drop
- ❑ Stabilized by quantum shell effects

What are superheavy nuclei?

As charged liquid drops, no existence of nuclei w/ $Z > 104$



SHE/SHN

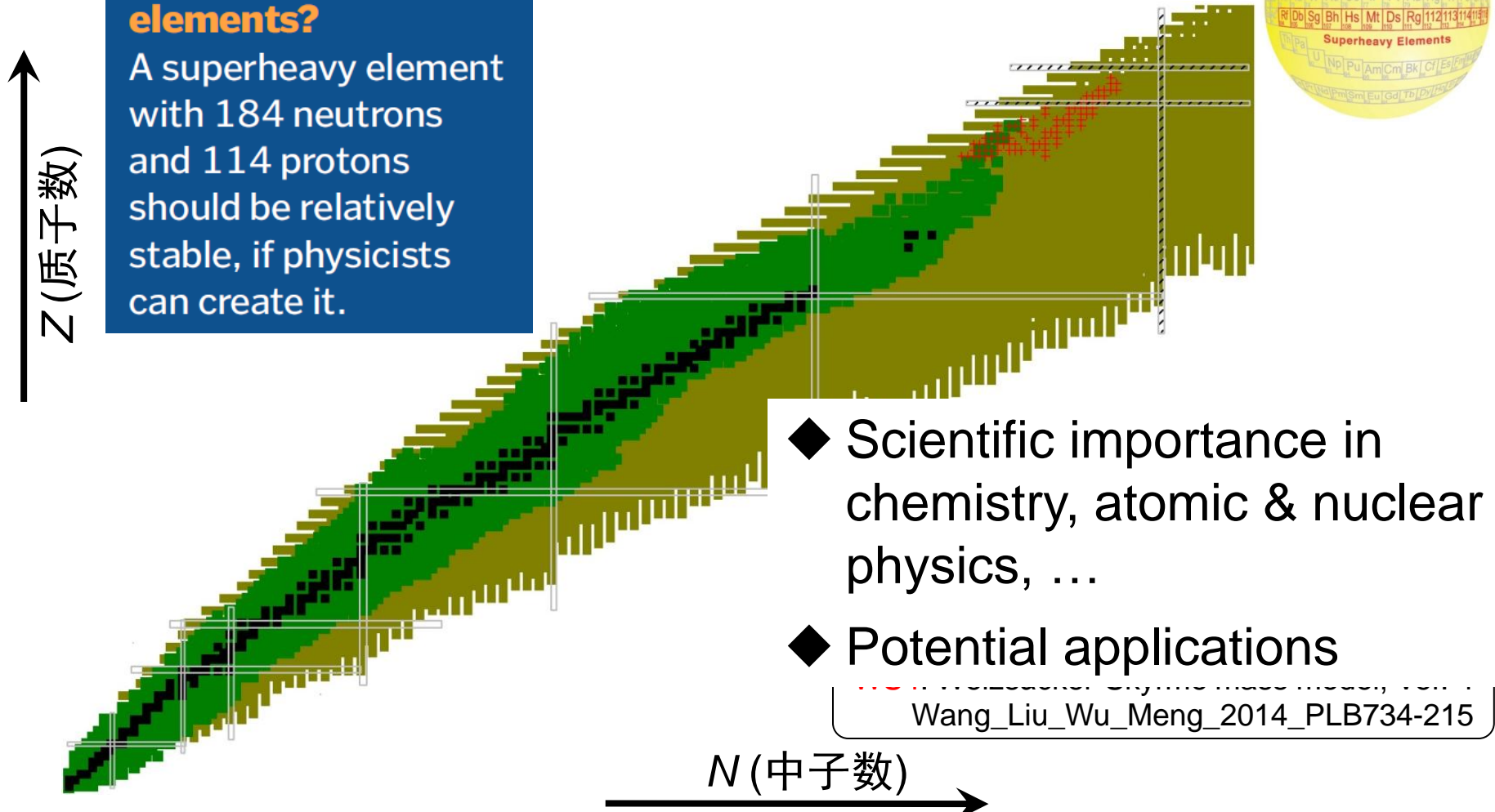
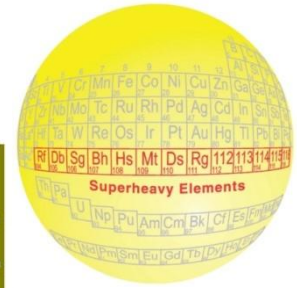
Based on quantum shell effects, predictions made of an island of stability of SHN around $(Z, N) = (114, 184)$

SHN: Charge & mass limits of nuclear existence

Are there stable high-atomic-number elements?

A superheavy element with 184 neutrons and 114 protons should be relatively stable, if physicists can create it.

Top 125 science questions
Science July 2005



- ◆ Scientific importance in chemistry, atomic & nuclear physics, ...
- ◆ Potential applications

Periodic table of elements: End?

	IA											0						
1	¹ 氢 H											² 氦 He						
2	³ 锂 Li	⁴ 铍 Be											⁵ 硼 B	⁶ 碳 C	⁷ 氮 N	⁸ 氧 O	⁹ 氟 F	¹⁰ 氖 Ne
3	¹¹ 钠 Na	¹² 镁 Mg	IIIB	IVB	VB	VIB	VIIB	VIII		IB	IIB	¹³ 铝 Al	¹⁴ 硅 Si	¹⁵ 磷 P	¹⁶ 硫 S	¹⁷ 氯 Cl	¹⁸ 氩 Ar	
4	¹⁹ 钾 K	²⁰ 钙 Ca	²¹ 钪 Sc	²² 钛 Ti	²³ 钒 V	²⁴ 铬 Cr	²⁵ 锰 Mn	²⁶ 铁 Fe	²⁷ 钴 Co	²⁸ 镍 Ni	²⁹ 铜 Cu	³⁰ 锌 Zn	³¹ 镓 Ga	³² 锗 Ge	³³ 砷 As	³⁴ 硒 Se	³⁵ 溴 Br	³⁶ 氪 Kr
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6	⁵⁵ 铯 Cs	⁵⁶ 钡 Ba	57-71 镧系 La-Lu	⁷² 铪 Hf	⁷³ 钽 Ta	⁷⁴ 钨 W	⁷⁵ 铼 Re	⁷⁶ 锇 Os	⁷⁷ 铱 Ir	⁷⁸ 铂 Pt	⁷⁹ 金 Au	⁸⁰ 汞 Hg	⁸¹ 铊 Tl	⁸² 铅 Pb	⁸³ 铋 Bi	⁸⁴ 钋 Po	⁸⁵ 砹* At	⁸⁶ 氡* Rn
7	⁸⁷ 钫* Fr	⁸⁸ 镭* Ra	89-103 锕系 Ac-Lr	¹⁰⁴ 𬬻* Rf	¹⁰⁵ 𬬺* Db	¹⁰⁶ 𬬻* Sg	¹⁰⁷ 𬬾* Bh	¹⁰⁸ 𬬽* Hs	¹⁰⁹ 𬬼* Mt	¹¹⁰ 𬬻* Ds	¹¹¹ 𬬺* Rg	¹¹² 𬬻* Cn	¹¹³ Uut* Uut	¹¹⁴ Uuq* Fl	¹¹⁵ Uup* Uup	¹¹⁶ Uuq* Lv	¹¹⁷ Uus* Uus	¹¹⁸ Uuo* Uuo

带*为放射性元素

其中黑色为天然放射性元素

红色为人造元素

蓝色为新命名元素

SGZ, SHN & SHE, Physics 43 (2014) 817-825 (in Chinese)

镧系	⁵⁷ 镧 La	⁵⁸ 铈 Ce	⁵⁹ 镨 Pr	⁶⁰ 钕 Nd	⁶¹ 钷* Pm	⁶² 钐 Sm	⁶³ 铕 Eu	⁶⁴ 钆 Gd	⁶⁵ 铽 Tb	⁶⁶ 镝 Dy	⁶⁷ 钬 Ho	⁶⁸ 铒 Er	⁶⁹ 铥 Tm	⁷⁰ 镱 Yb	⁷¹ 镱 Lu
锕系	⁸⁹ 锕* Ac	⁹⁰ 钍* Th	⁹¹ 镤* Pa	⁹² 铀* U	⁹³ 镎* Np	⁹⁴ 钚* Pu	⁹⁵ 镅* Am	⁹⁶ 锔* Cm	⁹⁷ 锿* Bk	⁹⁸ 镆* Cf	⁹⁹ 锿* Es	¹⁰⁰ 镆* Fm	¹⁰¹ 镎* Md	¹⁰² 锘* No	¹⁰³ 铹* Lr

Courtesy of Xu Meng (孟旭)

Lecture 1

- Predictions of the island of stability of SHN
- Experimental progress on synthesis of SHN
- Challenges in synthesizing SHN

Expt. exploration of SHN & island of stability

□ If half-life of SHN long enough (e.g, $T_{1/2} \sim 10^8$) & produced in nucleosynthesis

Herrmann1979_Nature280-543

➤ SHE/SHN may exist in Nature

➤ No confirmed evidence, but efforts still being made

- ✓ 1860 (1861), Bunsen & Kirchhoff found cesium (rubidium) in mineral water (lepidolite) by using the flame spectroscopy
- ✓ 1868, Janssen & Lockyer recorded helium spectral line during solar eclipse
- ✓ ...
- ✓ 1898, Curies discovered polonium (radium) in uranium ore pitchblende (uraninite) by identifying strong radioactivity

Expt. exploration of SHN & island of stability

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Herrmann1979_Nature280-543

- SHE/SHN may exist in Nature
- No confirmed evidence, but efforts still being made

- In Lab, via heavy ion fusion reactions

- GSI in Darmstadt, Germany
- Flerov Laboratory of Nuclear Reactions in Dubna, Russia
- Lawrence Berkeley National Laboratory, USA
- Lawrence Livermore National Laboratory, USA
- RIKEN in Wako, Japan
- GANIL in Caen, France

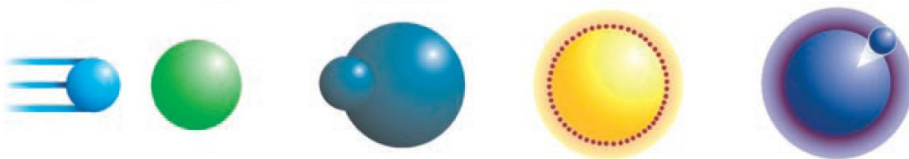
Hofmann_Münzenberg2000_RMP72-733

Morita...2004_JPSJ73-2593

Oganessian...2007_JPG34-R165

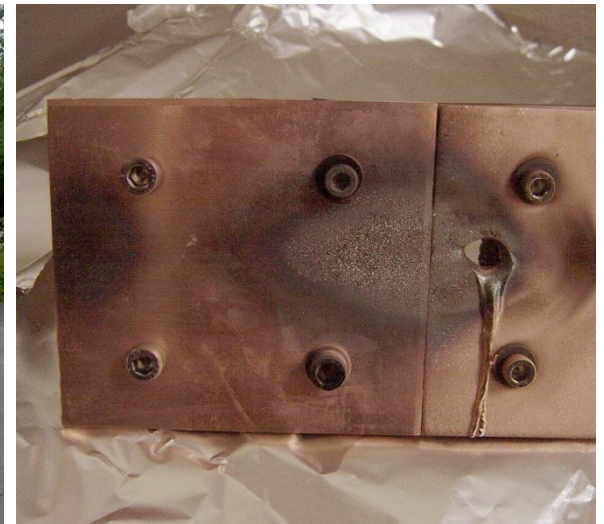
Oganessian...2010_PRL104-142502

Zhang...2012_CPL29-012502

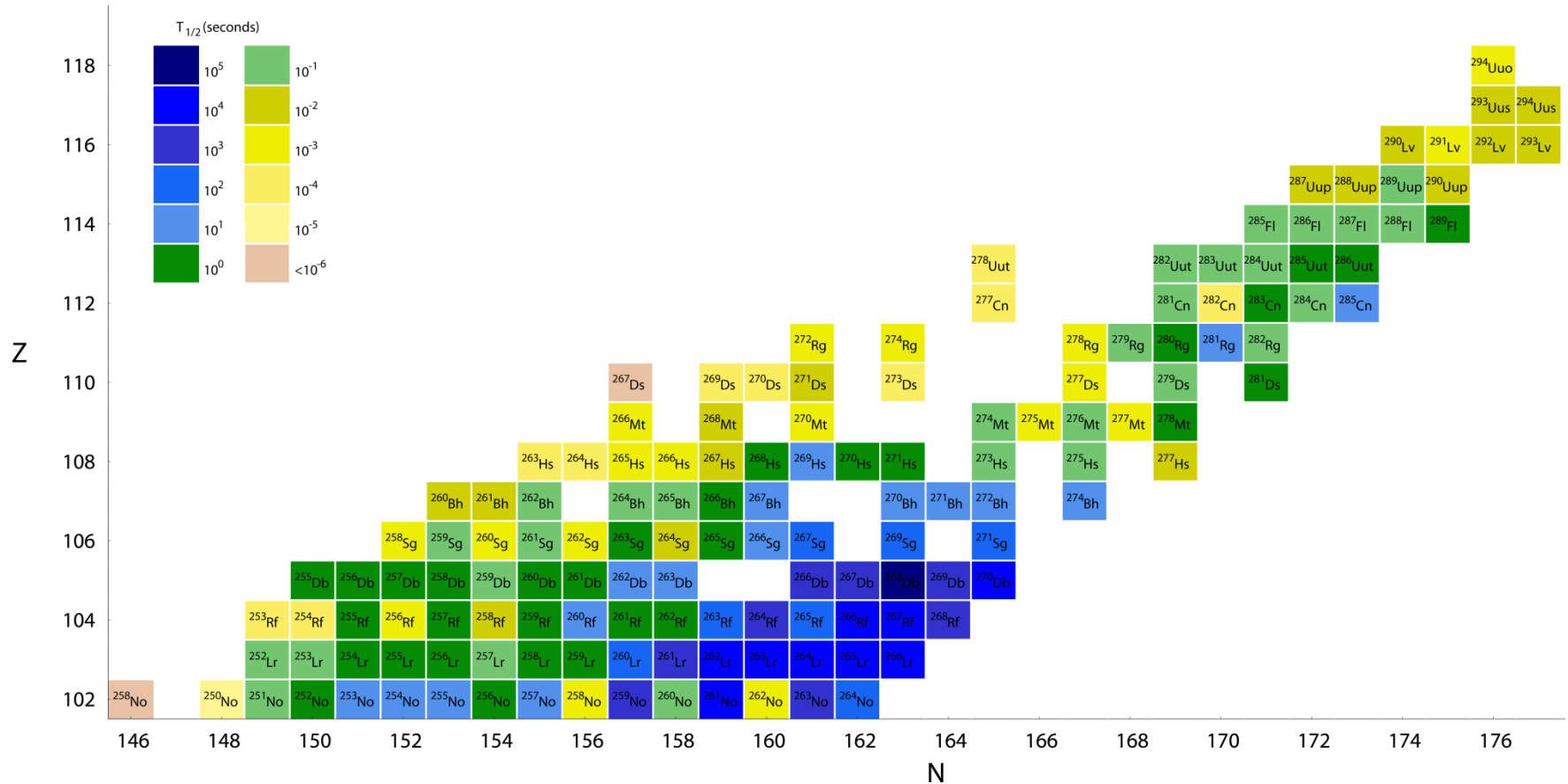


Synthesis of SHN in Labs

- ❑ Neutron capture followed by β decay
 - Up to fermium
- ❑ Light ions (H, D, T & α) as projectiles
 - Up to mendelevium
- ❑ Heavy ions as projectiles
 - **Cold fusion:** ^{208}Pb or ^{209}Bi as targets, up to nihonium
 - **Hot fusion:** ^{48}Ca as projectiles, up to oganesson



Expt. progress on synthesis of SHN



SHE with $Z \leq 118$ have been synthesized & named
Recent naming for elements with $Z = 113, 115, 117$ & 118

Periodic table of elements: End?

	IA																	0						
1	¹ 氢 H																	² 氦 He						
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其中黑色为天然放射性元素

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SGZ, SHN & SHE, Physics 43 (2014) 817-825 (in Chinese)

镧系	⁵⁷ 镧 La	⁵⁸ 铈 Ce	⁵⁹ 镨 Pr	⁶⁰ 钕 Nd	⁶¹ 钷* Pm	⁶² 钐 Sm	⁶³ 铕 Eu	⁶⁴ 钆 Gd	⁶⁵ 铽 Tb	⁶⁶ 镝 Dy	⁶⁷ 钬 Ho	⁶⁸ 铒 Er	⁶⁹ 铥 Tm	⁷⁰ 镱 Yb	⁷¹ 镱 Lu
锕系	⁸⁹ 锕* Ac	⁹⁰ 钍* Th	⁹¹ 镤* Pa	⁹² 铀* U	⁹³ 镎* Np	⁹⁴ 钚* Pu	⁹⁵ 镅* Am	⁹⁶ 锔* Cm	⁹⁷ 锿* Bk	⁹⁸ 镆* Cf	⁹⁹ 锘* Es	¹⁰⁰ 镆* Fm	¹⁰¹ 钔* Md	¹⁰² 锘* No	¹⁰³ 铹* Lr

Courtesy of Xu Meng (孟旭)

Element 113

JWP ASSESSMENT: Three chains of $^{278}113$ observed by the RIKEN collaborations, the first in 2004 [9], the second in 2007 [10], and the third in 2012 [16], are now construed as being consistent. Firm connection to established nuclides is provided. The remaining criterion achieved for acknowledgement of discovery is an identification of Z which is now embodied in the cross reaction production and characterization of the chain beginning with ^{266}Bh as found by the RIKEN collaboration in 2009 [14] and by Qin *et al.* in 2006 [12]. The Criteria for discovery have been met.

Karol+2016_PureApplChem88-139

Morita+2004

[9] K. Morita, K. Morimoto, D. Kaji, T. Akiyama, S. Goto, H. Haba, E. Ideguchi, R. Kanungo, K. Katori, H. Koura, H. Kudo, T. Ohnishi, A. Ozawa, T. Suda, K. Sueki, H. Xu, T. Yamaguchi, A. Yoneda, A. Yoshida, Y.-L. Zhao. *J. Phys. Soc. Jpn.* **73**, 2593 (2004).

Morita+2007

[10] K. Morita, K. Morimoto, D. Kaji, T. Akiyama, S. Goto, H. Haba, E. Ideguchi, K. Katori, H. Koura, H. Kikunaga, H. Kudo, T. Ohnishi, A. Ozawa, N. Sato, T. Suda, K. Sueki, F. Tokanai, T. Yamaguchi, A. Yoneda, A. Yoshida. *J. Phys. Soc. Jpn.* **76**, 045001 (2007).

Qin+2006

[11] P. A. Wilk, K. E. Gregorich, A. Türler, C. A. Laue, R. Eichler, V. Ninov, J. L. Adams, U. W. Kirbach, M. R. Lane, D. M. Lee, J. B. Patin, D. A. Shaughnessy, D. A. Strellis, H. Nitsche, D. C. Hoffman. *Phys. Rev. Lett.* **85**, 2697 (2000).

[12] Z. Qin, X. Wu, H. Ding, W. Wu, W. Huang, X. Lei, Y. Xu, X. Yuan, B. Guo, W. Yang, Z. Gan, H. Fan, J. Guo, H. Xu, G. Xiao. *Nucl. Phys. Rev.* **23**, 400 (2006).

Morita+2009

[13] Yu. Ts. Oganessian, V. K. Utyonkov, Yu. V. Lobanov, F. Sh. Abdullin, A. N. Polyakov, R. N. Sagaidak, I. V. Shirokovsky, Yu. S. Tsyganov, A. A. Voinov, G. G. Gulbekian, S. L. Bogomolov, B. N. Gikal, A. N. Mezentsev, V. G. Subbotin, A. M. Sukhov, K. Subotic, V. I. Zagrebaev, G. K. Vostokin, M. G. Itkis, R. A. Henderson, J. M. Kenneally, J. H. Landrum, K. J. Moody, D. A. Shaughnessy, M. A. Stoyer, N. J. Stoyer, P. A. Wilk. *Phys. Rev. C* **76**, 011601(R) (2007).

[14] K. Morita, K. Morimoto, D. Kaji, H. Haba, K. Ozeki, Y. Kudou, N. Sato, T. Sumita, A. Yoneda, T. Ichikawa, Y. Fujimori, S. Goto, E. Ideguchi, Y. Kasamatsu, K. Katori, Y. Komori, H. Koura, H. Kudo, K. Ooe, A. Ozawa, F. Tokanai, K. Tsukada, T. Yamagichi, A. Yoshida. *J. Phys. Soc. Jpn.* **78**, 064201 (2009).

Morita+2012

[15] Y. A. Akevali. *Nucl. Data Sheets* **94**, 131 (2001).

[16] K. Morita, K. Morimoto, D. Kaji, H. Haba, K. Ozeki, Y. Kudou, T. Sumita, Y. Wakabayashi, A. Yoneda, K. Tanaka, S. Yamaki, R. Sakai, T. Akiyama, S. Goro, H. Hasebe, M. Huang, T. Huang, E. Ideguchi, Y. Kasamatsu, K. Katori, Y. Kariya, H. Kikunaga, H. Koura, H. Kudo, A. Mashiko, K. Mayama, S. Mitsuoka, T. Moriya, M. Murakami, H. Murayama, S. Namai, A. Ozawa, N. Sato, K. Sueki, M. Takeyajma, F. Tokanai, T. Yamaguchi, A. Yoshida. *J. Phys. Soc. Jpn.* **81**, 103201 (2012).

Periodic table of elements: End?

	IA																	0
1	1 氢 H																	2 氦 He
2	3 锂 Li	4 铍 Be																10 氖 Ne
3	11 钠 Na	12 镁 Mg																18 氩 Ar
4	19 钾 K	20 钙 Ca																36 氪 Kr
5	37 铷 Rb	38 锶 Sr																54 氙 Xe
6	55 铯 Cs	56 钡 Ba	57 镧系 La-Lu															86 氡 Rn
7	87 钫* Fr	88 镭* Ra	89-103 镧系 Ac-Lr	104 𬬻* Rf	105 𬬾* Db	106 𬬿* Sg	107 𬬰* Bh	108 𬬱* Hs	109 𬬲* Mt	110 𬬳* Ds	111 𬬴* Rg	112 𬬵* Cn	113 𬬶* Nh	114 𬬷* Fl	115 𬬸* Mc	116 𬬹* Lv	117 𬬺* Ts	118 𬬻* Og

带*为放射性元素

其中黑色为天然放射性元素

红色为人造元素

113: Nihonium (Nh, 铈) 2004-2006-2007-2009-2012

115: Moscovium (Mc, 镆) 2004-2010-2013

117: Tennessine (Ts, 𫟼) 2010-2012-2013

118: Oganesson (Og, 𫟽) 2006-2012

SGZ, SHN & New Elements, Nucl. Phys. Rev. 34 (2017) 318-331 (in, Chinese)

镧系	57 镧 La	58 铈 Ce	59 镨 Pr	60 钕 Nd	61 钷* Pm	62 钐 Sm	63 铕 Eu	64 钆 Gd	65 铽 Tb	66 镝 Dy	67 钬 Ho	68 铒 Er	69 铥 Tm	70 镱 Yb	71 镱 Lu
锕系	89 锕* Ac	90 钍* Th	91 镤* Pa	92 铀* U	93 镎* Np	94 钚* Pu	95 镅* Am	96 锔* Cm	97 锫* Bk	98 锿* Cf	99 镄* Es	100 镆* Fm	101 钔* Md	102 镎* No	103 铹* Lr

Courtesy of Xu Meng (孟旭)

Periodic table of elements: End?

带*为放射性元素
其中黑色为天然放射性元素
红色为人造元素

	IA																		0	
1	¹ 氢 H																			² 氦 He
2	³ 锂 Li	⁴ 铍 Be																		¹⁰ 氖 Ne
3	¹¹ 钠 Na	¹² 镁 Mg																		¹⁸ 氩 Ar
4	¹⁹ 钾 K	²⁰ 钙 Ca	²¹ 钪 Sc	²² 钛 Ti	²³ 钒 V	²⁴ 铬 Cr	²⁵ 锰 Mn	²⁶ 铁 Fe	²⁷ 钴 Co	²⁸ 镍 Ni	²⁹ 铜 Cu	³⁰ 锌 Zn	³¹ 镓 Ga	³² 锗 Ge	³³ 砷 As	³⁴ 硒 Se	³⁵ 溴 Br		³⁶ 氪 Kr	
5	³⁷ 铷 Rb	³⁸ 锶 Sr	³⁹ 钇 Y	⁴⁰ 锆 Zr	⁴¹ 铌 Nb	⁴² 钼 Mo	⁴³ 锝* Tc	⁴⁴ 钌 Ru	⁴⁵ 铑 Rh	⁴⁶ 钯 Pd	⁴⁷ 银 Ag	⁴⁸ 镉 Cd	⁴⁹ 铟 In	⁵⁰ 锡 Sn	⁵¹ 锑 Sb	⁵² 碲 Te	⁵³ 碘 I		⁵⁴ 氙 Xe	
6	⁵⁵ 铯 Cs	⁵⁶ 钡 Ba	⁵⁷⁻⁷¹ 镧系 La-Lu	⁷² 铪 Hf	⁷³ 钽 Ta	⁷⁴ 钨 W	⁷⁵ 铼 Re	⁷⁶ 锇 Os	⁷⁷ 铱 Ir	⁷⁸ 铂 Pt	⁷⁹ 金 Au	⁸⁰ 汞 Hg	⁸¹ 铊 Tl	⁸² 铅 Pb	⁸³ 铋 Bi	⁸⁴ 钋 Po	⁸⁵ 砹* At		⁸⁶ 氡* Rn	
7	⁸⁷ 钫* Fr	⁸⁸ 镭* Ra	⁸⁹⁻¹⁰³ 锕系 Ac-Lr	¹⁰⁴ 𬬻* Rf	¹⁰⁵ 𬬼* Db	¹⁰⁶ 𬬽* Sg	¹⁰⁷ 𬬾* Bh	¹⁰⁸ 𬬿* Hs	¹⁰⁹ 𬭀* Mt	¹¹⁰ 𬬻* Ds	¹¹¹ 𬬼* Rg	¹¹² 𬬽* Cn	¹¹³ 𬬾* Nh	¹¹⁴ 𬬿* Fl	¹¹⁵ 𬭀* Mc	¹¹⁶ 𬬻* Lv	¹¹⁷ 𬬼* Ts		¹¹⁸ 𬬽* Og	

SGZ, SHN & New Elements, Nucl. Phys. Rev. 34 (2017) 318-331 (in, Chinese)

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锕系	⁸⁹ 锕* Ac	⁹⁰ 钍* Th	⁹¹ 镤* Pa	⁹² 铀* U	⁹³ 镎* Np	⁹⁴ 钚* Pu	⁹⁵ 镅* Am	⁹⁶ 锔* Cm	⁹⁷ 锇* Bk	⁹⁸ 锿* Cf	⁹⁹ 镅* Es	¹⁰⁰ 镆* Fm	¹⁰¹ 钔* Md	¹⁰² 锘* No	¹⁰³ 铱* Lr

Courtesy of Xu Meng (孟旭)

Chinese names of new elements

□ 113: Nihonium (Nh)

➤ Proposals of Chinese name: 铈, 铊, 铋 & 铌

□ 115: Moscovium (Mc)

➤ Proposal of Chinese name: 镆

□ 117: Tennessine (Ts)

➤ Proposals of Chinese name: 砹 & 碛

□ 118: Oganesson (Og)

➤ Proposals of Chinese name: 奥 & 𫖇



Chinese names of new elements

核データニュース, No.118 (2017)

話題・解説 (I)

Chinese Names of New Elements with $Z = 113, 115, 117$ & 118

Shan-Gui Zhou (周善贵/周善貴)¹

Institute of Theoretical Physics, Chinese Academy of Sciences, Beijing 100190, China

(中国科学院理论物理研究所/中國科學院理論物理研究所)

sgzhou@itp.ac.cn

HIRFL China: ^{271}Ds ($Z = 110$)

2011.01.15 $^{64}\text{Ni} + ^{208}\text{Pb} \rightarrow ^{272}\text{Ds}^*$ 7 days

2011.03.15 $^{64}\text{Ni} + ^{208}\text{Pb} \rightarrow ^{272}\text{Ds}^*$ 13 days

IMP/CAS, ITP/CAS

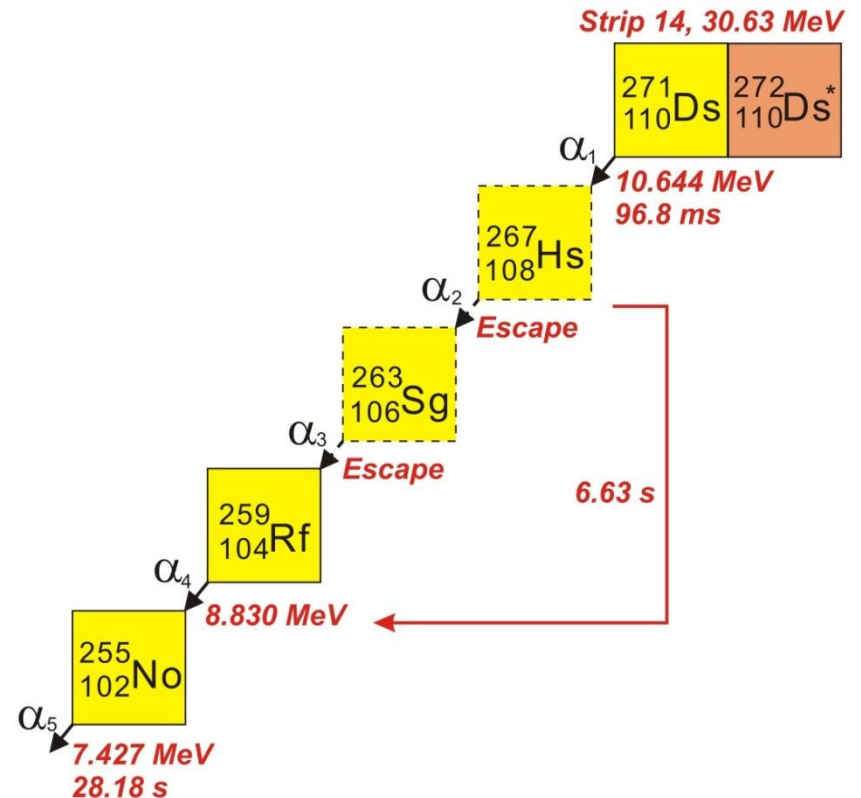
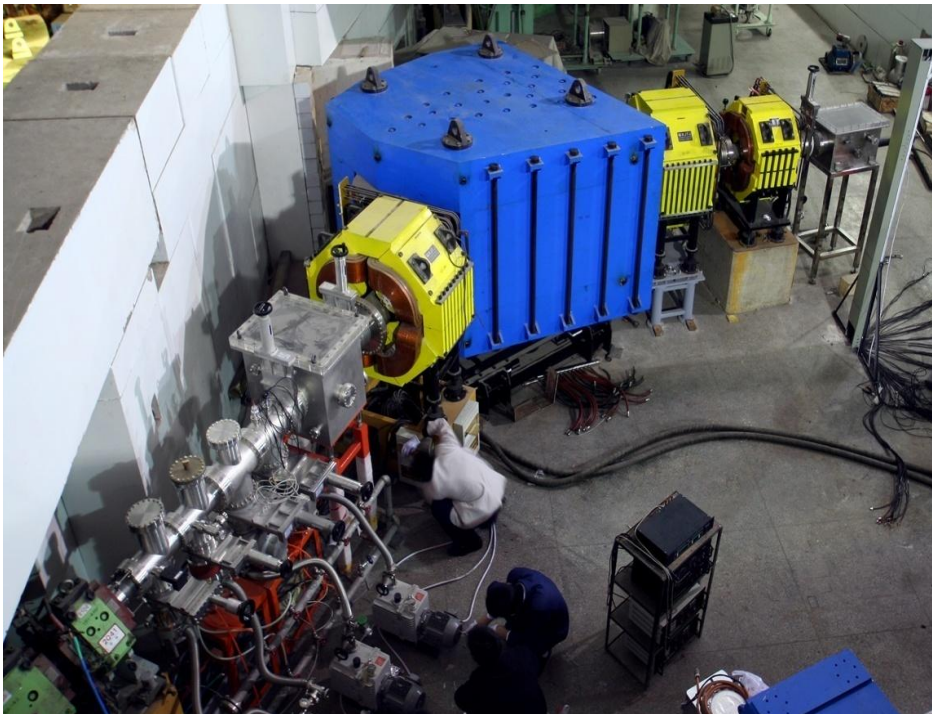
Nanjing Univ, CIAE

Zhang, Gan, Ma ...

Chinese Physics Letters 29 (2012) 012502

2001: ^{259}Db ($Z=105$)

2004: ^{265}Bh ($Z=107$)



Courtesy of Zai-Guo Gan (甘再国)

Lecture 1

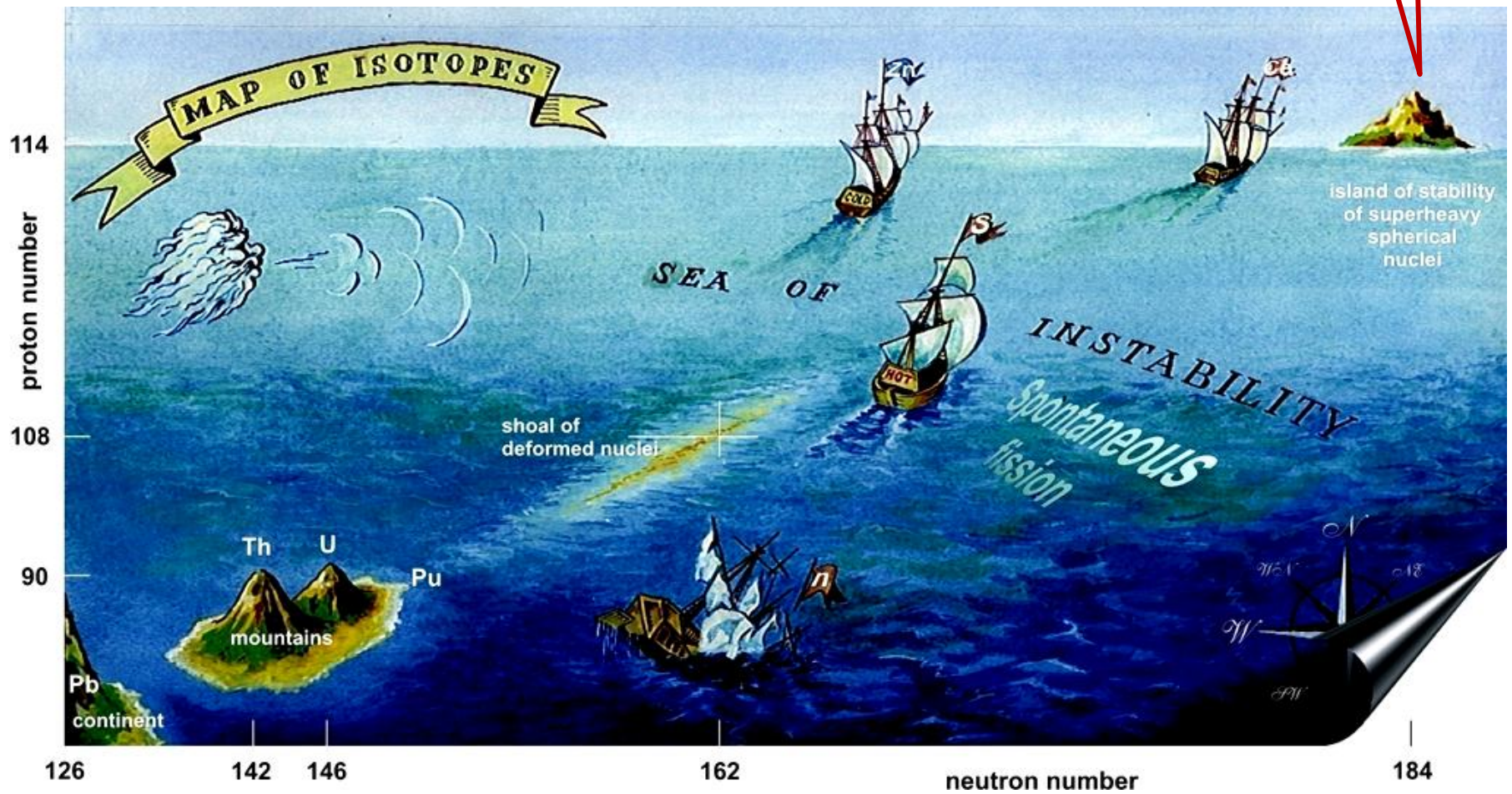
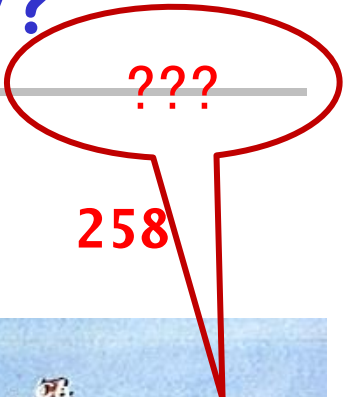
- Predictions of the island of stability of SHN
- Experimental progress on synthesis of SHN
- Challenges in synthesizing SHN

Challenges to synthesize SHN & new SHEs

- Elusive island of stability

Where is the island of stability?

Z = ?	114	116	120	126	132	138		
N = ?	172	176	178	184	198	228	238	258



Challenges to synthesize SHN & new SHEs

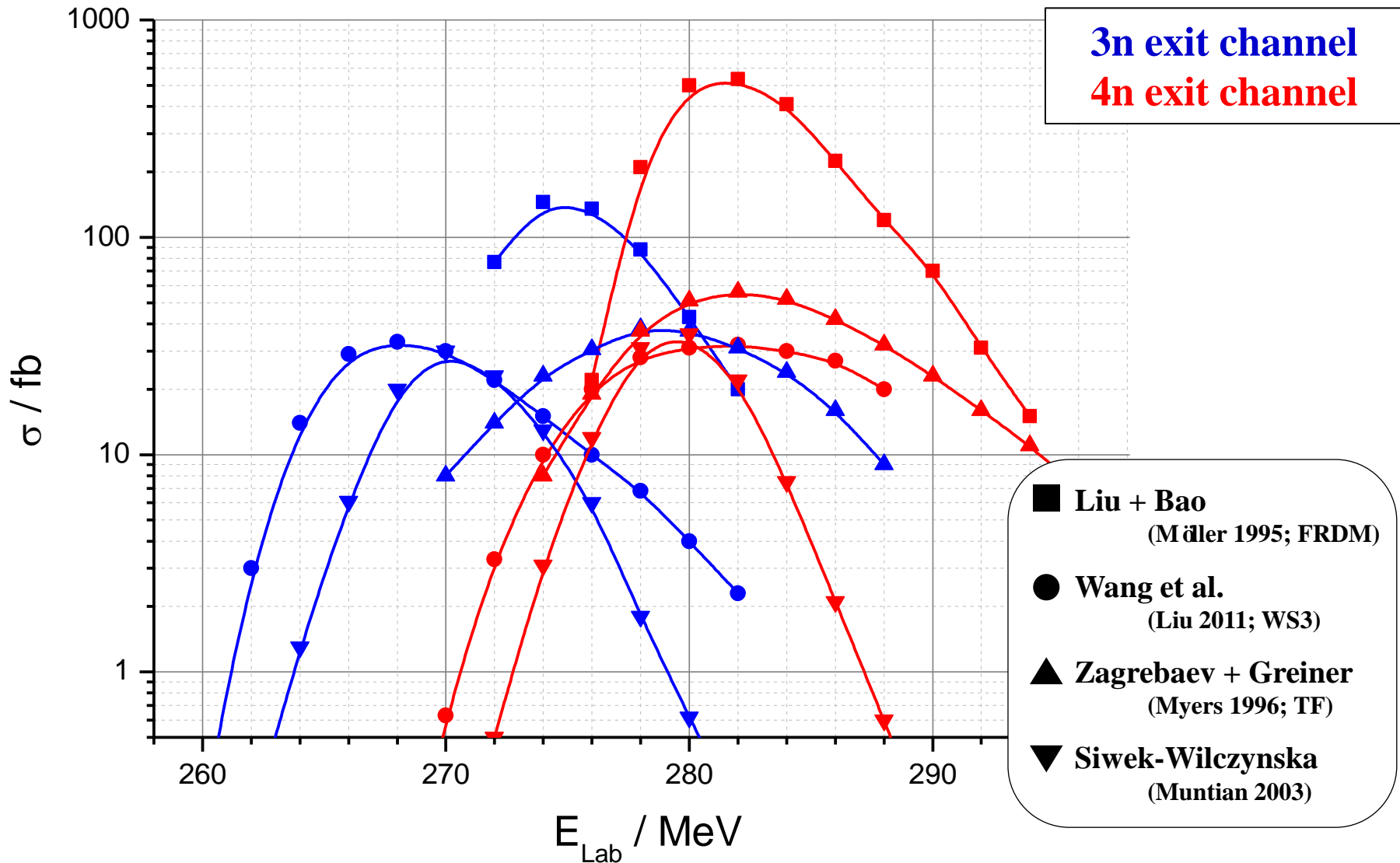
□ Elusive island of stability

□ Tiny cross sections w/ huge uncertainties

➤ ^{271}Ds : $\sigma \sim 10$ pb; HIRFL: 1 event in 20 days

➤ ^{278}Nh : $\sigma \sim 0.02$ pb; RIKEN: 3 events in 553 days

Uncertainties in predicted xsections for $Z = 119$

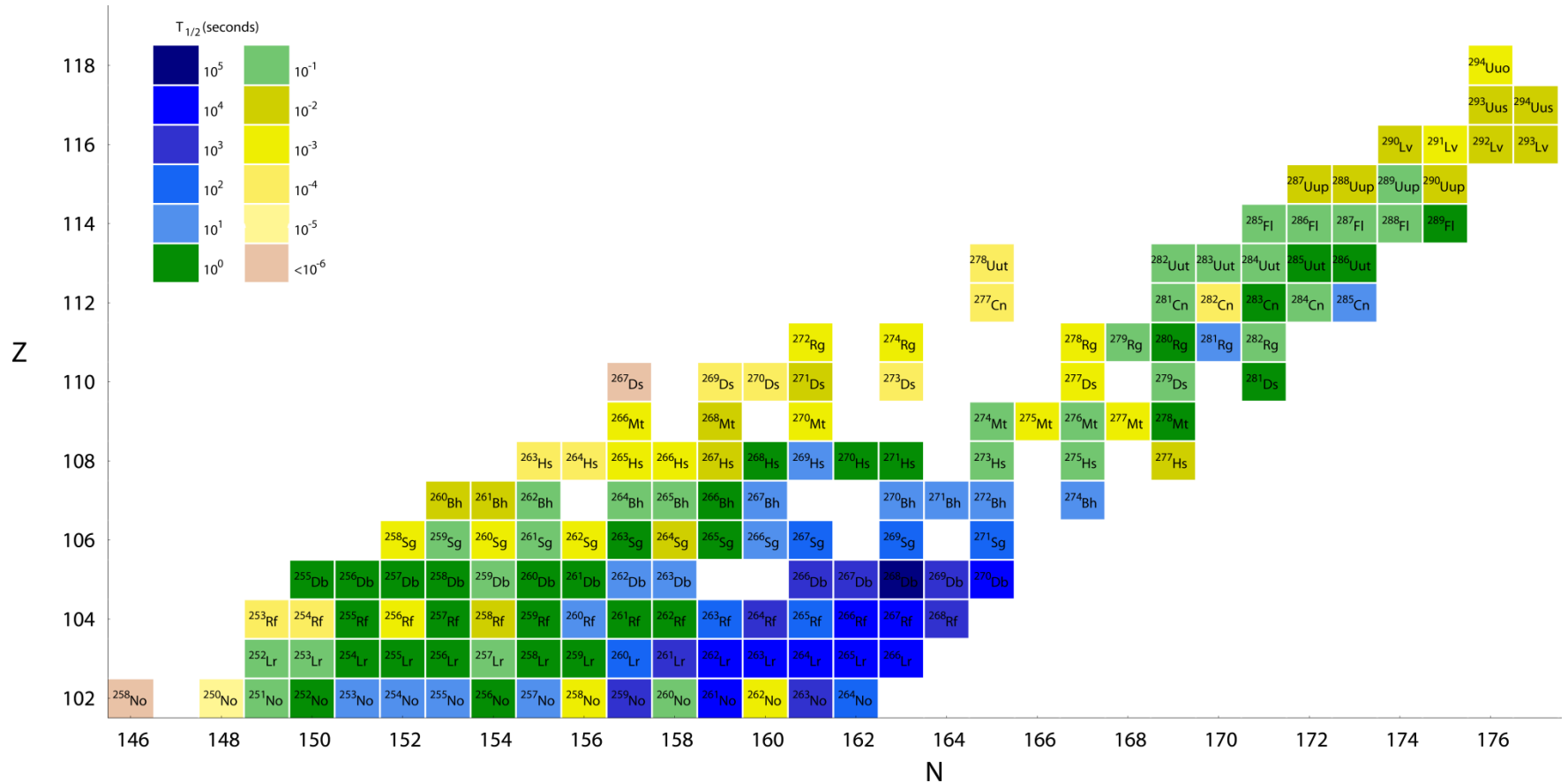


From Ch. E. Duellman (FUSHE2012)

Challenges to synthesize SHN & new SHEs

- Elusive island of stability
- Tiny cross sections w/ huge uncertainties
 - ^{271}Ds : $\sigma \sim 10$ pb; HIRFL: 1 event in 20 days
 - ^{278}Nh : $\sigma \sim 0.02$ pb; RIKEN: 3 events in 553 days
- Only neutron-deficient SHN, far away from $N = 184$

Only neutron deficient SHN produced via HI fusion



Life-time longer w/ increasing N !

Lecture 2

- Challenges in synthesizing SHN

- Theoretical study of structure of SHN
 - Nuclear models
 - Next shell closures beyond ^{208}Pb as seen from single particle spectra, shell correction energy & nuclear shapes
 - Exotic shapes in SHN
 - Low-lying spectra of SHN & magicities

- Theoretical study of decay of SHN

Study of SHN: Theory & experiment

- Theory and experiment has progressed hand in hand
 - 1960s, heavy ion accelerators & detectors built for SHN following theoretical predictions; 1980s, important progress made in experiments
 - ...
 - 2010, element with $Z = 117$

PRL **104**, 142502 (2010)

 Selected for a *Viewpoint* in *Physics*
PHYSICAL REVIEW LETTERS

week ending
9 APRIL 2010



Synthesis of a New Element with Atomic Number $Z = 117$

Yu. Ts. Oganessian,^{1,*} F. Sh. Abdullin,¹ P. D. Bailey,² D. E. Benker,² M. E. Bennett,³ S. N. Dmitriev,¹ J. G. Ezold,² J. H. Hamilton,⁴ R. A. Henderson,⁵ M. G. Itkis,¹ Yu. V. Lobanov,¹ A. N. Mezentsev,¹ K. J. Moody,⁵ S. L. Nelson,⁵ A. N. Polyakov,¹ C. E. Poit,² R. N. Sagaidak,¹ D. A. Shaughnessy,² Yu. S. Tsybenko,¹ V. I. Zinchenko,¹ and S. N. Dmitriev,¹

[12] C. Shen *et al.*, *Int. J. Mod. Phys. E* **17**, 66 (2008).
[13] V. Zagrebaev *et al.*, *Phys. Rev. C* **78**, 034610 (2008).
[14] Z. H. Liu *et al.*, *Phys. Rev. C* **80**, 034601 (2009).

³University of Nevada Las Vegas, Las Vegas, Nevada 89154, USA

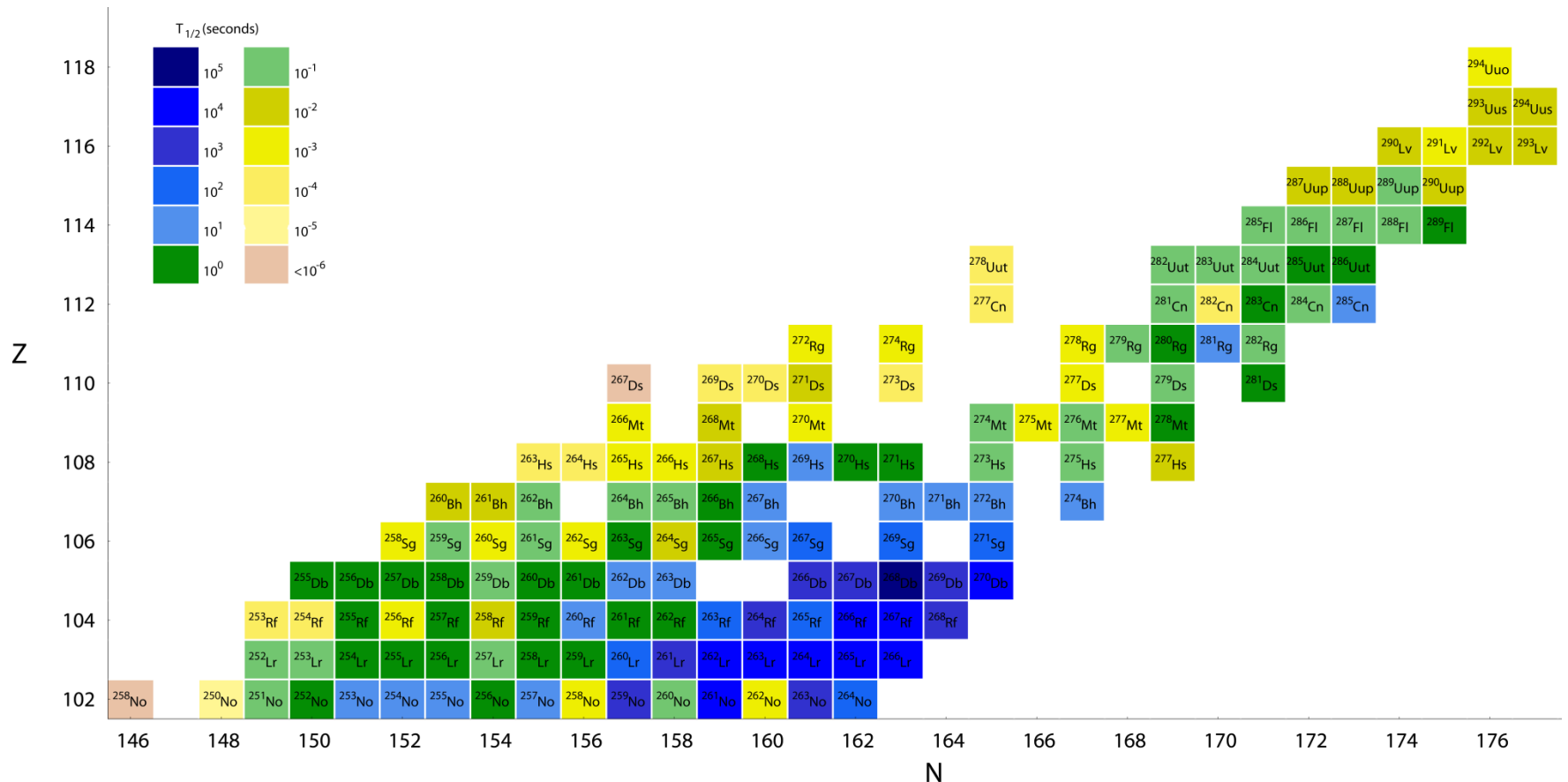
⁴Department of Physics and Astronomy, Vanderbilt University, Nashville, Tennessee 37235, USA

⁵Lawrence Livermore National Laboratory, Livermore, California 94551, USA

⁶Research Institute of Atomic Reactors, RU-433510 Dimitrovgrad, Russian Federation

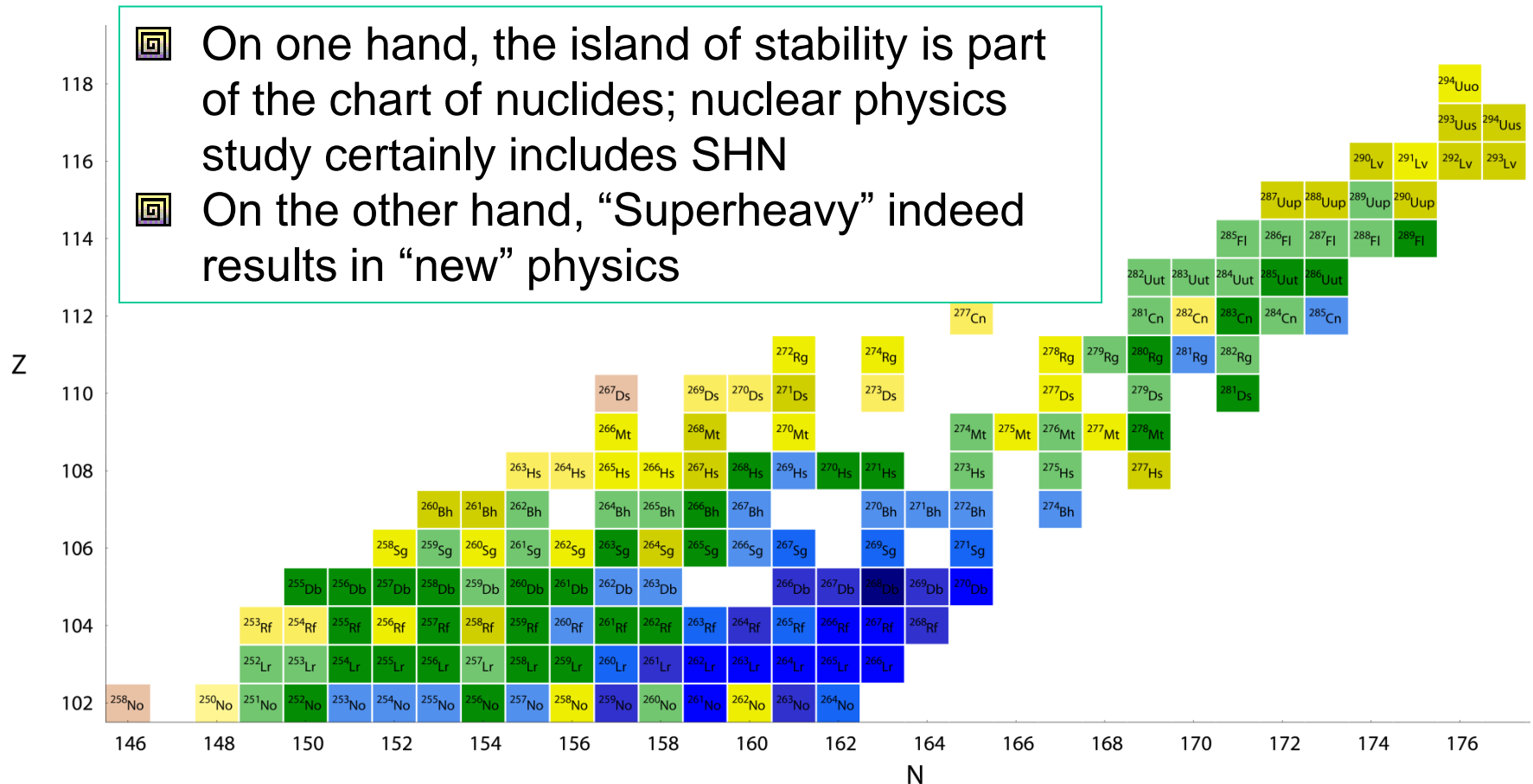
(Received 15 March 2010; published 9 April 2010)

Is there new physics on the island of stability?



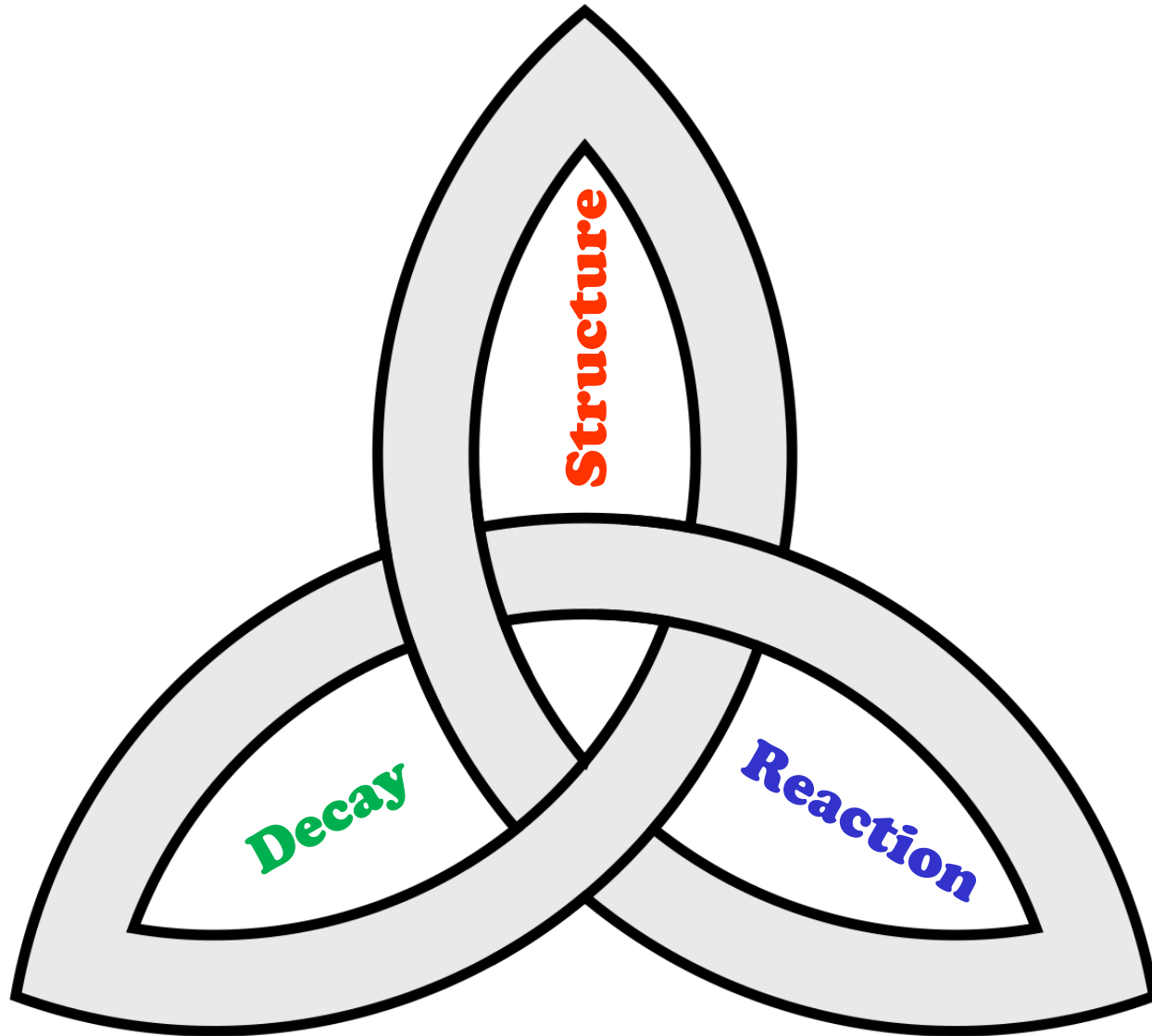
SHE with $Z \leq 118$ have been synthesized & named
Recent naming for elements with $Z = 113, 115, 117 & 118$

Is there new physics on the island of stability?



SHE with $Z \leq 118$ have been synthesized & named
 Recent naming for elements with $Z = 113, 115, 117 \& 118$

Trinity in nucl. phys.: **Structure**, **decay** & **reaction**



Lecture 2

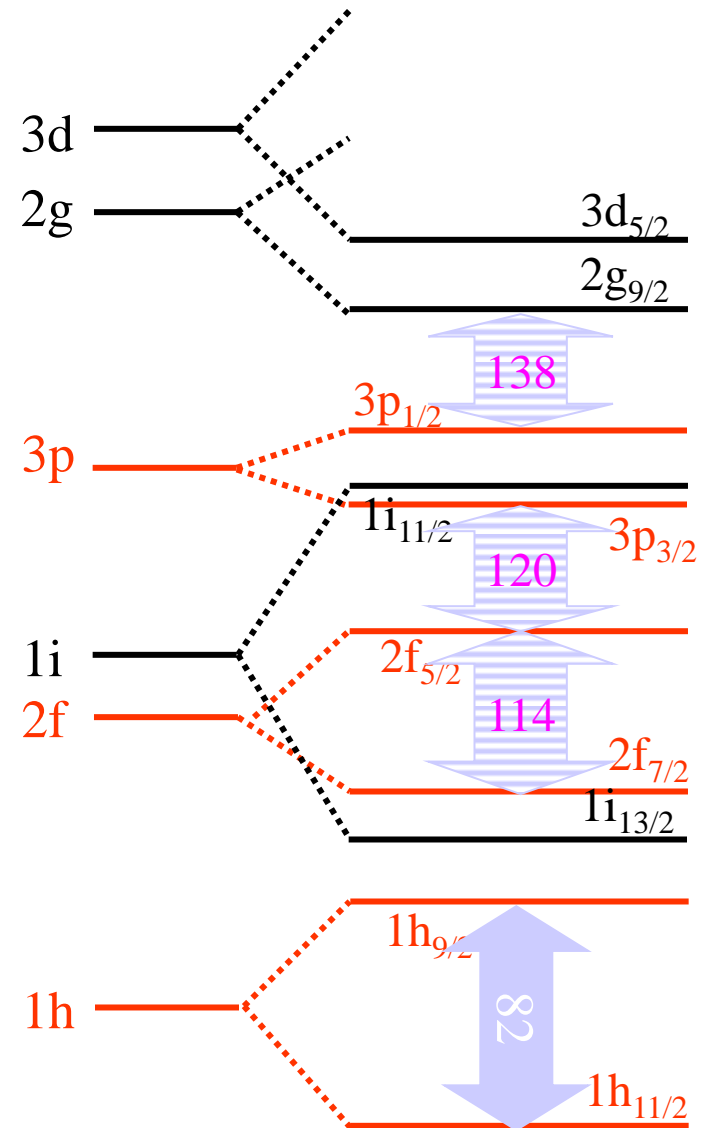
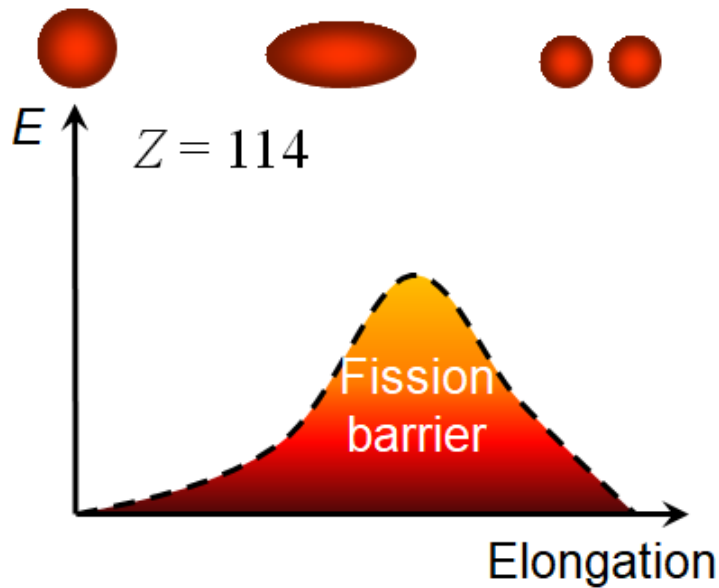
- Challenges in synthesizing SHN

- Theoretical study of structure of SHN
 - Nuclear models
 - Next shell closures beyond ^{208}Pb as seen from single particle spectra, shell correction energy & nuclear shapes
 - Exotic shapes in SHN
 - Low-lying spectra of SHN & magicities

- Theoretical study of decay of SHN

Structure of SHN

- Structure**
- Ground state properties
 - Low-E spectroscopy
 - Fission barrier
 - Isomeric states
 - ...



Lecture 2

- Challenges in synthesizing SHN

- Theoretical study of structure of SHN
 - Nuclear models
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Nuclear models

□ Macroscopic models

- Global behavior & average effects

□ Macroscopic-microscopic models

- Global behavior & average effects
- Correction from single particle motion on whole nucleus

□ Microscopic models

- Starting from nucleons; correlations among nucleons
- NN interaction, mean-field approximation, residual interaction

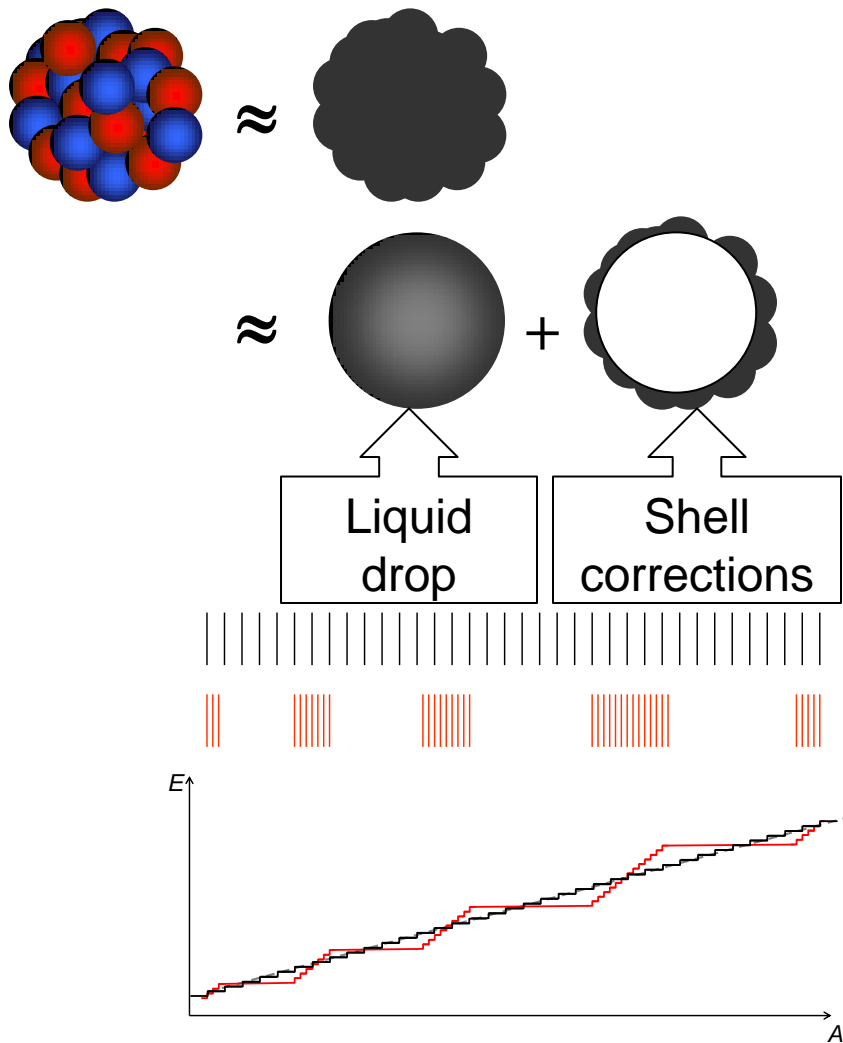
Hu & Zhong, Macroscopic models for nuclei, 1998 (in Chinese)

Greiner & Maruhn, Nuclear models, Springer, 1995

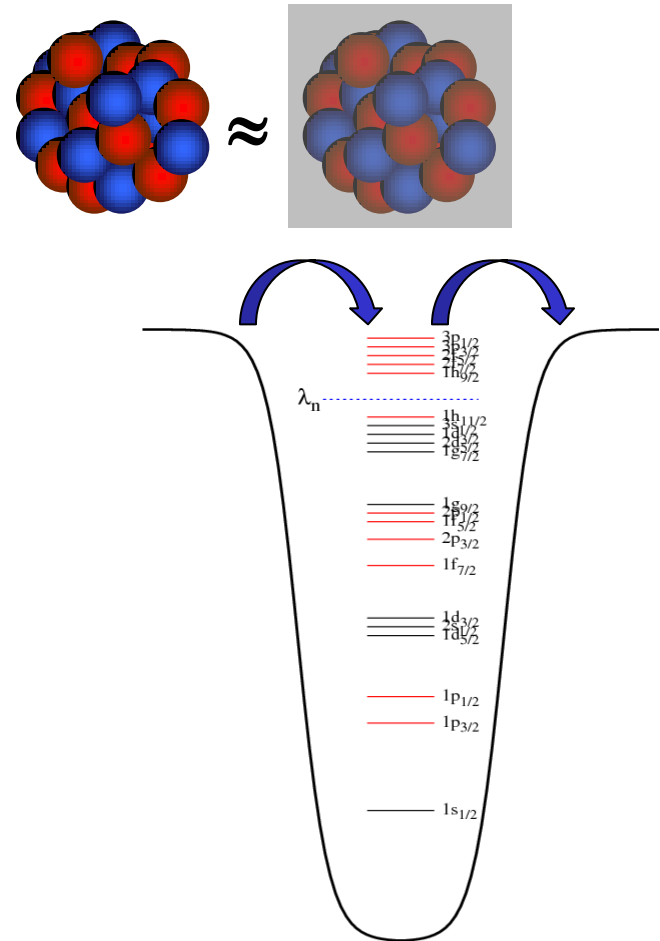
Ring & Schuck, The Nuclear Many-Body Problem, Springer, 1980

Nuclear models

Mac-mic models



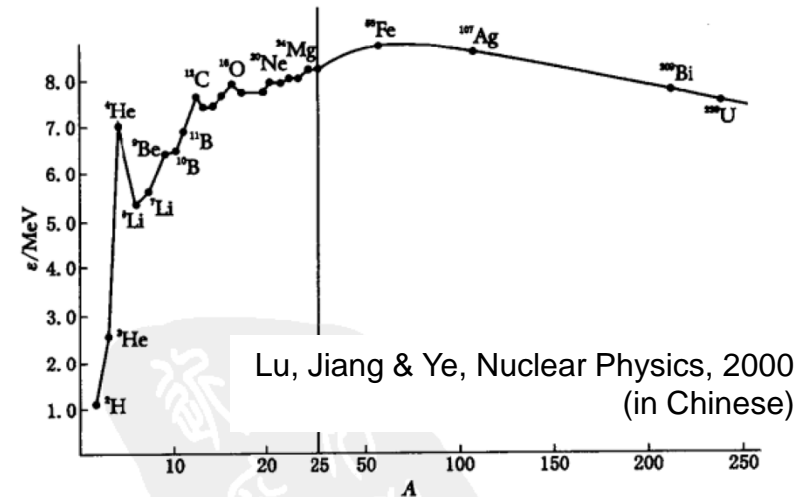
Self-consistent models



Liquid drop model: Nucleus ~ charged LD

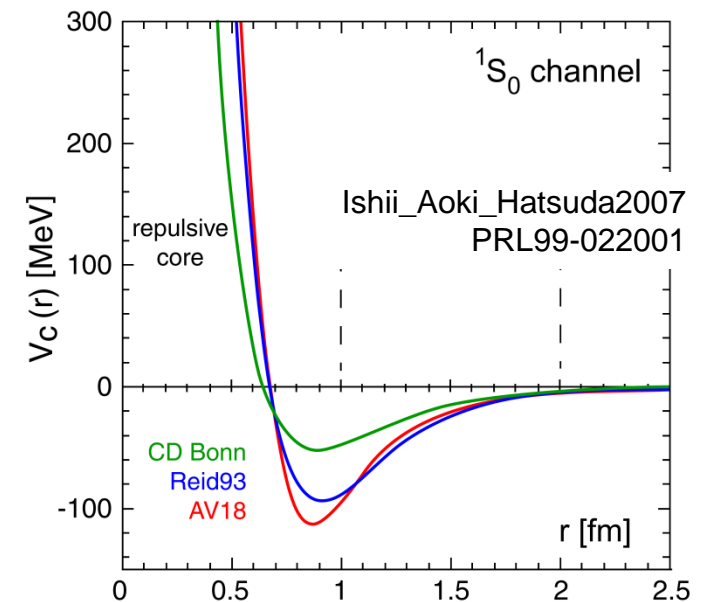
□ Similarities

- $B/A \sim 8 \text{ MeV} \Rightarrow$ Saturation of nuclear force
- $R \sim A^{1/3} \Rightarrow$ Incompressibility of nuclear matter



□ Differences

- Mean distance between nucleons larger than that corresponding to the minimum in NN potential
- Mean free path of nucleon comparable to nuclear size (Fermi gas model)



Liquid drop model: Nucleus ~ charged LD

$$B(Z, A) = a_v A$$

$$- a_s A^{2/3}$$

$$- a_c Z^2 A^{-1/3}$$

$$- a_a \left(\frac{A}{2} - Z \right)^2 A^{-1}$$

$$+ \delta a_p A^{-1/2}$$

□ Volume

□ Surface

□ Coulomb

□ Symmetry

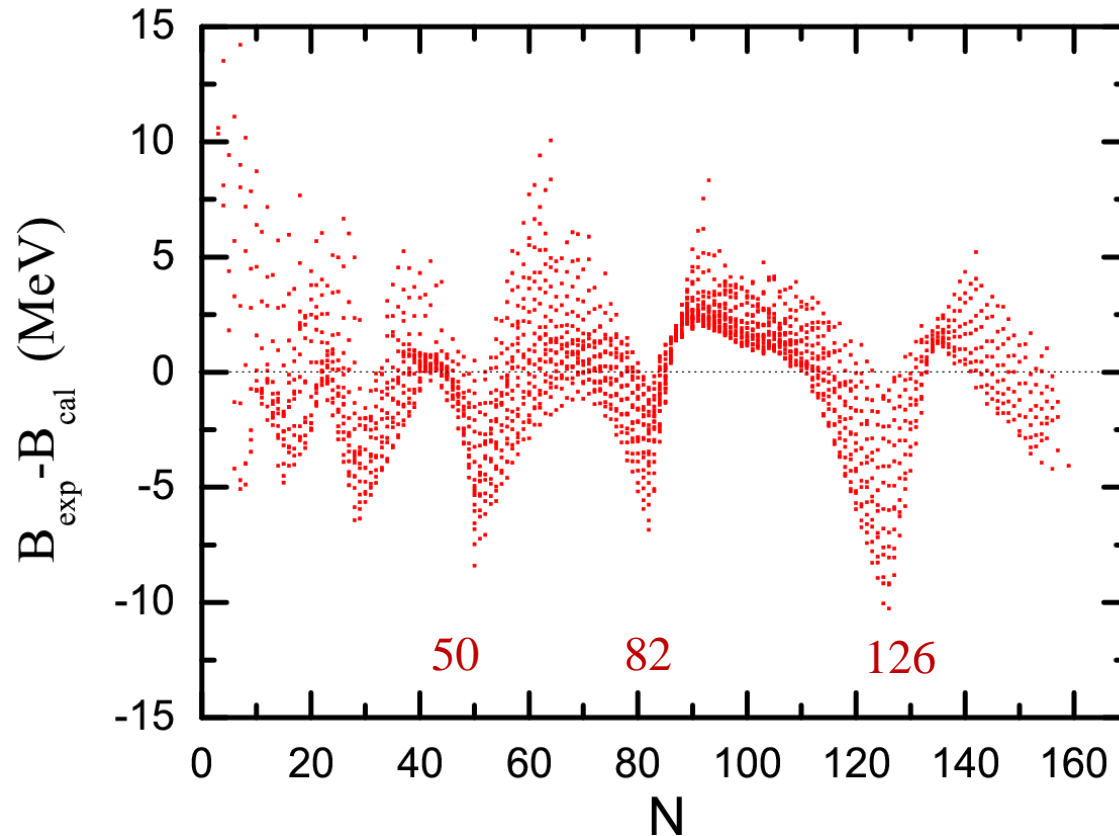
□ Pairing

Quantum effects

$$\delta = \begin{cases} 1 & \text{Even-even} \\ 0 & \text{Odd-A} \\ -1 & \text{Odd-odd} \end{cases}$$

Lu, Jiang & Ye, Nuclear Physics, 2000
(in Chinese)

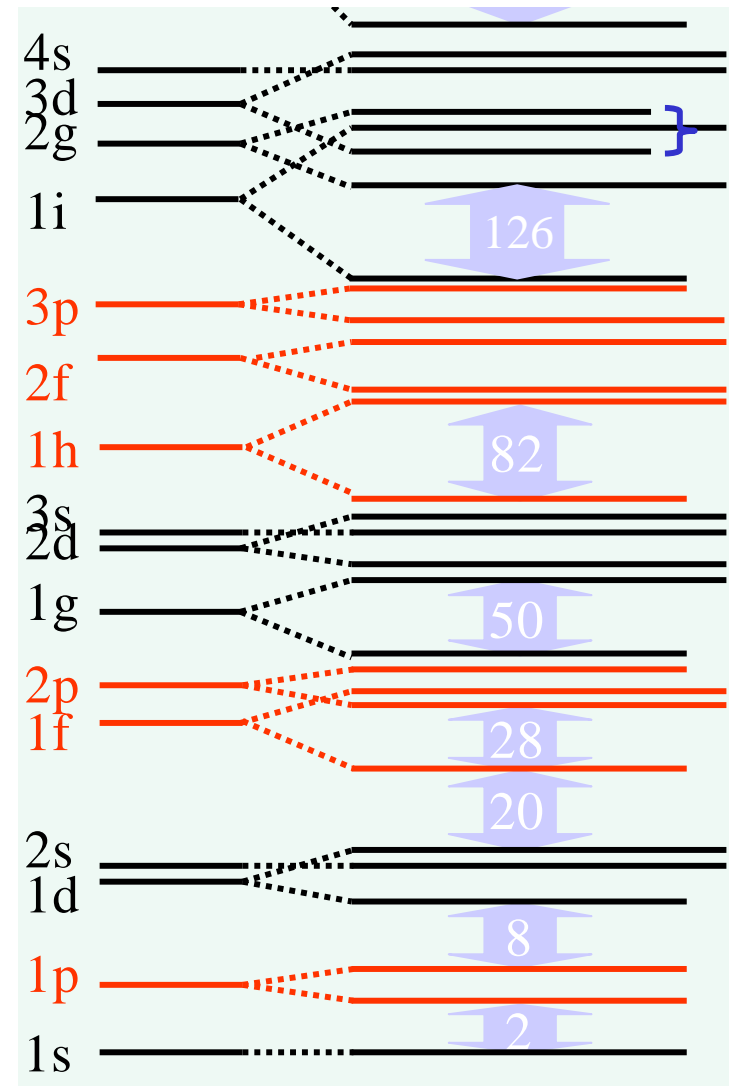
Comparison w/ experiment



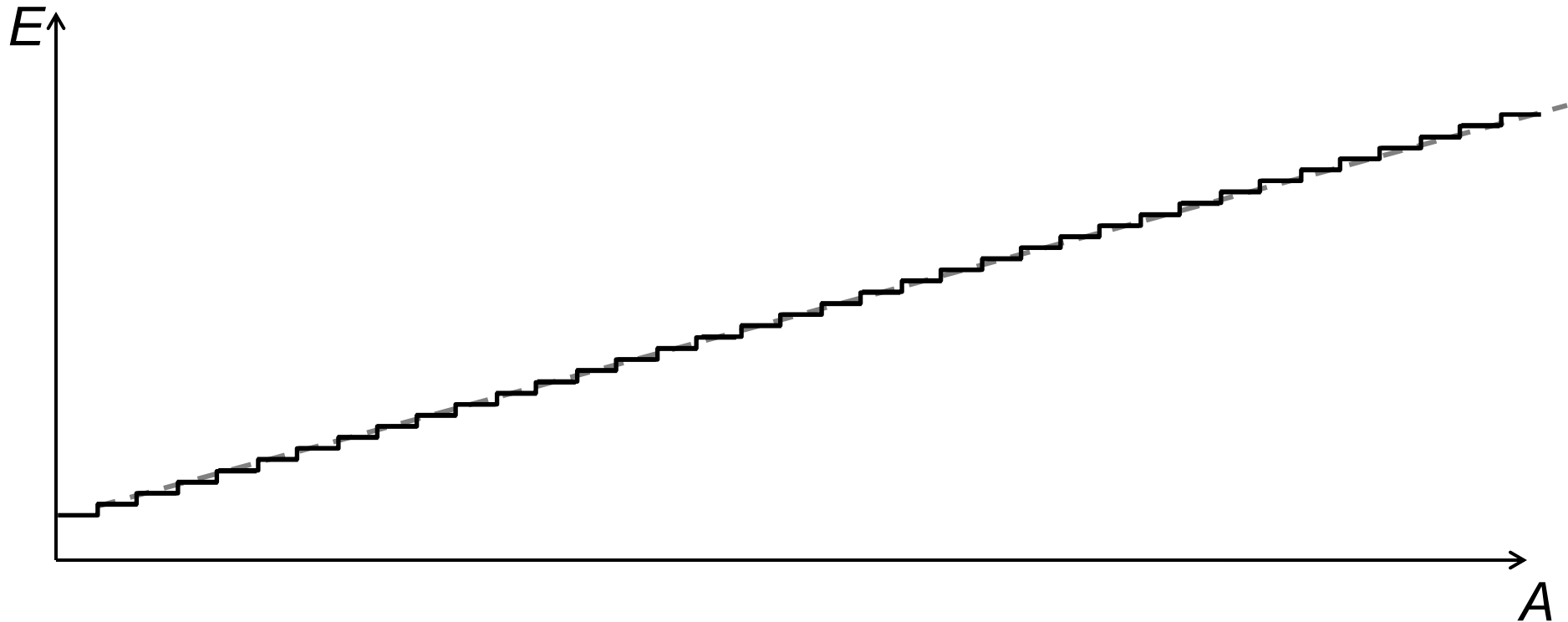
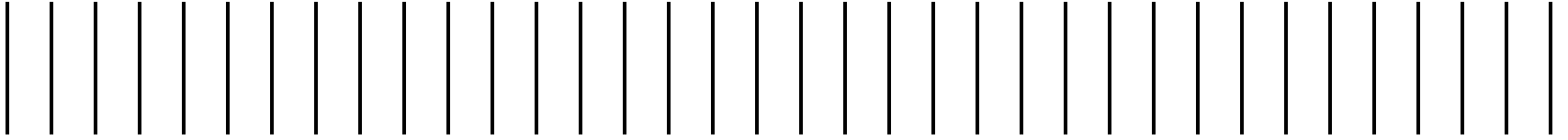
$$B(Z, N) = a_V A - a_S A^{2/3} - \frac{a_C Z^{5/3}}{1 + 0.452(T_z - T_z^*)/A} - a_a T_z (T_z + 1) A^{-0.86} + a_P A^{-1/2}$$

Single nucleon potential, spectra & magicities

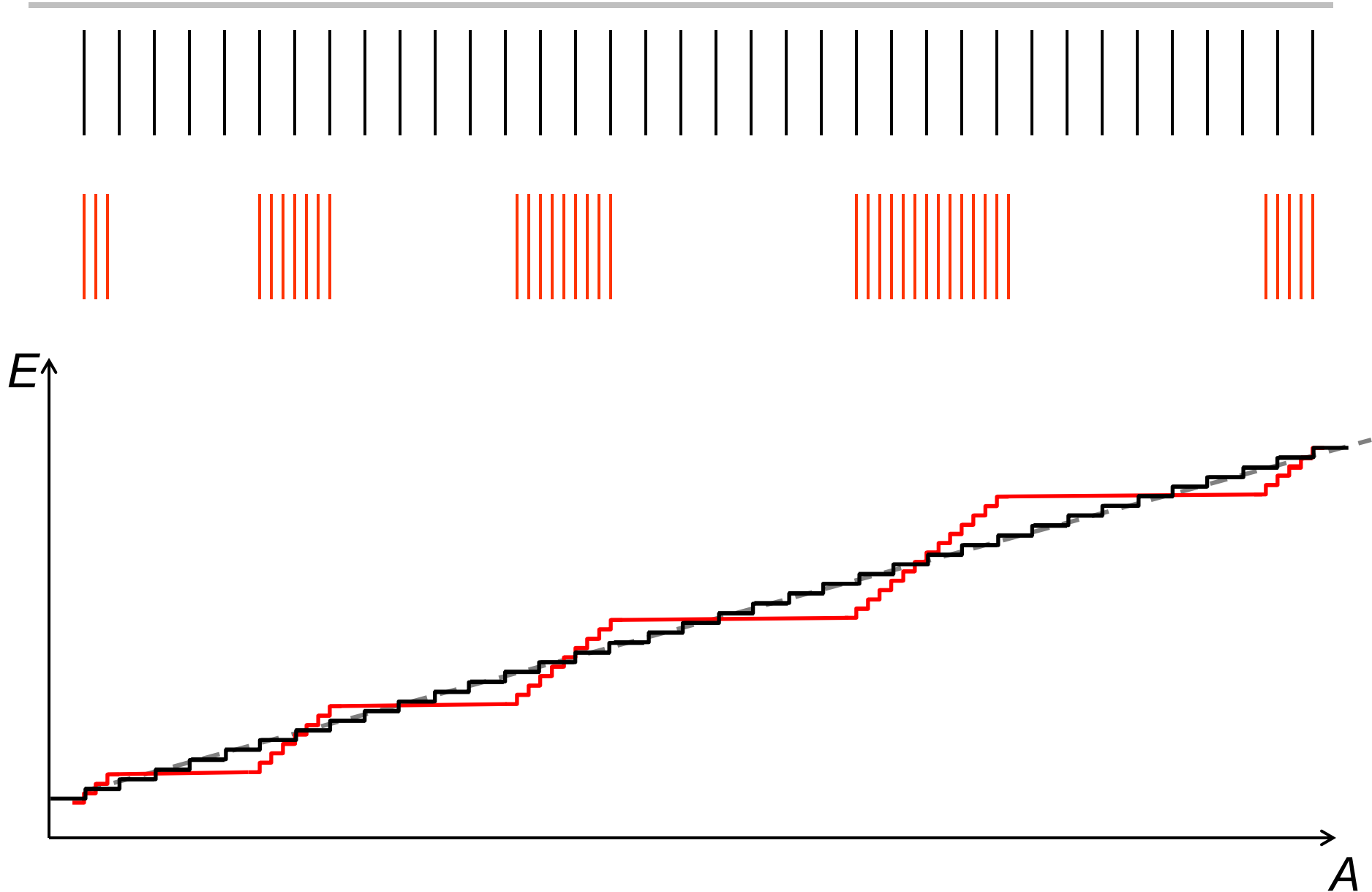
- Harmonic oscillator & square well
- Woods-Saxon
- Self-consistent (self-bound)



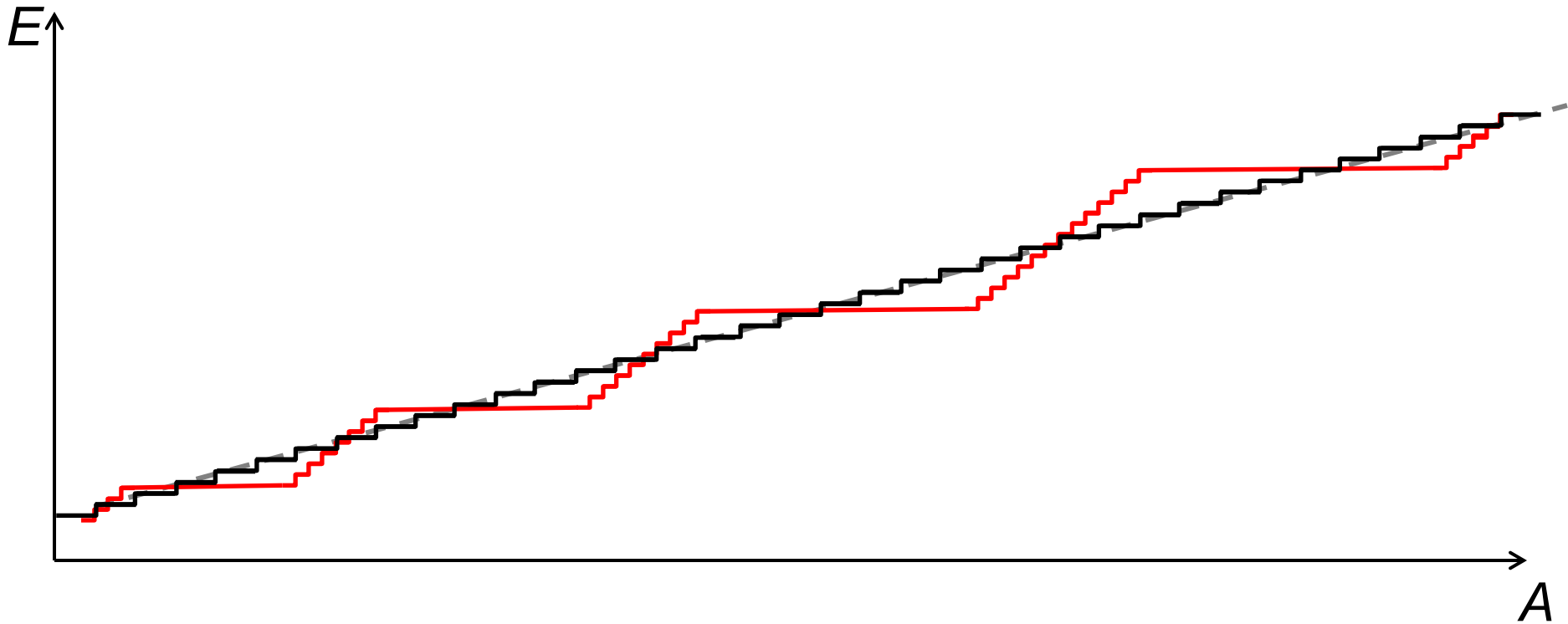
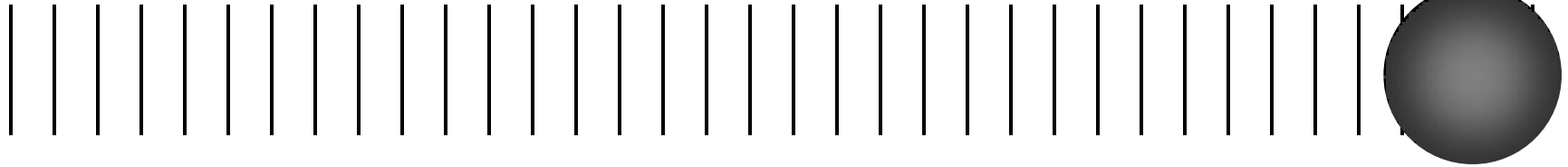
Quantum shell effects



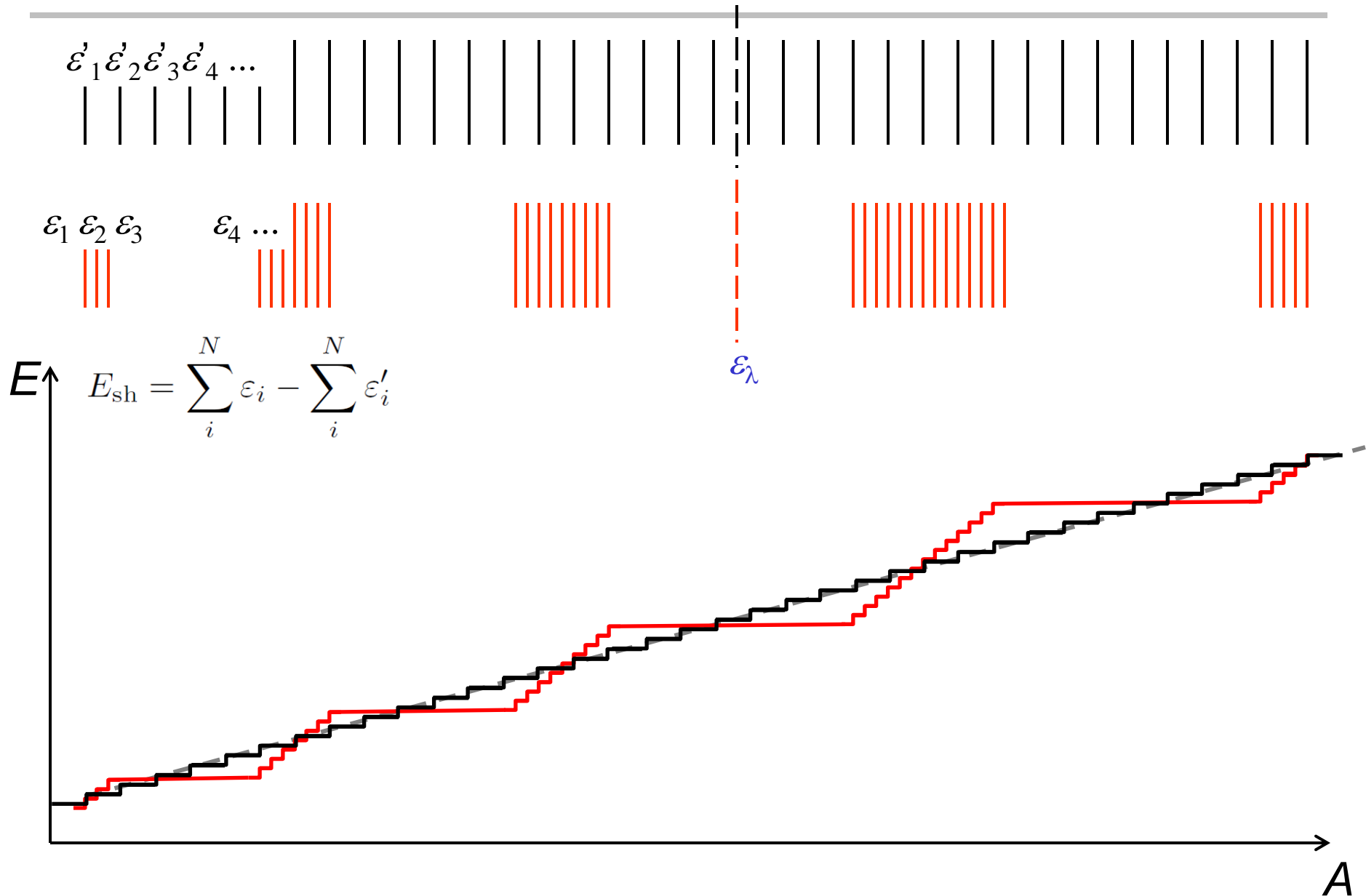
Quantum shell effects



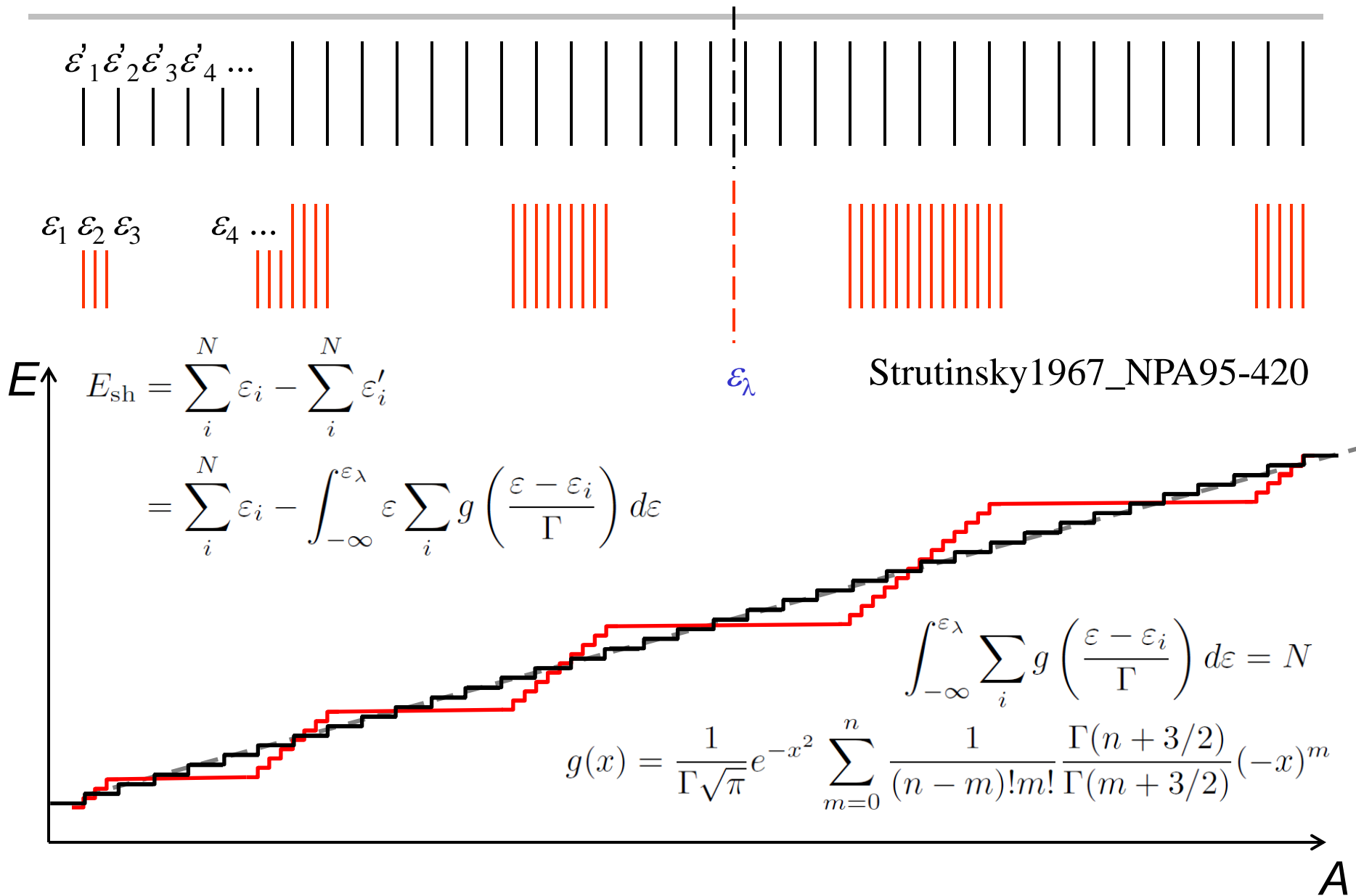
Quantum shell effects



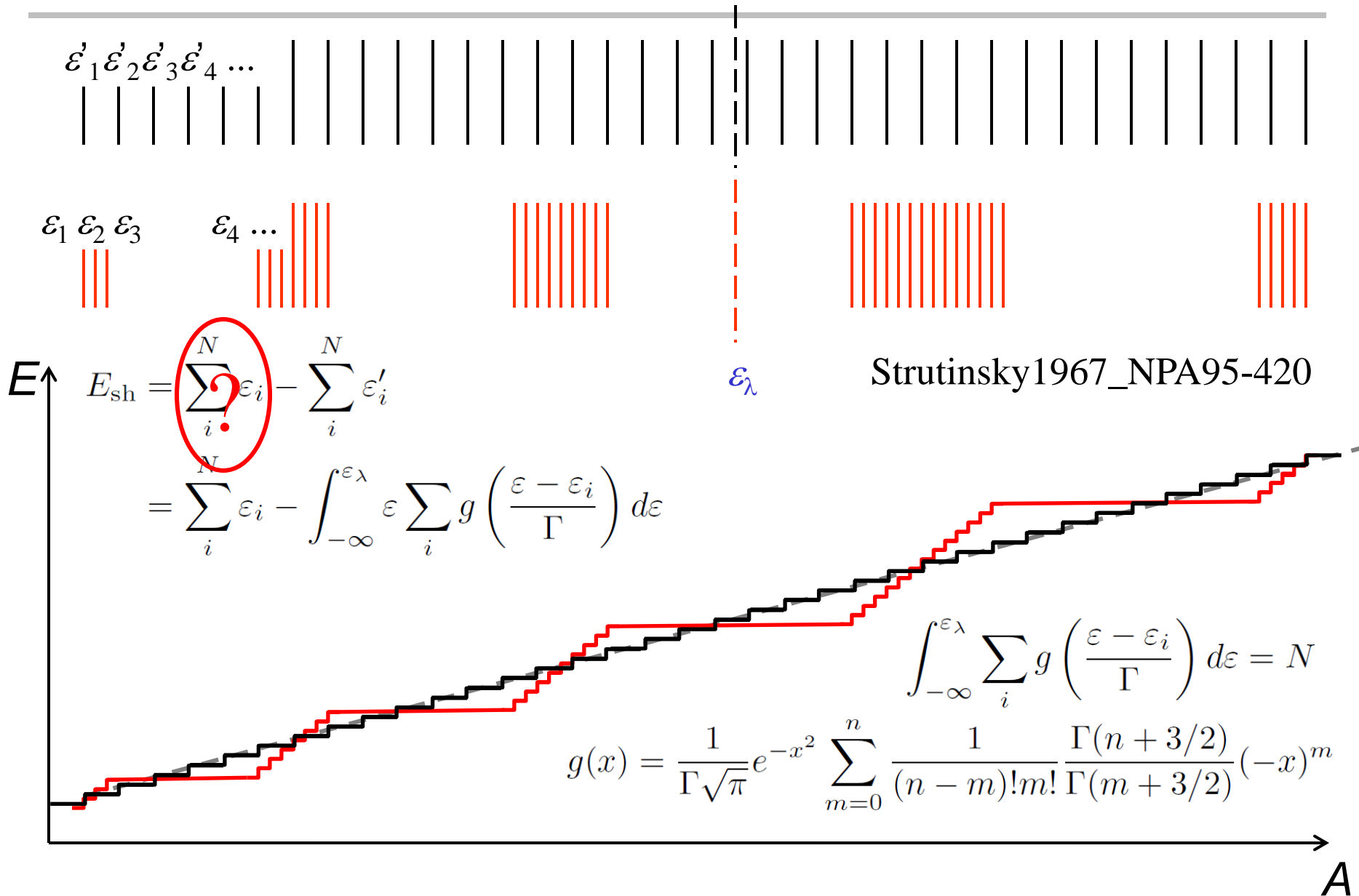
Microscopic shell corrections



Shell corrections: Strutinsky method



Shell corrections: Strutinsky method



Finite range droplet model (FRDM)

$$E_{\text{mac}}(Z, N, \text{shape}) =$$

$$M_{\text{H}}Z + M_{\text{n}}N$$

mass excesses of Z hydrogen atoms and N neutrons

$$+ \left(-a_1 + J\bar{\delta}^2 - \frac{1}{2}K\bar{\epsilon}^2 \right) A \quad \text{volume energy}$$

$$+ \left(a_2B_1 + \frac{9J^2}{4Q}\bar{\delta}^2\frac{B_s^2}{B_1} \right) A^{2/3} \quad \text{surface energy}$$

$$+ a_3A^{1/3}B_k \quad \text{curvature energy}$$

$$+ a_0A^0 \quad A^0 \text{ energy}$$

$$+ c_1\frac{Z^2}{A^{1/3}}B_3 \quad \text{Coulomb energy}$$

$$- c_2Z^2A^{1/3}B_r \quad \text{volume redistribution energy}$$

$$- c_4\frac{Z^{4/3}}{A^{1/3}} \quad \text{Coulomb exchange correction}$$

$$- c_5Z^2\frac{B_wB_s}{B_1} \quad \text{surface redistribution energy}$$

$$+ f_0\frac{Z^2}{A} \quad \text{proton form-factor correction to the Coulomb energy}$$

Moeller+2016_ADNDT109-110-1

$$- c_a(N - Z) \quad \text{charge-asymmetry energy}$$

$$+ W \left(|I| + \begin{cases} 1/A, & Z \text{ and } N \text{ odd and equal} \\ 0, & \text{otherwise} \end{cases} \right)$$

Wigner energy

$$+ \begin{cases} +\bar{\Delta}_p + \bar{\Delta}_n - \delta_{np}, & Z \text{ and } N \text{ odd} \\ +\bar{\Delta}_p, & Z \text{ odd and } N \text{ even} \\ +\bar{\Delta}_n, & Z \text{ even and } N \text{ odd} \\ +0, & Z \text{ and } N \text{ even} \end{cases}$$

average pairing energy

$$- a_{\text{el}}Z^{2.39} \quad \text{energy of bound electrons}$$

Weizsäcker-Skyrme (WS) mass model

$$E(A, Z, \beta) = E_{\text{LD}}(A, Z) \prod_{k \geq 2} (1 + b_k \beta_k^2) + \Delta E(A, Z, \beta) + \Delta_{\text{res}}$$

Liquid drop

Deformation

Shell

Residual

$$E_{\text{LD}}(A, Z) = a_v A + a_s A^{2/3} + E_C + a_{\text{sym}} I^2 A + a_{\text{pair}} A^{-1/3} \delta_{np}$$

$\Delta E(A, Z, \beta)$: Shell corrections

Δ_{res} (residual): Mirror, pairing, Wigner corrections, ...

Macro-micro concept & Skyrme energy density functional

Covariant Density Functional Theory (CDFT)

$$\begin{aligned} \mathcal{L} = & \bar{\psi}_i (i\partial - M) \psi_i + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - U(\sigma) - g_\sigma \bar{\psi}_i \sigma \psi_i \\ & - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - g_\omega \bar{\psi}_i \psi \psi_i \\ & - \frac{1}{4} \vec{R}_{\mu\nu} \vec{R}^{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \vec{\rho}^\mu - g_\rho \bar{\psi}_i \vec{\rho} \vec{\tau} \psi_i \\ & - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - e \bar{\psi}_i \frac{1 - \tau_3}{2} A \psi_i, \end{aligned}$$

Serot_Walecka1986_ANP16-1

Reinhard1989 RPP52-439

Ring1996_PPNP37-193

Vretenar_Afanasjev_Lalazissis_Ring2005_PR409-101

Meng_Toki_SGZ_Zhang_Long_Geng2006_PPNP57-470

$$(\alpha \cdot \mathbf{p} + \beta(M + S(\mathbf{r})) + V(\mathbf{r})) \psi_i = \epsilon_i \psi_i$$

Liang_Meng_SGZ2015_PR570-1

$$(-\nabla^2 + m_\sigma^2) \sigma = -g_\sigma \rho_S - g_2 \sigma^2 - g_3 \sigma^3$$

Meng_SGZ2015_JPG42-093101

$$(-\nabla^2 + m_\omega^2) \omega = g_\omega \rho_V - c_3 \omega^3$$

$$(-\nabla^2 + m_\rho^2) \rho = g_\rho \rho_3$$

$$-\nabla^2 A = e \rho_C$$

Various nuclear mass models: Comparison



^c Sobiczewski_Litvinov&Palczewski2018_ADNDT119-1
Atomic Data and Nuclear Data Tables

journal homepage: www.elsevier.com/locate/adt



Detailed illustration of the accuracy of currently used nuclear-mass models

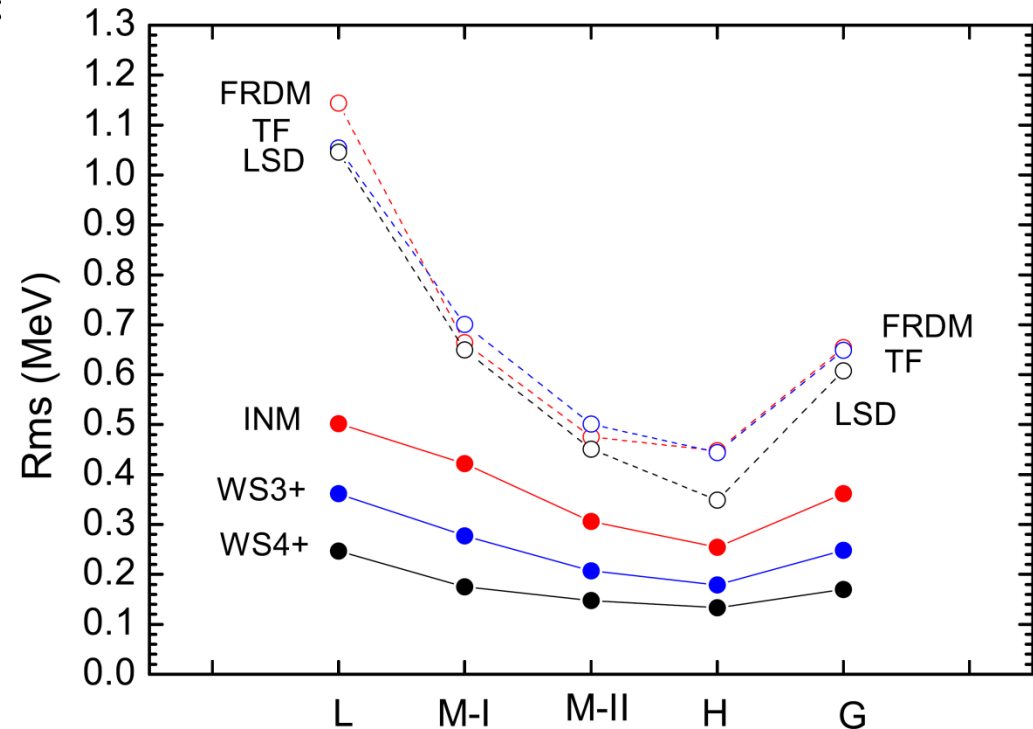


A. Sobiczewski ^{a,b,c,*}, Yu.A. Litvinov ^c, M. Palczewski ^c

^a National Centre for Nuclear Research, Hoza 69, 00-681 Warsaw, Poland

^b Joint Institute for Nuclear Research, 141980 Dubna, Russian Federation

^c GSI Helmholtzzentrum für Schwerionenforschung GmbH, 64291 Darmstadt, Germany



Relativistic continuum Hartree-Bogoliubov (RCHB) theory



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Xia+2018_ADNDT121-122-1

Atomic Data and Nuclear Data Tables

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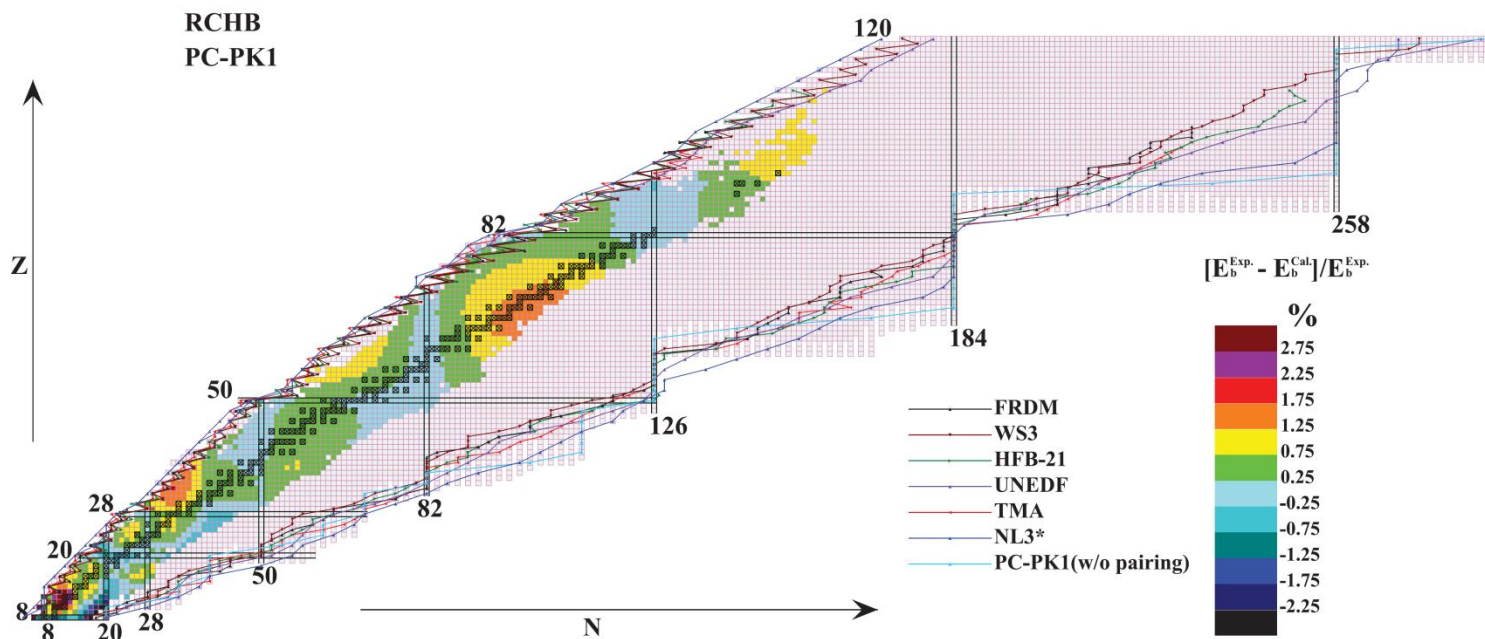


The limits of the nuclear landscape explored by the relativistic continuum Hartree–Bogoliubov theory



X.W. Xia^a, Y. Lim^{b,c}, P.W. Zhao^{d,e}, H.Z. Liang^f, X.Y. Qu^{a,g}, Y. Chen^{d,h}, H. Liu^d, L.F. Zhang^d, S.Q. Zhang^d, Y. Kim^c, J. Meng^{d,a,i,*}

- ^a School of Physics and Nuclear Engineering, Cyclotron Institute, Texas A&M University, TX 77703, USA
- ^b Cyclotron Institute, Texas A&M University, TX 77703, USA
- ^c Rare Isotope Science Project, Institute of Physics, Chinese Academy of Sciences, Beijing 100049, China
- ^d State Key Laboratory of Nuclear Physics and Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, China
- ^e Physics Division, Argonne National Laboratory, Lemont, IL 60469, USA
- ^f RIKEN Nishina Center, Wako 351-0199, Japan
- ^g School of Mechatronics Engineering, Beijing University of Aeronautics and Astronautics, Beijing 100023, China
- ^h Institute of Materials, China Academy of Engineering and Technology, Beijing 100049, China
- ⁱ Department of Physics, University of Jilin, Changchun 130012, China



RCHB + radial basis function approach

Chinese Physics C Vol. 43, No. 7 (2019) 074104

Shi_Niu_Liang2019_ChinPhysC43-074104

Mass predictions of the relativistic continuum Hartree-Bogoliubov model with radial basis function approach*

Min Shi(仕敏)¹ Zhong-Ming Niu(牛中明)^{2,3;1)} Hao-Zhao Liang(梁豪兆)^{4,5}

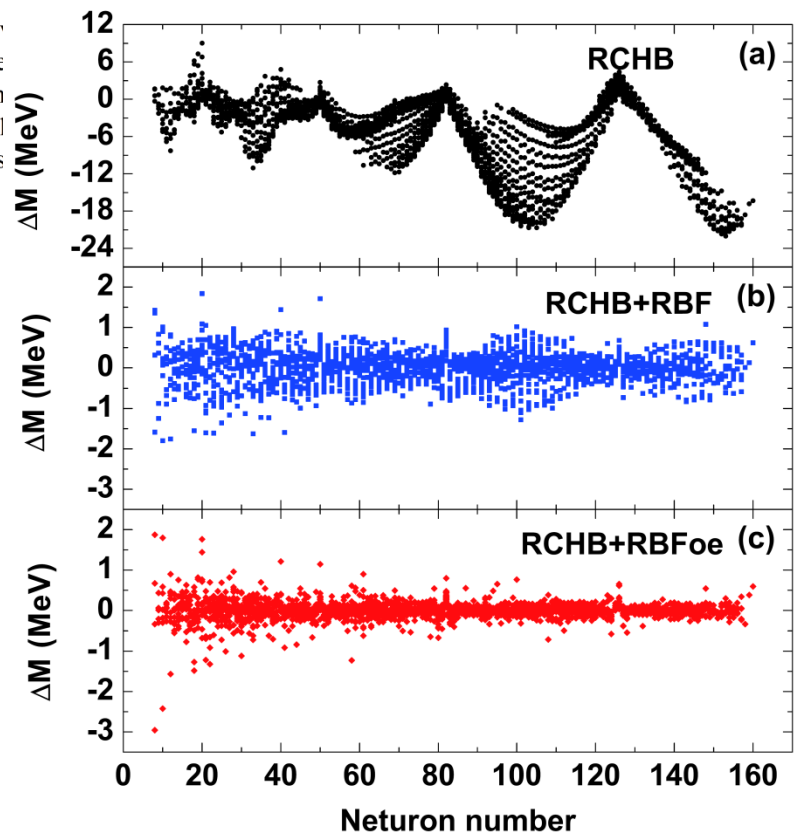
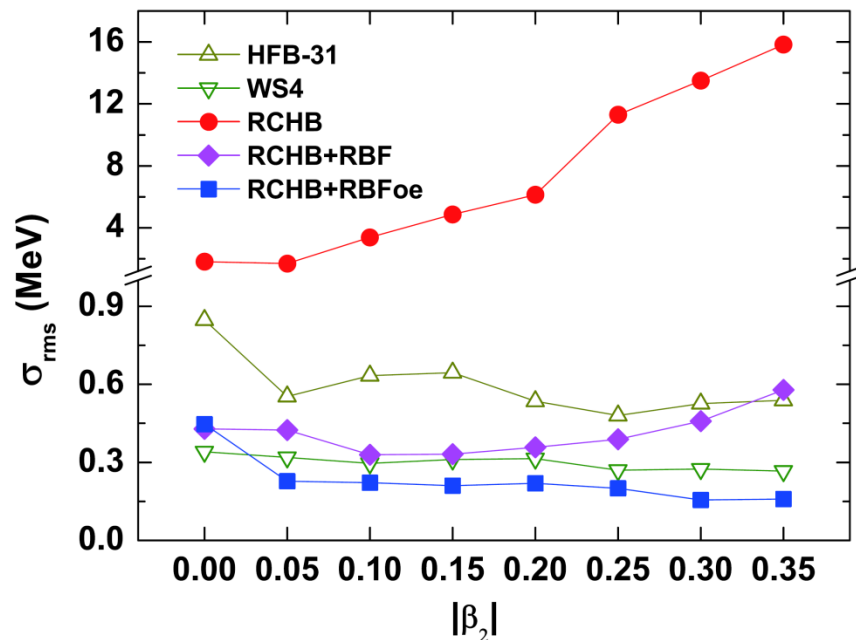
¹School of Mathematics and Physics, Anhui Jianzhu Uni

²School of Physics and Materials Science, Anhui Unive

³Institute of Physical Science and Information Technology, Anh

⁴RIKEN Nishina Center, Wako 351-01

⁵... The Univers



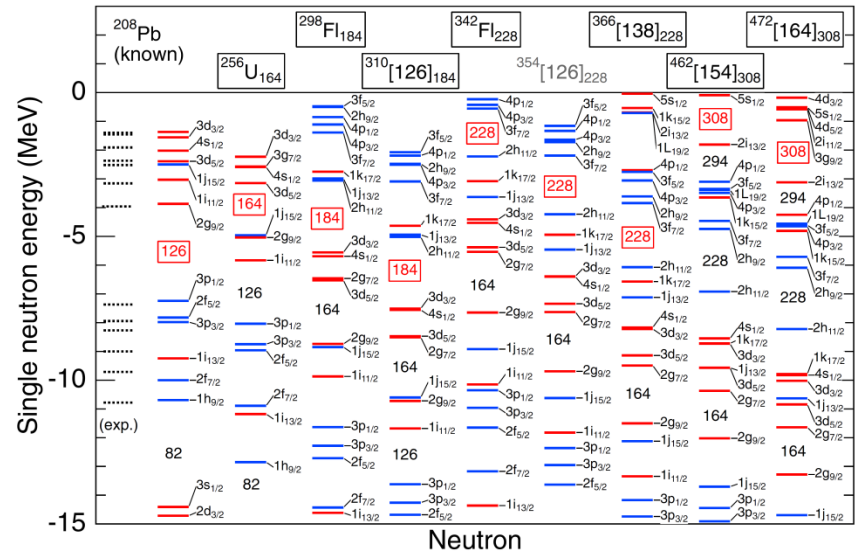
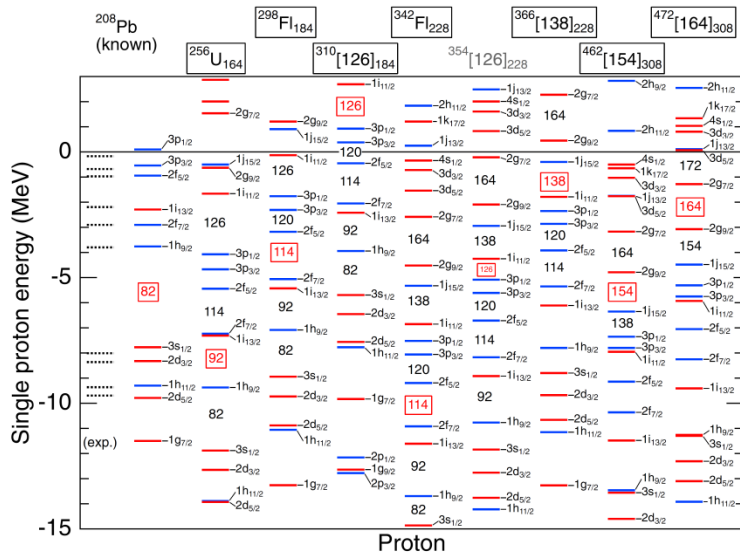
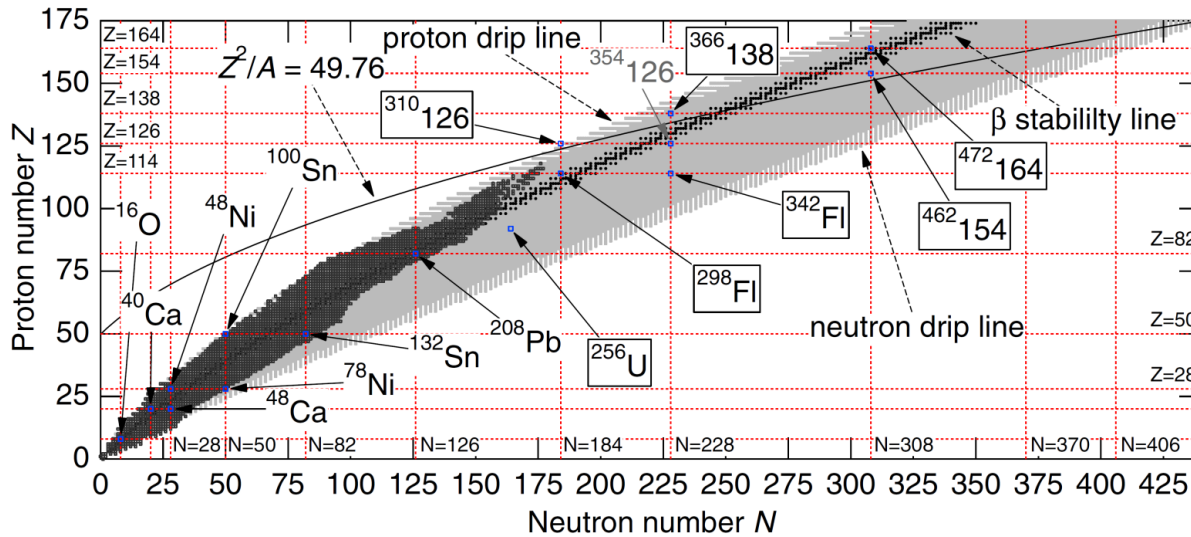
Lecture 2

- Challenges in synthesizing SHN

- Theoretical study of structure of SHN
 - Nuclear models
 - Next shell closures beyond ^{208}Pb as seen from single particle spectra, shell correction energy & nuclear shapes
 - Exotic shapes in SHN
 - Low-lying spectra of SHN & magicities

- Theoretical study of decay of SHN

Modified Woods-Saxon potential



Hyperheavy nuclei



ELSEVIER

Contents list Afanasjev_Agbemava_Ring2018_PLB782-533

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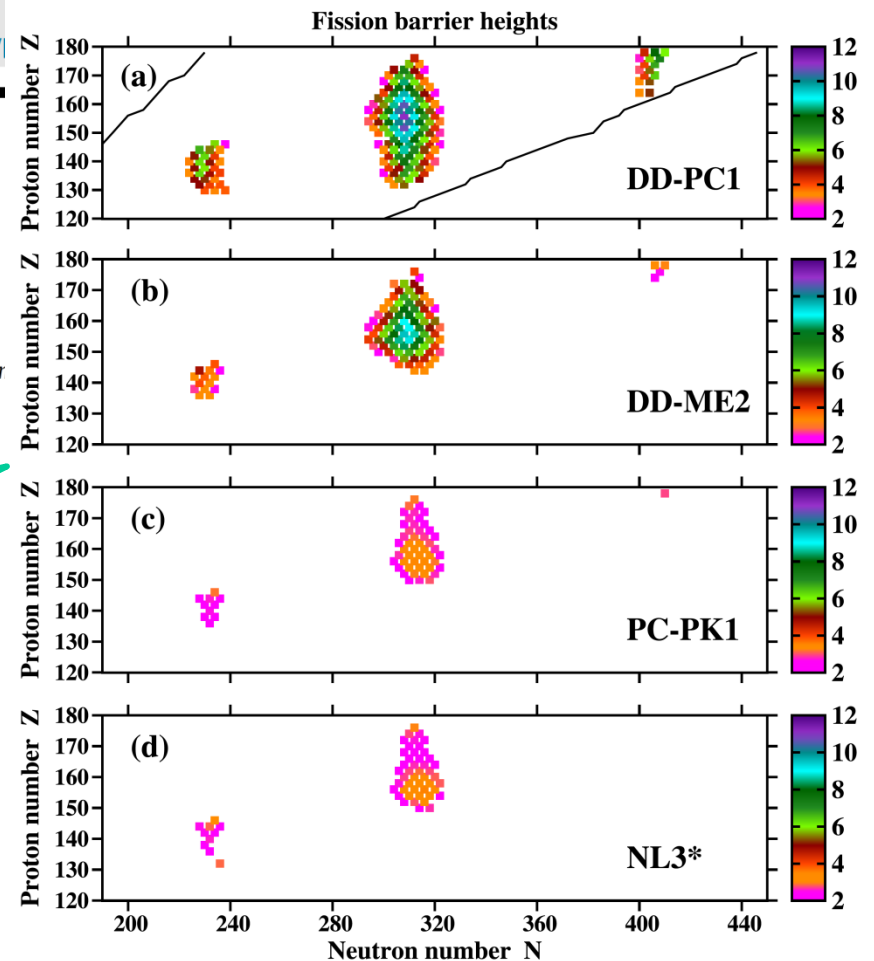
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Hyperheavy nuclei: Existence and stability

A.V. Afanasjev*, S.E. Agbemava, A. Gyawali

Department of Physics and Astronomy, Mississippi State University, MS 39762, United States of Ar

Z	N
138	230
156	310
174	410

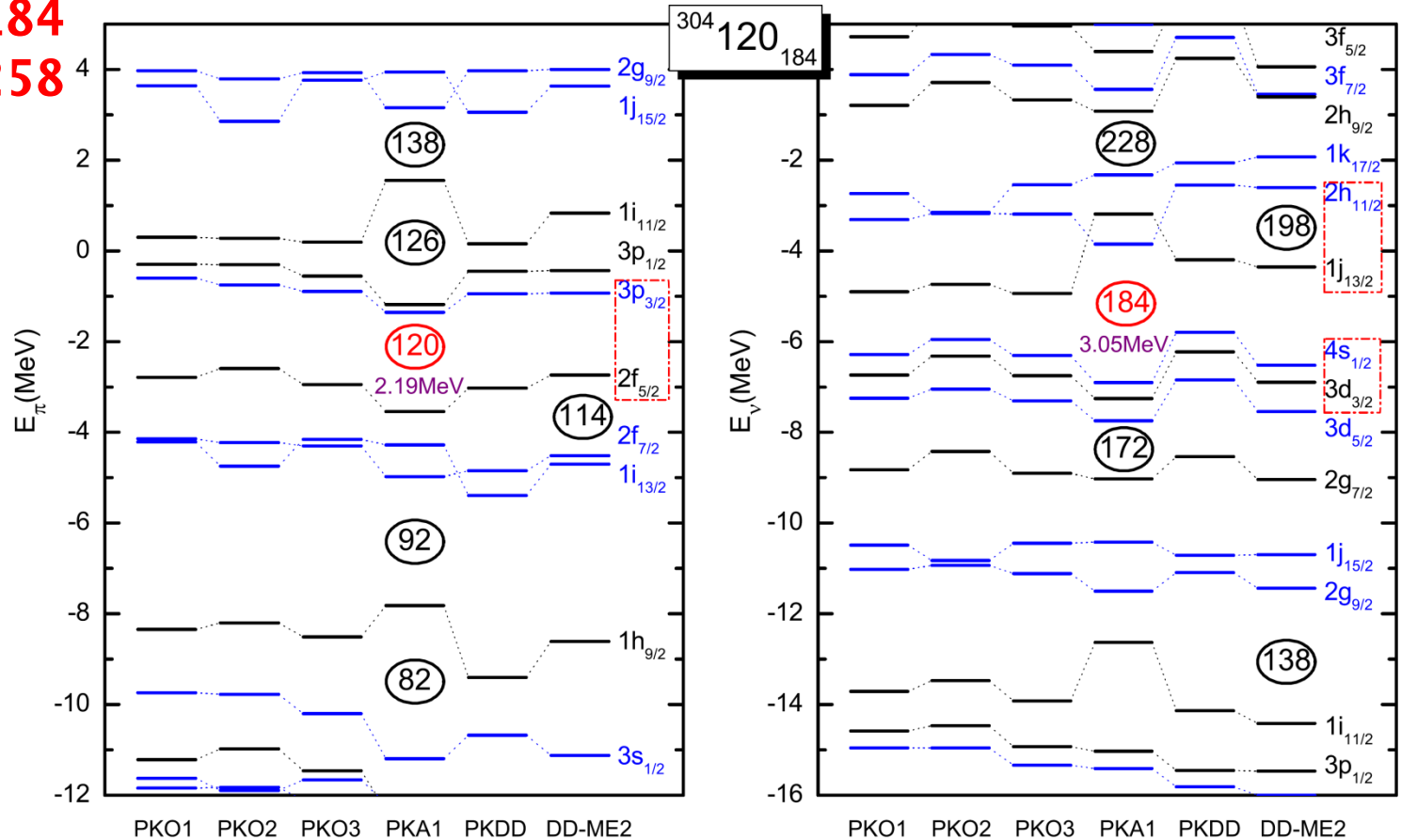


Relativistic Hartree-Fock-Bogoliubov model

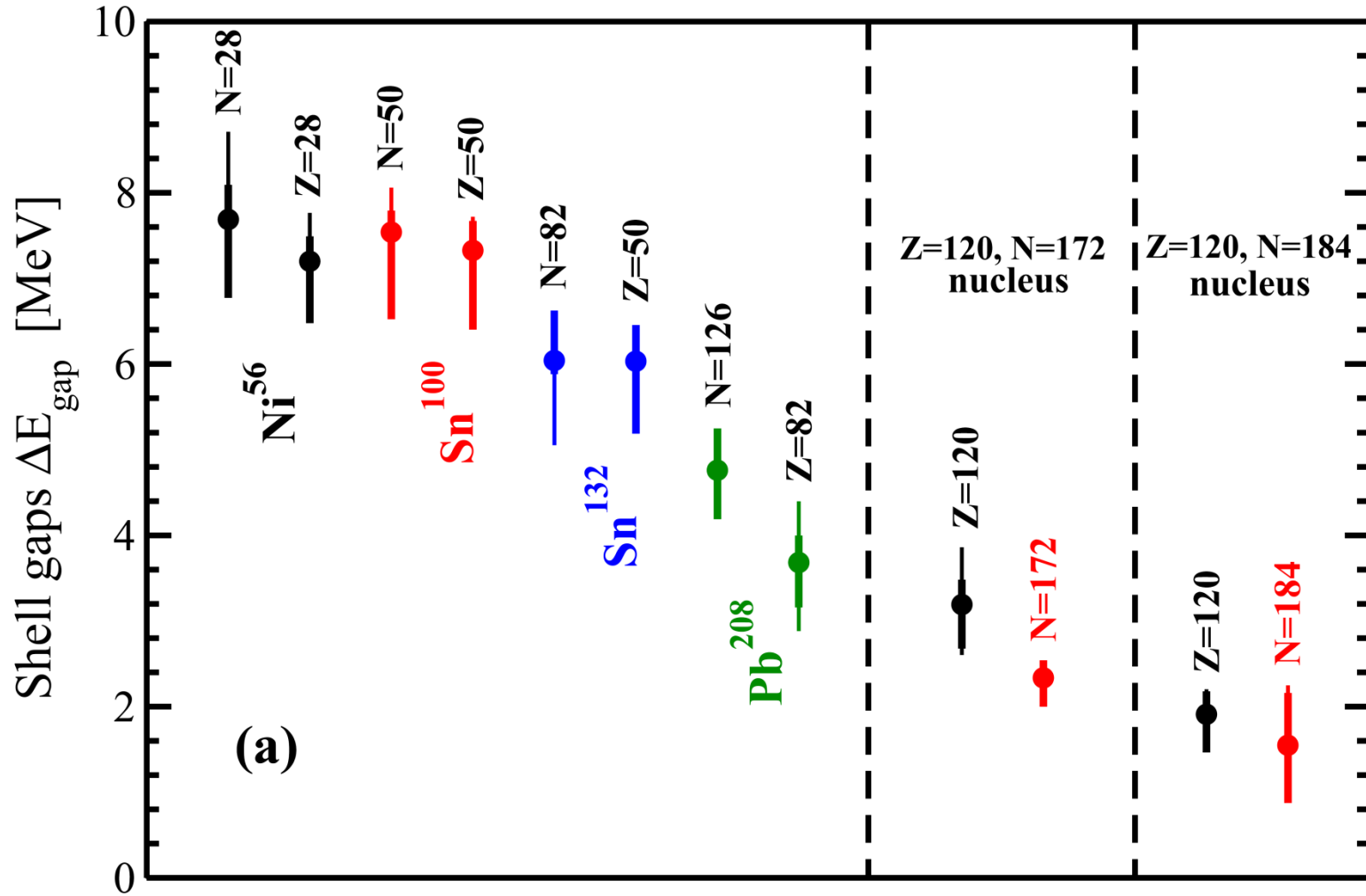
Z = 120 138

Li_Long_Margueron_Giai2014_PLB732-169

**N = 172 184
228 258**



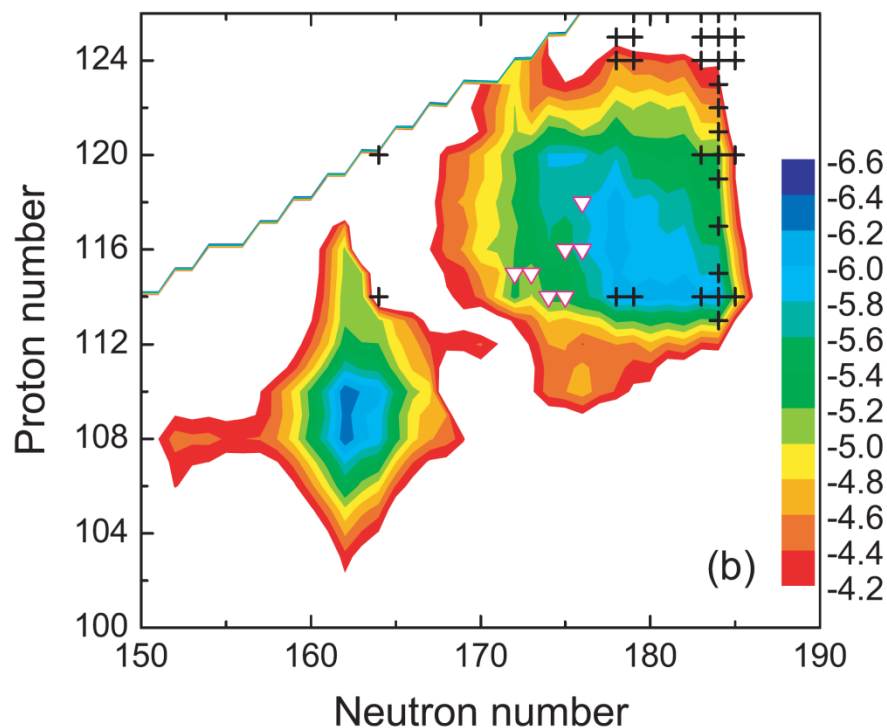
Single-particle level densities & gaps



Microscopic correction energies & deformations

Z = 114

N = 178

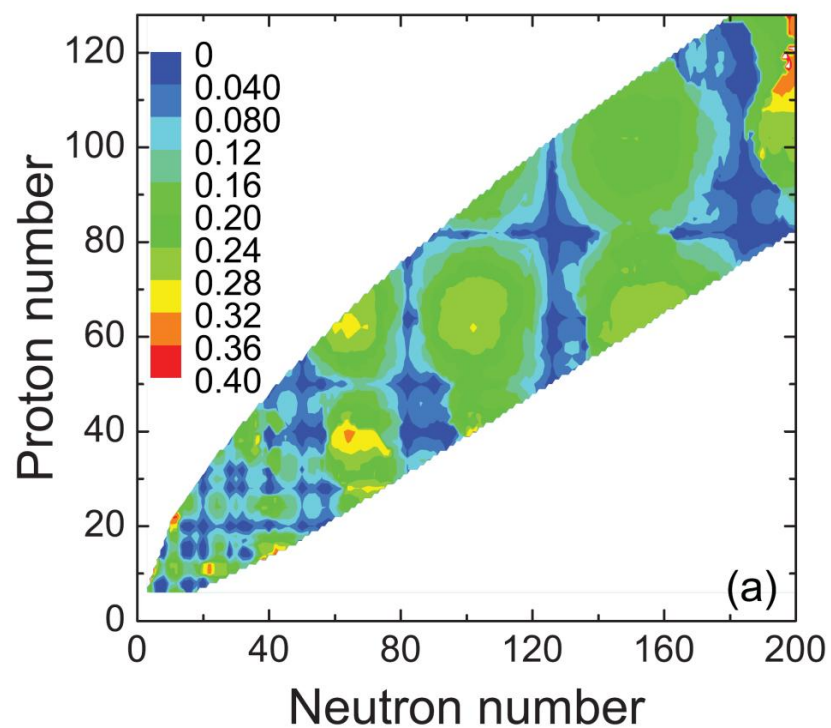


Shell correction energy

Wang_Liu_Wu2010_PRC81-044322

Z = 116~120

N = 176~178

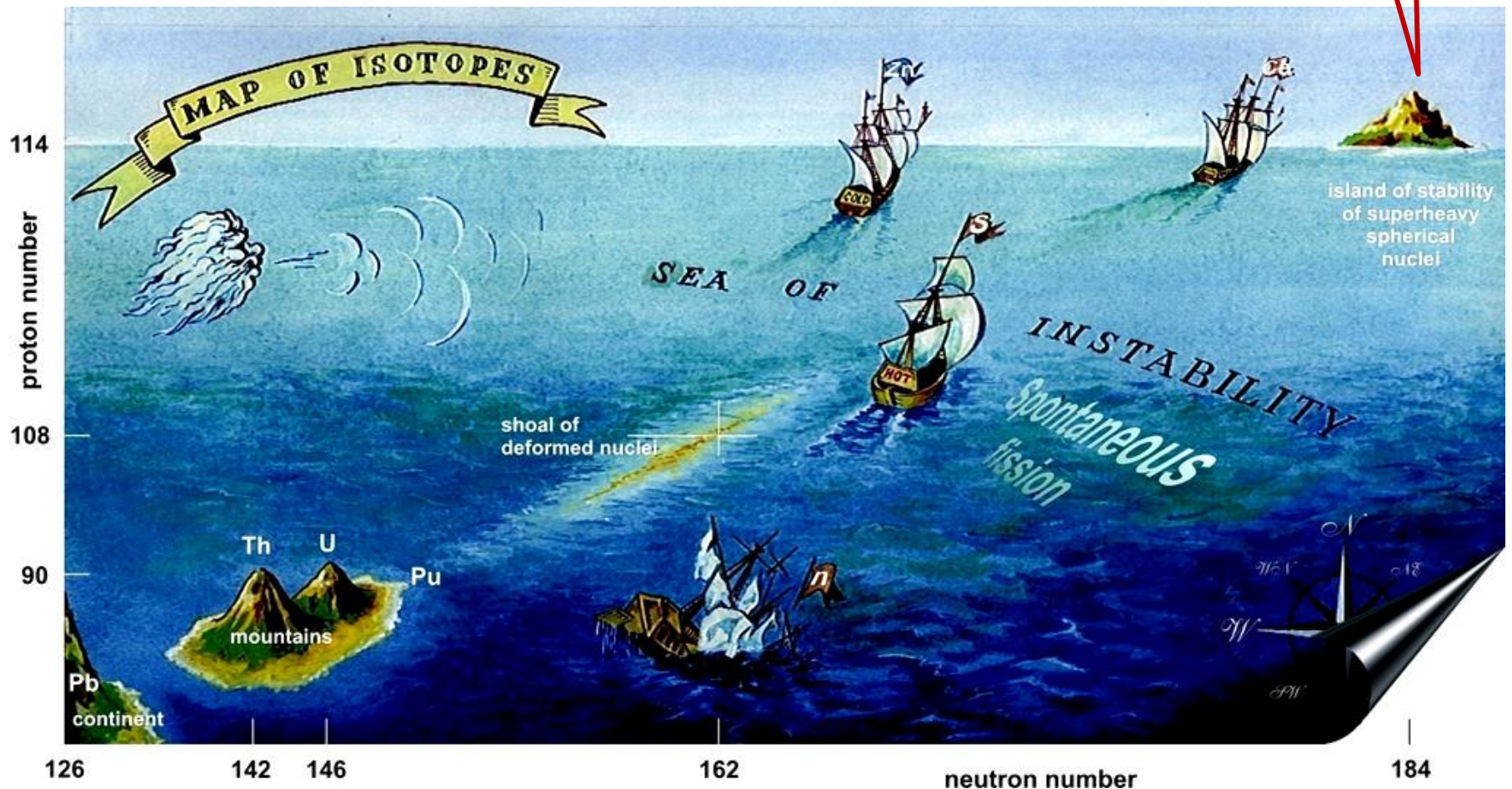
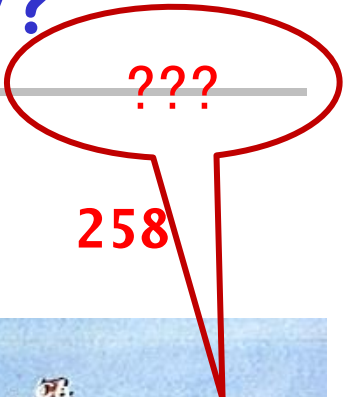


Quadrupole def. $|\beta_2|$

Wang_Liang_Liu_Wu2010PRC82-044304

Where is the island of stability?

Z = ?	114	116	120	126	132	138		
N = ?	172	176	178	184	198	228	238	258



Lecture 2

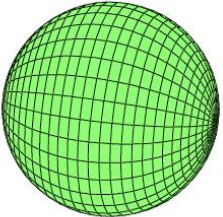
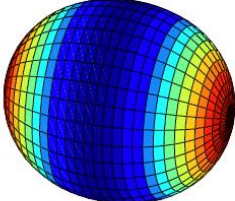
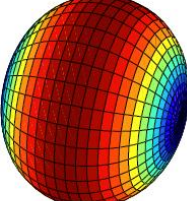
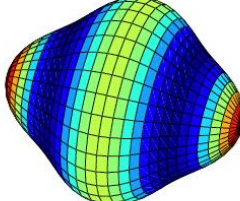
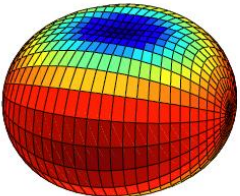
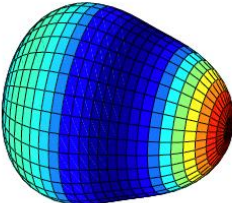
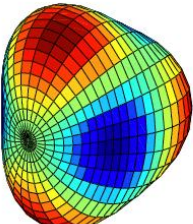
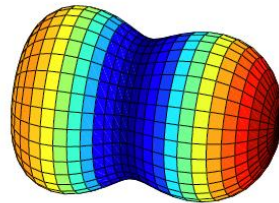
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Nuclear shapes

$$R(\theta, \varphi) = R_0 \left[1 + \beta_{00} + \sum_{\lambda=1}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \beta_{\lambda\mu}^* Y_{\lambda\mu}(\theta, \varphi) \right]$$

(a) $\beta_{\lambda\mu} = 0$	(b) $\beta_{20} > 0$	(c) $\beta_{20} < 0$	(d) $\beta_{40} > 0$
			
(e) $\beta_{22} \neq 0$	(f) $\beta_{30} \neq 0$	(g) $\beta_{32} \neq 0$	(h) $\beta_{20} \gg 0$
			

Shapes of SHN: Triaxial & octupole

Cwiok_Heenen&Nazarewicz2005_Nature433-705

Shape coexistence and triaxiality in the superheavy nuclei

S. Cwiok^{1*}, P.-H. Heenen² & W. Nazarewicz^{3,4,5}

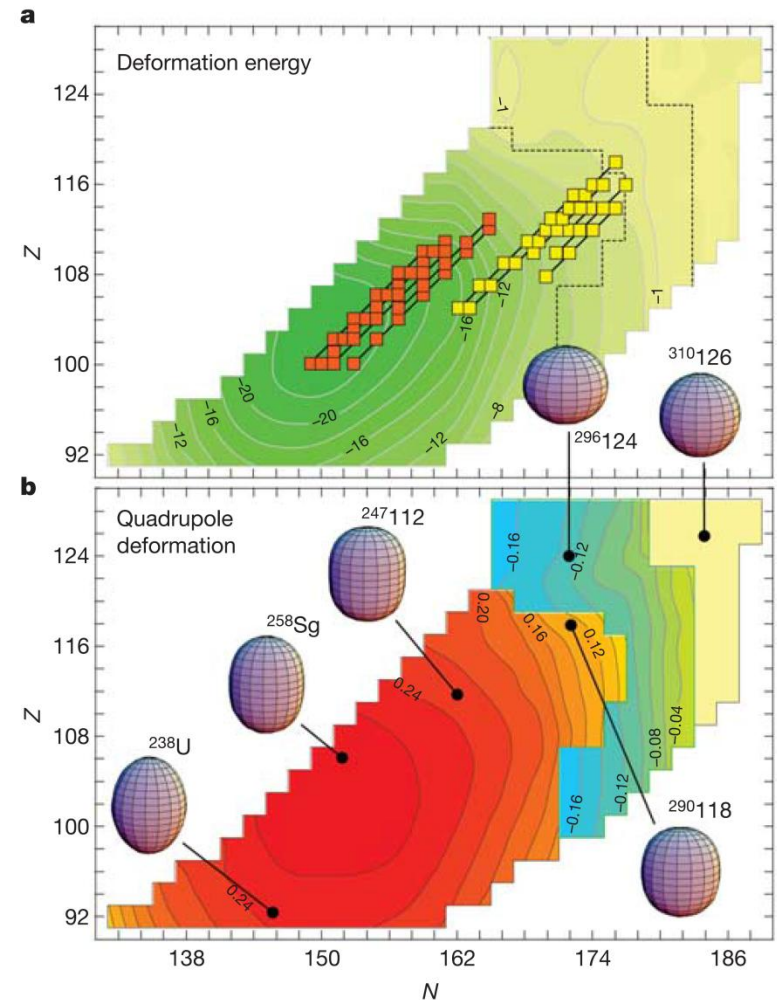
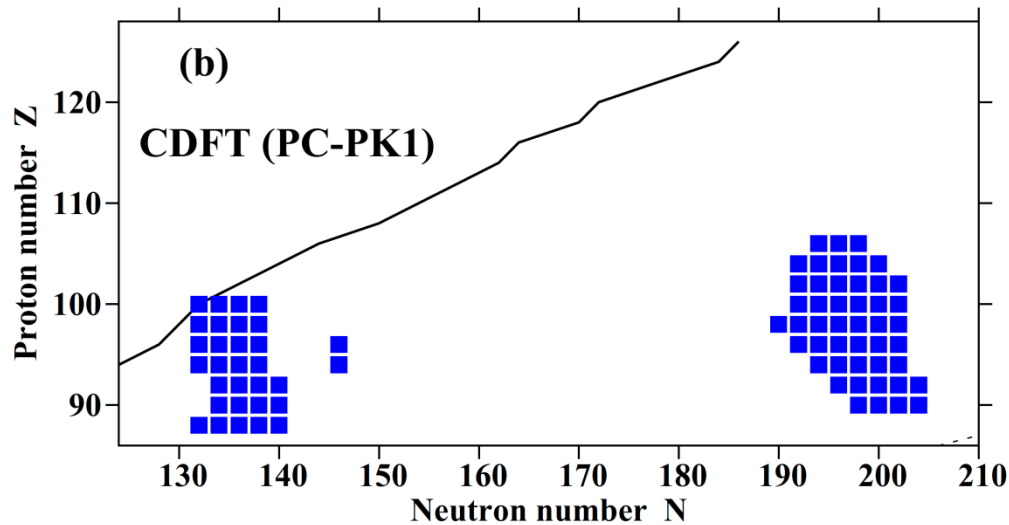
¹Institute of Physics, Warsaw University of Technology, ul. Koszykowa 75, PL-00662, Warsaw, Poland

²Service de Physique Nucléaire Théorique, Université Libre de Bruxelles, CP 229, B-1050 Brussels, Belgium

³Department of Physics and Astronomy, The University of Tennessee, Knoxville, Tennessee 37996, USA

⁴Physics Division, Oak Ridge National Laboratory, PO Box 2008, Oak Ridge, Tennessee 37831, USA

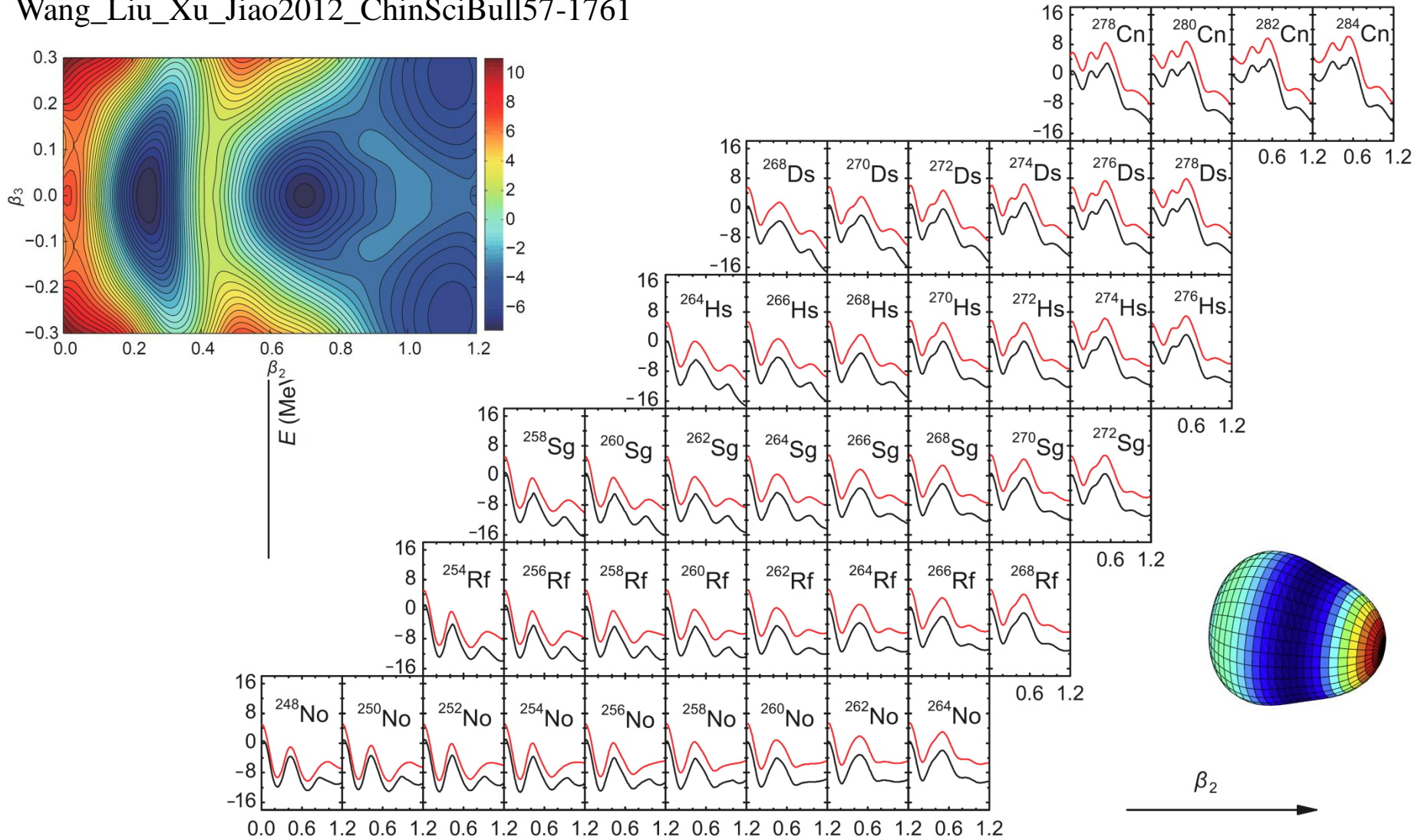
⁵Institute of Theoretical Physics, Warsaw University, ul. Hoza 69, PL-00681, Warsaw, Poland



Agbemava&Afanasjev2017_PRC96-024301

Octupole correlations in SHN

Wang_Liu_Xu_Jiao2012_ChinSciBull57-1761



Possible exotic shapes in SHN

□ Bubble or toroidal

1960s, H. A. Bethe, J. A. Wheeler, ...

Wong1973_APNY77-279

Dietrich_Pomorski1998_PRL80-37

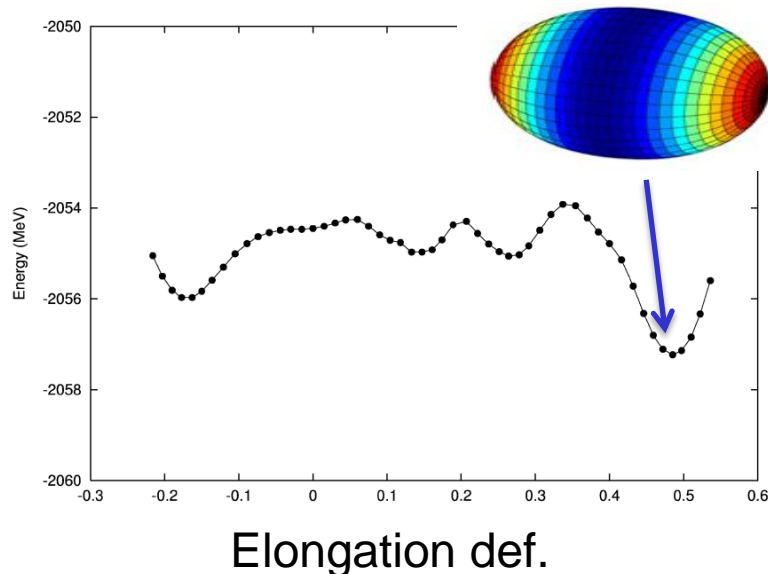
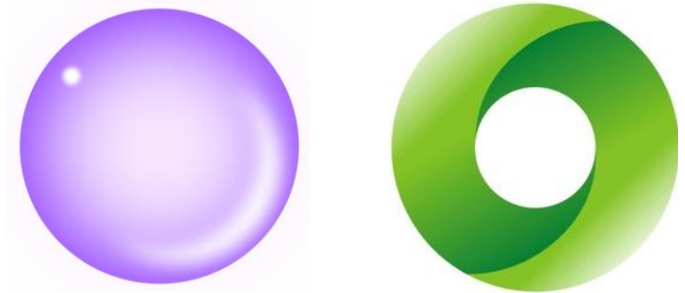
Pei_Xu_Stevenson 2005_PRC71-034302

...

□ Superdeformed or tetrahedral

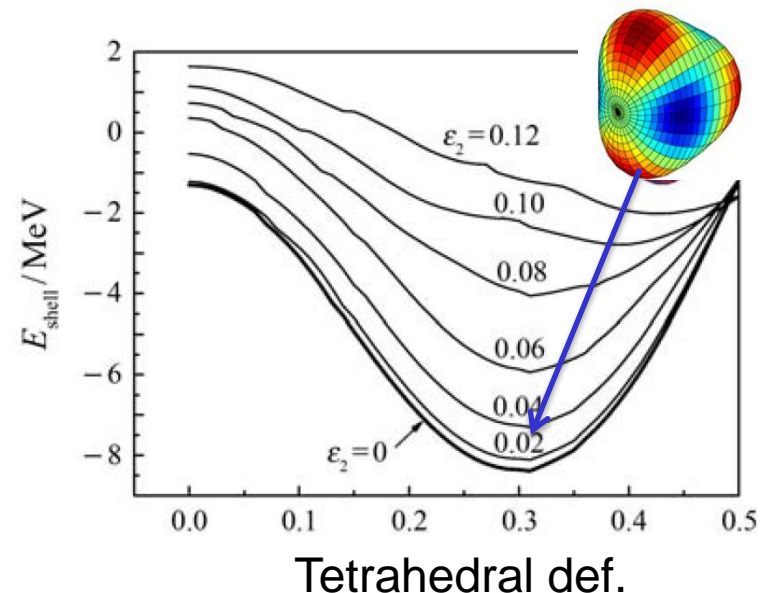
Ren_Toki2001_NPA689-691

Ren2002_PRC65-051304R



Chen_Gao2010_NPA834-380c

2013_NPR30-278



Lecture 2

- Challenges in synthesizing SHN

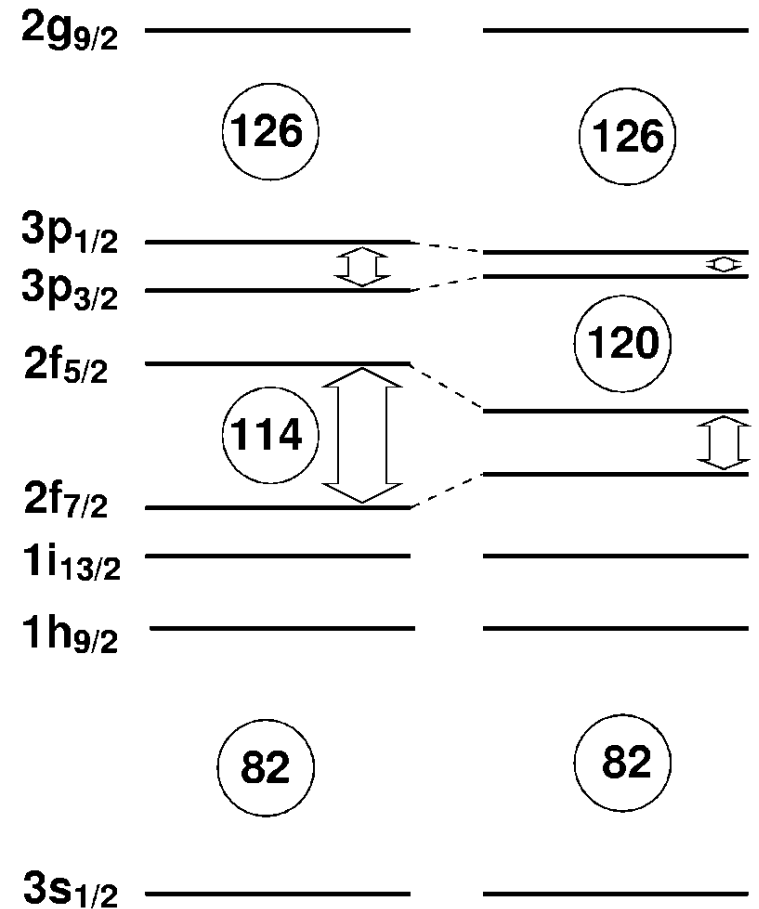
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 - **Low-lying spectra of SHN & magicities**

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Spectroscopy of nuclei with $Z \sim 100$

Synthesis of SHN \Rightarrow Decay modes & energies; X-sections, ...

- Spectroscopy of SHN
 - Detailed structure & stability
- Spectroscopy of deformed nuclei with $Z \sim 100$ & $N \sim 152$
 - Of interest in itself --- occurrence of deformation & K -isomerism
 - Orbitals around the Fermi level in these nuclei stem from those connected to the spherical shell gaps in SHN (1/2-[521])



Experimental facilities & status

RITU@JYFL

SHIP@GSI

VASSILISSA@FLNR

FMA@ANL

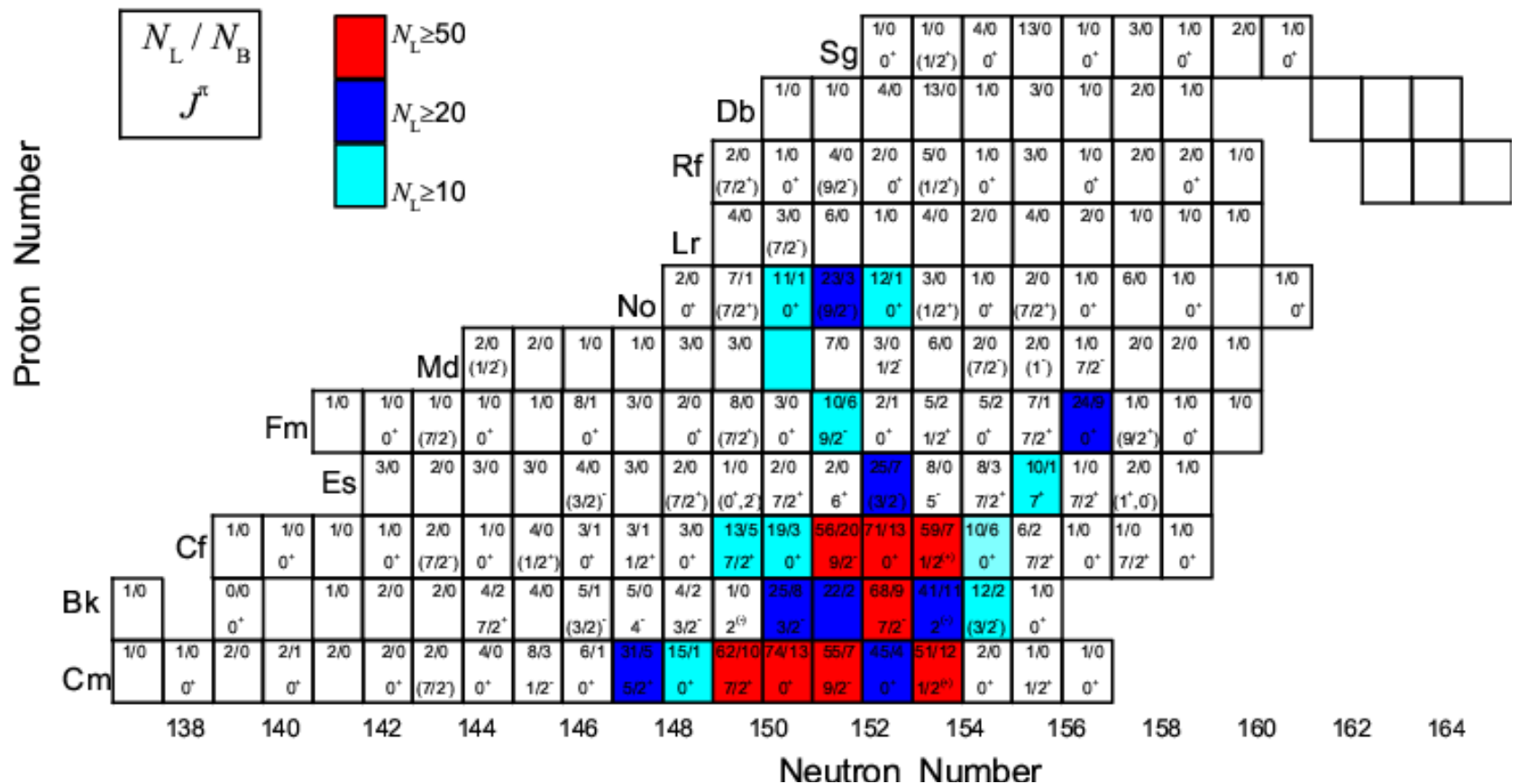
LISE3@GANIL

RMS@JAEA

HIRFL@IMP

...

Data from ENSDF (Apr., 2012) by Zhen-Hua Zhang (张振华)



Theoretical study of low-lying spectra

□ Self-consistent approaches

Egido_Robledo2000_PRL85-1198 Afanasjev...2003_PRC67-024309
Delaroche...2006_NPA771-103 Bender...2003_NPA723-354
[Adamian...2011_PRC84-024324](#)

□ Macroscopic-Microscopic models

Cwiok...1994_NPA573-356 Muntian...1999_PRC60-041302R
Sobiczewski...2001_PRC63-034306
Parkhomenko Sobiczewski2004 APPB35-2447
Parkhomenko_Sobiczewski2005_APPB36-3115
[Adamian...2011_PRC 84-024324](#)

□ Projected shell model

Sun...2008_PRC77-044307 Chen...2008_PRC77-061305
Al-Khudair...2009_PRC79-034320

□ Cranking shell model

He...2009_NPA817-45 [Zhang...2011_PRC83-011304R](#)
Liu...2012_PRC86-011301R [Zhang...2012_PRC85_014324](#)
 [Zhang...2013_PRC87-054308](#)

Cranked Nilsson model w/ pairing treated by a particle number conserving method

$$H_{\text{CSM}} = H_0 + H_P = H_{\text{Nil}} - \omega J_x + H_P$$

$$H_{\text{Nil}} = \frac{1}{2} \hbar \omega_0 (\varepsilon_2, \varepsilon_4) \left[-\nabla_\rho^2 + \frac{1}{3} \left(2 \frac{\partial^2}{\partial \zeta^2} - \frac{\partial^2}{\partial \xi^2} - \frac{\partial^2}{\partial \eta^2} \right) + \rho^2 \right. \\ \left. - \frac{2}{3} \varepsilon_2 \rho^2 P_2(\cos \theta_t) + 2 \varepsilon_4 \rho^2 P_4(\cos \theta_t) \right] \\ - 2 \kappa \hbar \omega_0 \left(\vec{s} \cdot \vec{l}_t - \mu (\rho^4 - \langle \rho^4 \rangle_N) \right),$$

Zhang...2011_PRC83-011304R

Zhang...2012_PRC85_014324

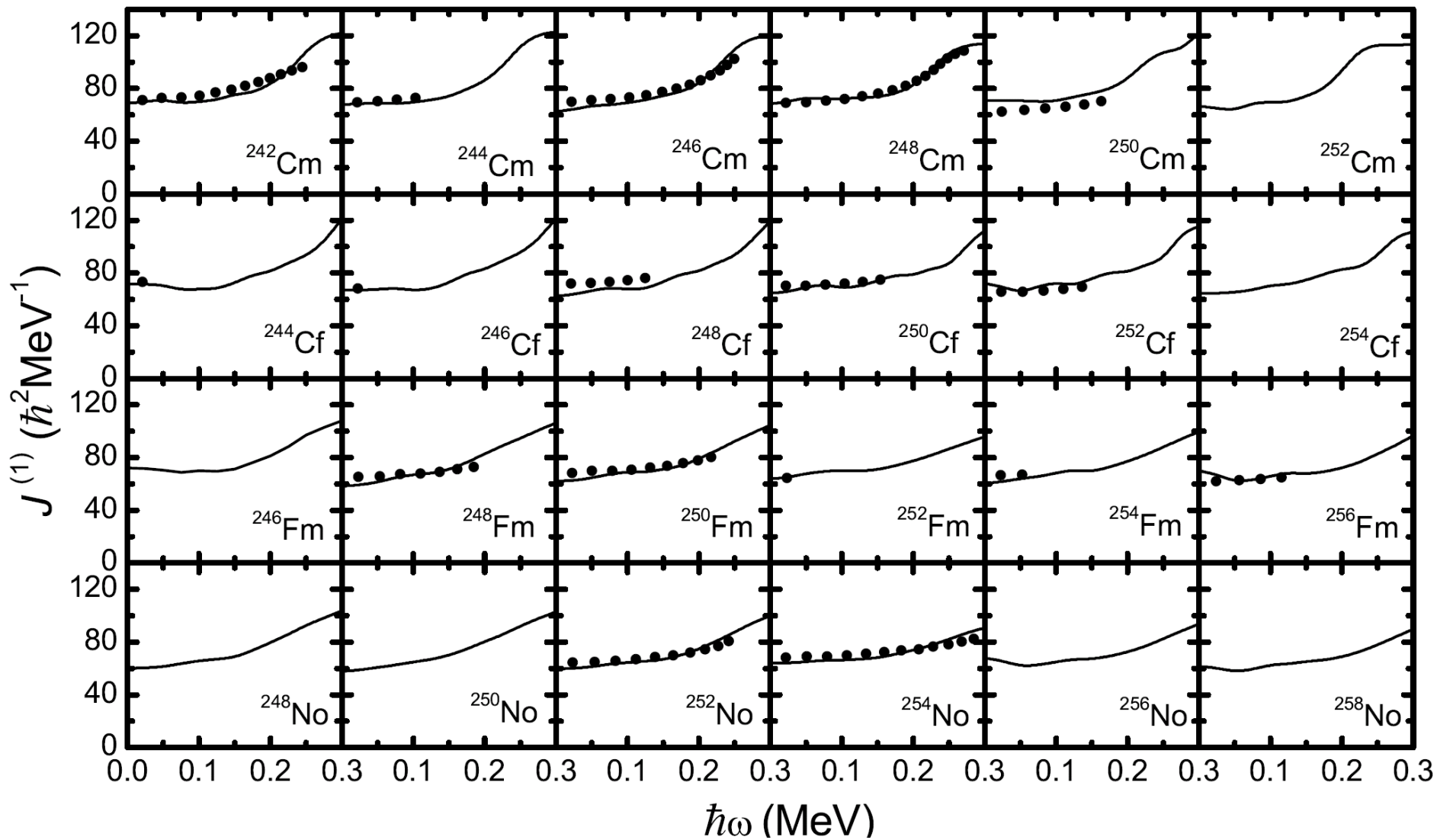
Z. H. Zhang (张振华)

PhD Thesis, ITP (2012)

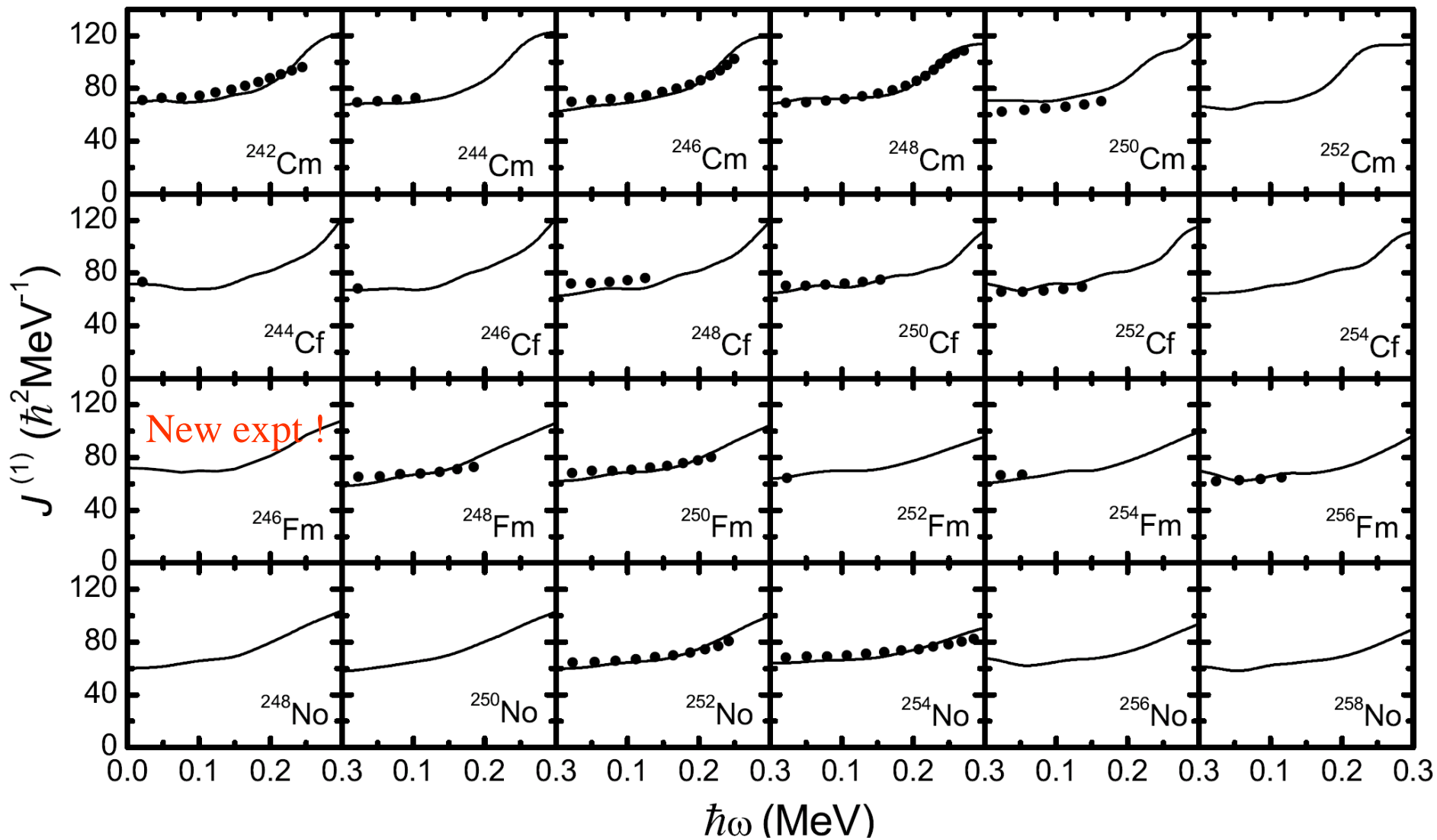
$$H_P(0) = -G_0 \sum_{\xi \eta} a_\xi^\dagger a_\xi^\dagger a_{\bar{\eta}} a_\eta \quad H_P(2) = -G_2 \sum_{\xi \eta} q_2(\xi) q_2(\eta) a_\xi^\dagger a_\xi^\dagger a_{\bar{\eta}} a_\eta$$

N	l	κ_p	μ_p	N	l	κ_n	μ_n
4	0,2,4	0.0670	0.654				
5	1	0.0250	0.710	6	0	0.1600	0.320
	3	0.0570	0.800		2	0.0640	0.200
	5	0.0570	0.710		4,6	0.0680	0.260
6	0,2,4,6	0.0570	0.654	7	1,3,5,7	0.0634	0.318

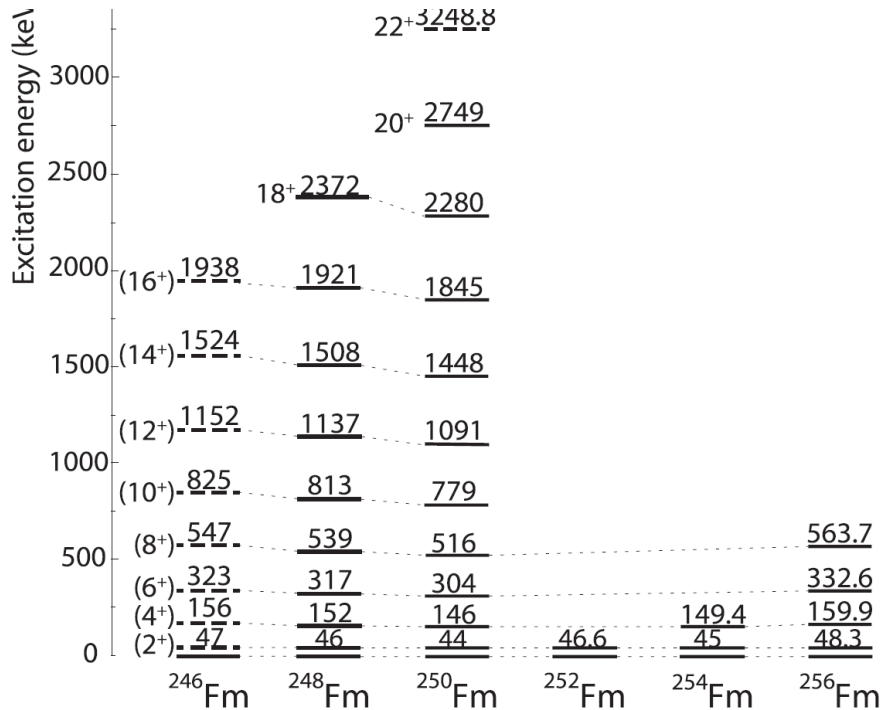
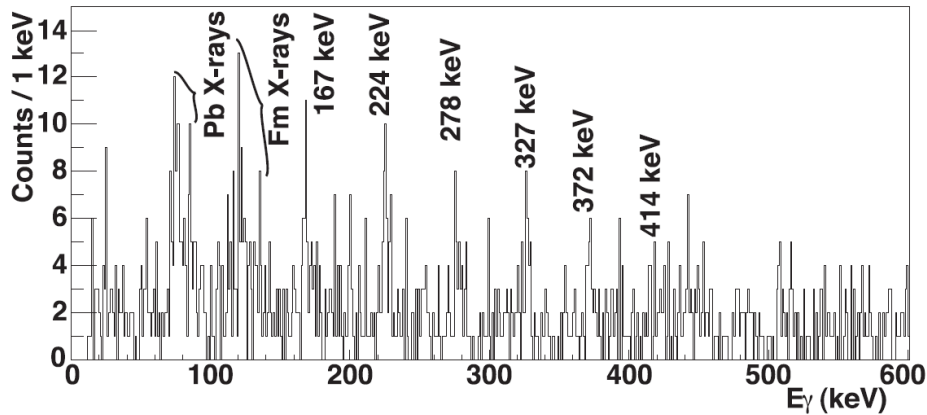
Moments of inertia of even-even nuclei



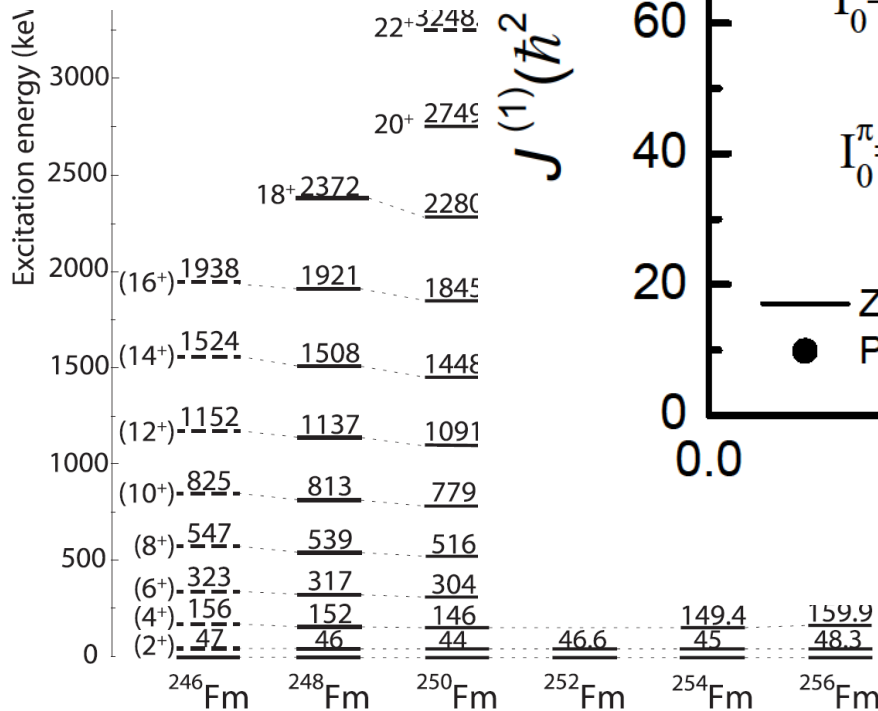
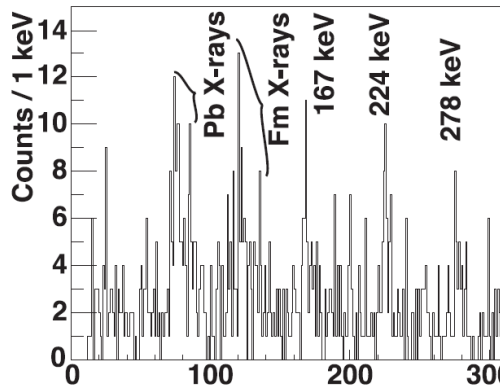
Moments of inertia of even-even nuclei



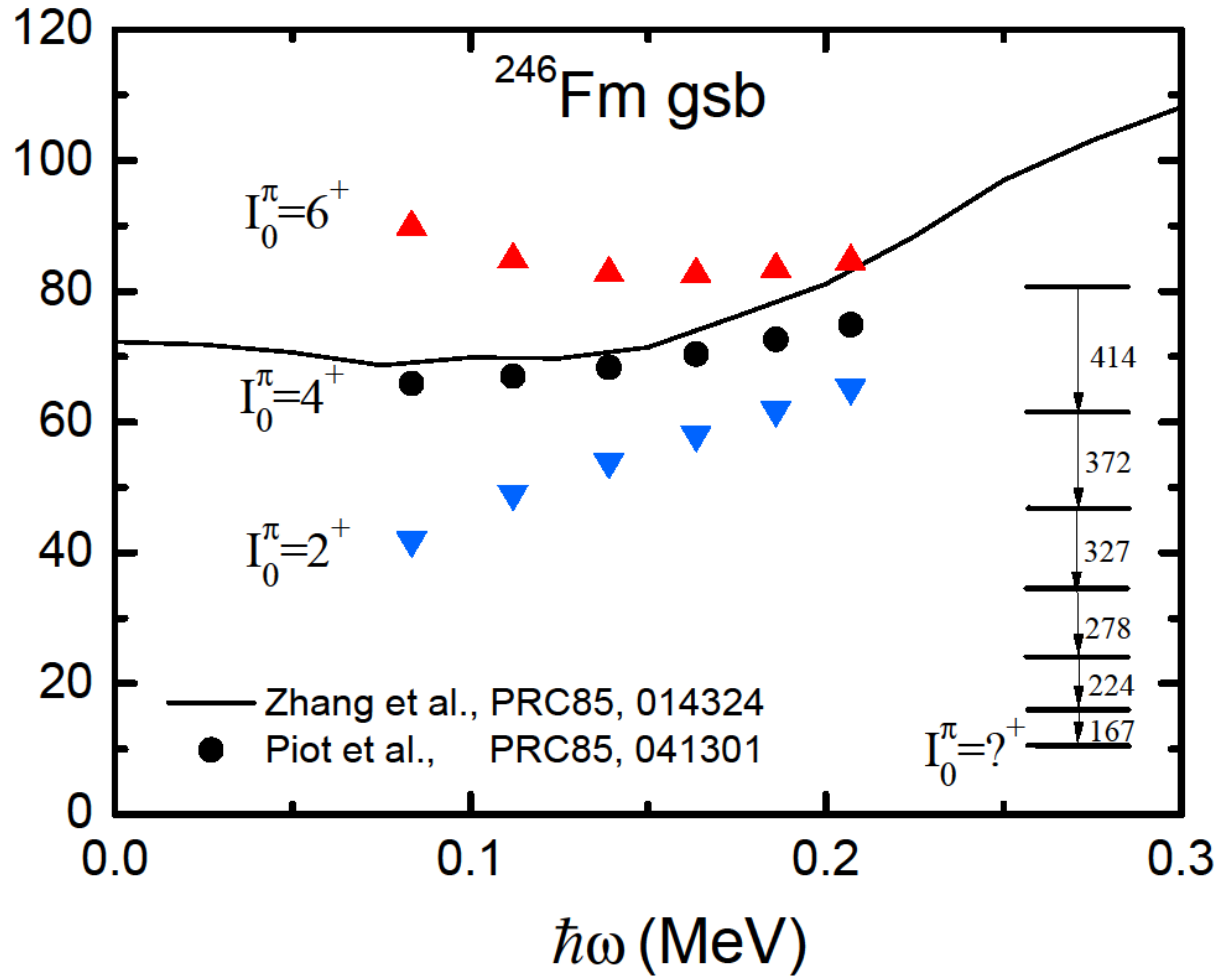
^{246}Fm : ground state band observed @ Jyvaskyla



^{246}Fm : ground state band observed @ Jyvaskyla




$J^{(1)} (\hbar^2 \text{MeV}^{-1})$



^{256}Rf : ground state band observed @ Jyvaskyla

PRL **109**, 012501 (2012)

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Shell-Structure and Pairing Interaction in Superheavy Nuclei: Rotational Properties of the $Z=104$ Nucleus ^{256}Rf

P. T. Greenlees,^{1,*} J. Rubert,² J. Piot,² B. J. P. Gall,² L. L. Andersson,³ M. Asai,⁴ Z. Asfari,² D. M. Cox,³ F. Dechery,⁵
O. Dorvaux,² T. Grahn,¹ K. Hauschild,⁶ G. Henning,^{6,7} A. Herzan,¹ R.-D. Herzberg,³ F. P. Heßberger,⁸
U. Jakobsson,¹ P. Jones,^{1,†} R. Julin,¹ S. Juutinen,¹ S. Ketelhut,¹ T.-L. Khoo,⁷ M. Leino,¹ J. Ljungvall,⁶
A. Lopez-Martens,⁶ R. Lozeva,² P. Nieminen,¹ J. Pakarinen,⁹ P. Papadakis,³ E. Parr,³ P. Peura,¹
P. Rahkila,¹ S. Rinta-Antila,¹ P. Ruotsalainen,¹ M. Sandzelius,¹ J. Sarén,¹ C. Scholey,¹
D. Seweryniak,⁷ J. Sorri,¹ B. Sulignano,⁵ Ch. Theisen,⁵ J. Uusitalo,¹ and M. Venhart¹⁰

^{256}Rf : ground state band observed @ Jyvaskyla

PRL **109**, 012501 (2012)

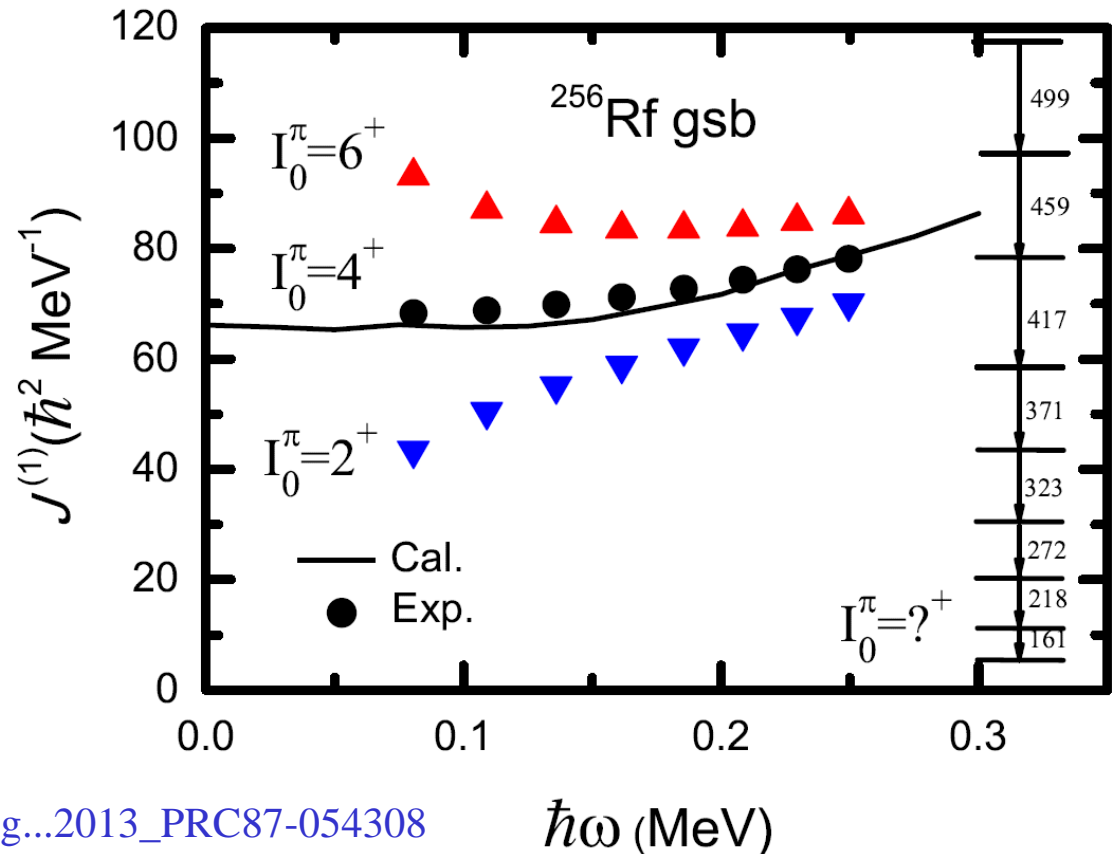
Selected for a **Viewpoint** in *Physics*
PHYSICAL REVIEW LETTERS

week ending
6 JULY 2012



Shell-Structure and Pairing Interaction in Superheavy Nuclei: Rotational Properties of the $Z=104$ Nucleus ^{256}Rf

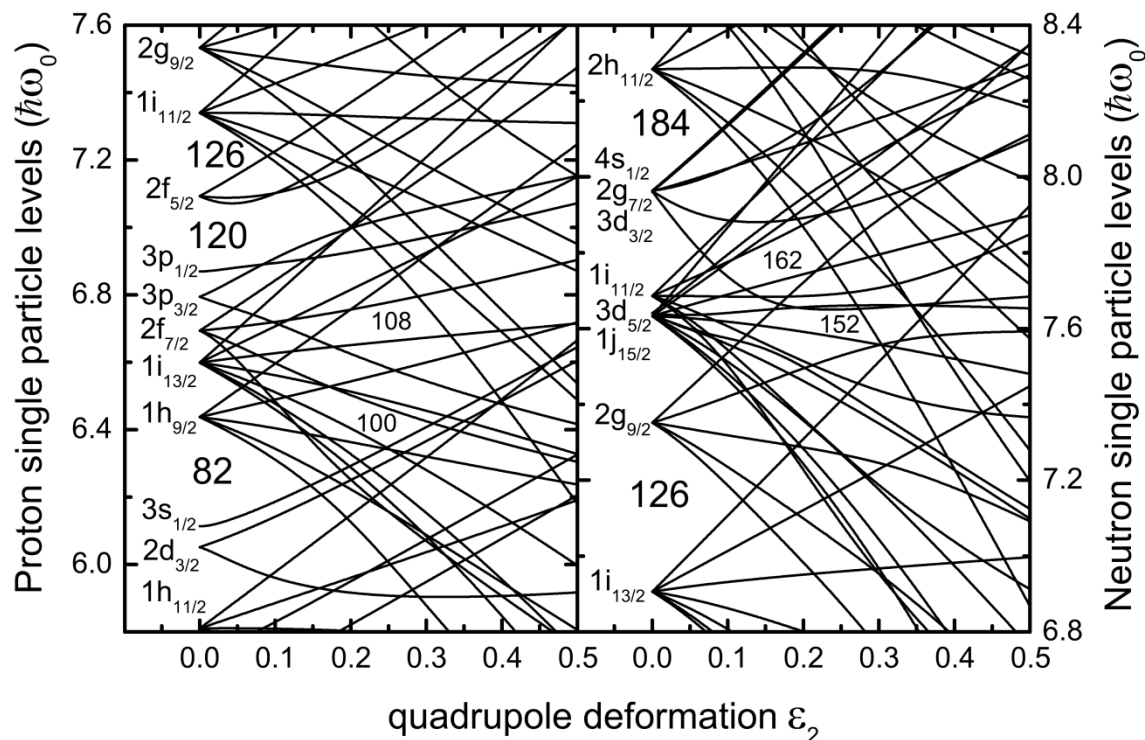
P. T. Greenlees,^{1,*} J. Rubert,² J. Piot,² B. J.
O. Dorvaux,² T. Grahn,¹ K. Hausch
U. Jakobsson,¹ P. Jones,^{1,†} R. Julin,
A. Lopez-Martens,⁶ R. Lozeva,²
P. Rahkila,¹ S. Rinta-Antila,
D. Seweryniak,⁷ J. Sorri,¹ B



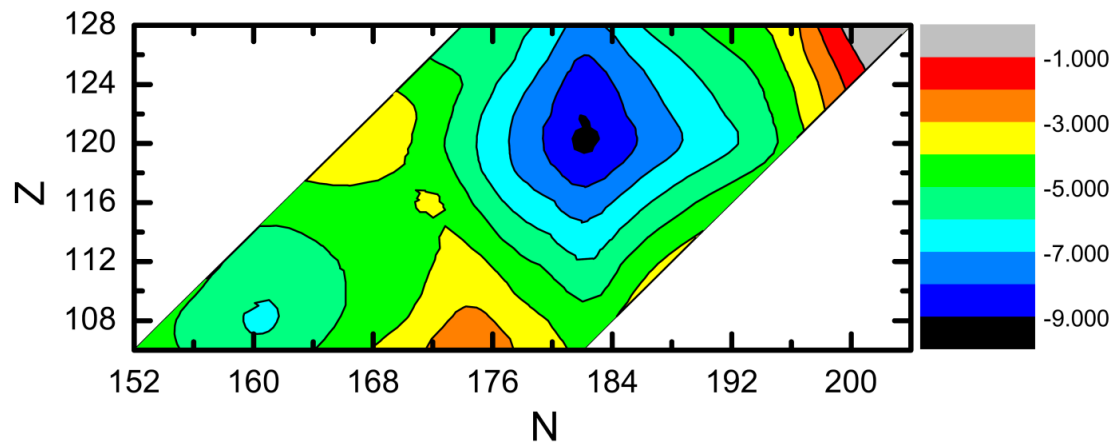
Zhang...2013_PRC87-054308

$\hbar\omega$ (MeV)

Nilsson diagrams & E_{mic}



Z. H. Zhang (张振华)
PhD Thesis, ITP (2012)



Lecture 2

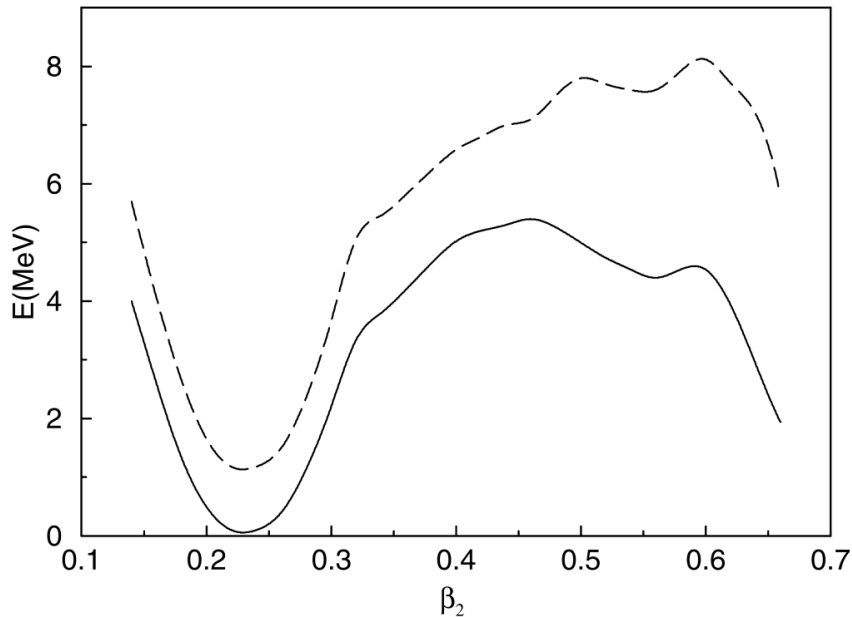
- Challenges in synthesizing SHN

- Theoretical study of structure of SHN
 - Nuclear models
 - Next shell closures beyond ^{208}Pb as seen from single particle spectra, shell correction energy & nuclear shapes
 - Exotic shapes in SHN
 - Low-lying spectra of SHN & magicities

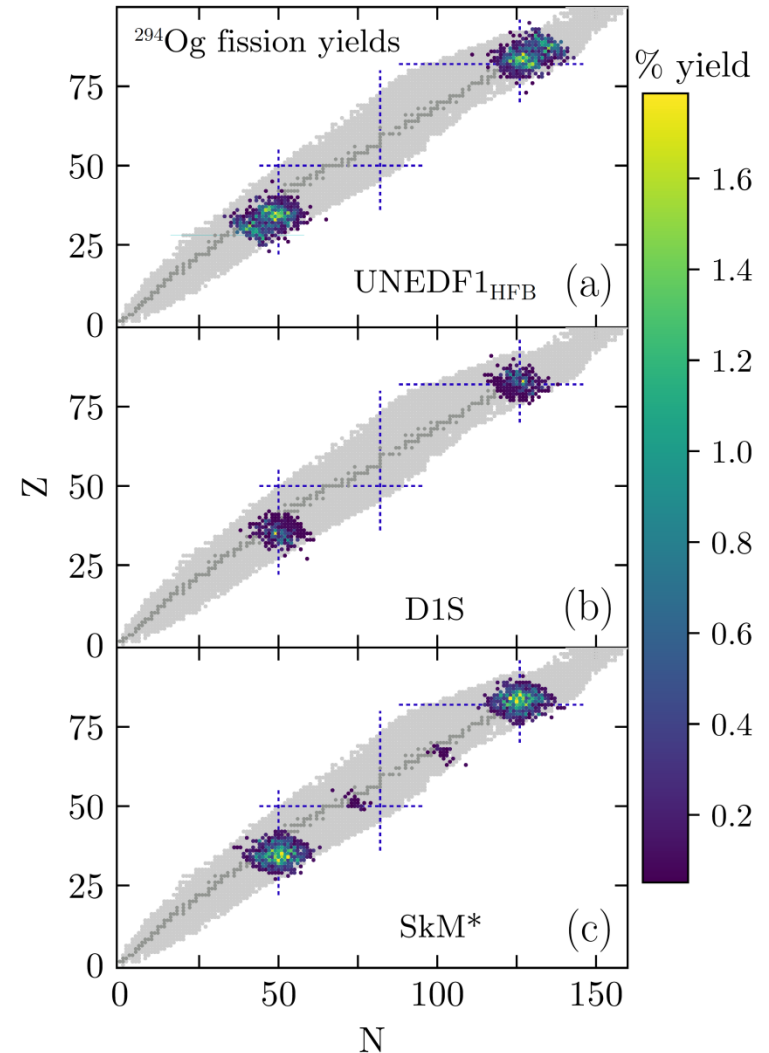
- Theoretical study of decay of SHN

Decay & spontaneous fission of SHN

- α , β & γ decays
- Spontaneous fission
- Long-lived isomers
- Cluster radioactivity



Xu_Zhao_Wyss&Walker2004_PRL92-252501



Matheson...2018_arXiv1812.06490

Zhang&Wang2018_PRC97-014318

Lectures 3 & 4

- Challenges in synthesizing SHN

- Synthesis mechanism of SHN
 - Large uncertainties in predicted Xsections
 - Heavy ion fusion reactions
 - Capture
 - Fusion
 - Survival against fission
 - Multi-nucleon transfer reactions
 - Reactions w/ radioactive ion beams

Lectures 3 & 4

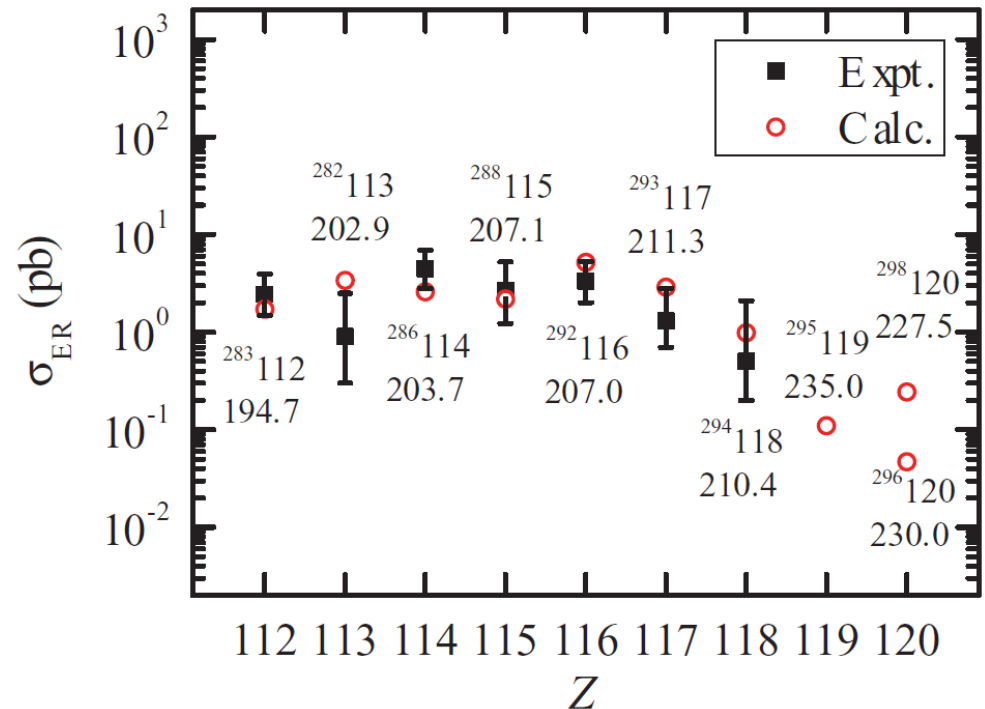
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How to synthesize SHE w/ $Z > 118$?

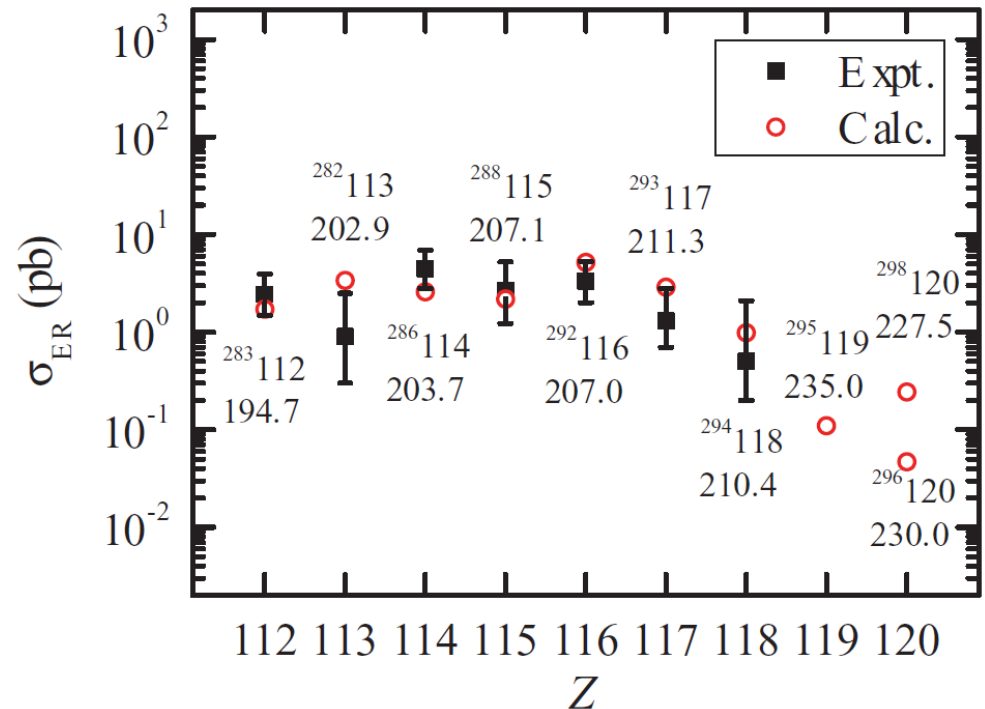


How to synthesize SHE w/ $Z > 118$?



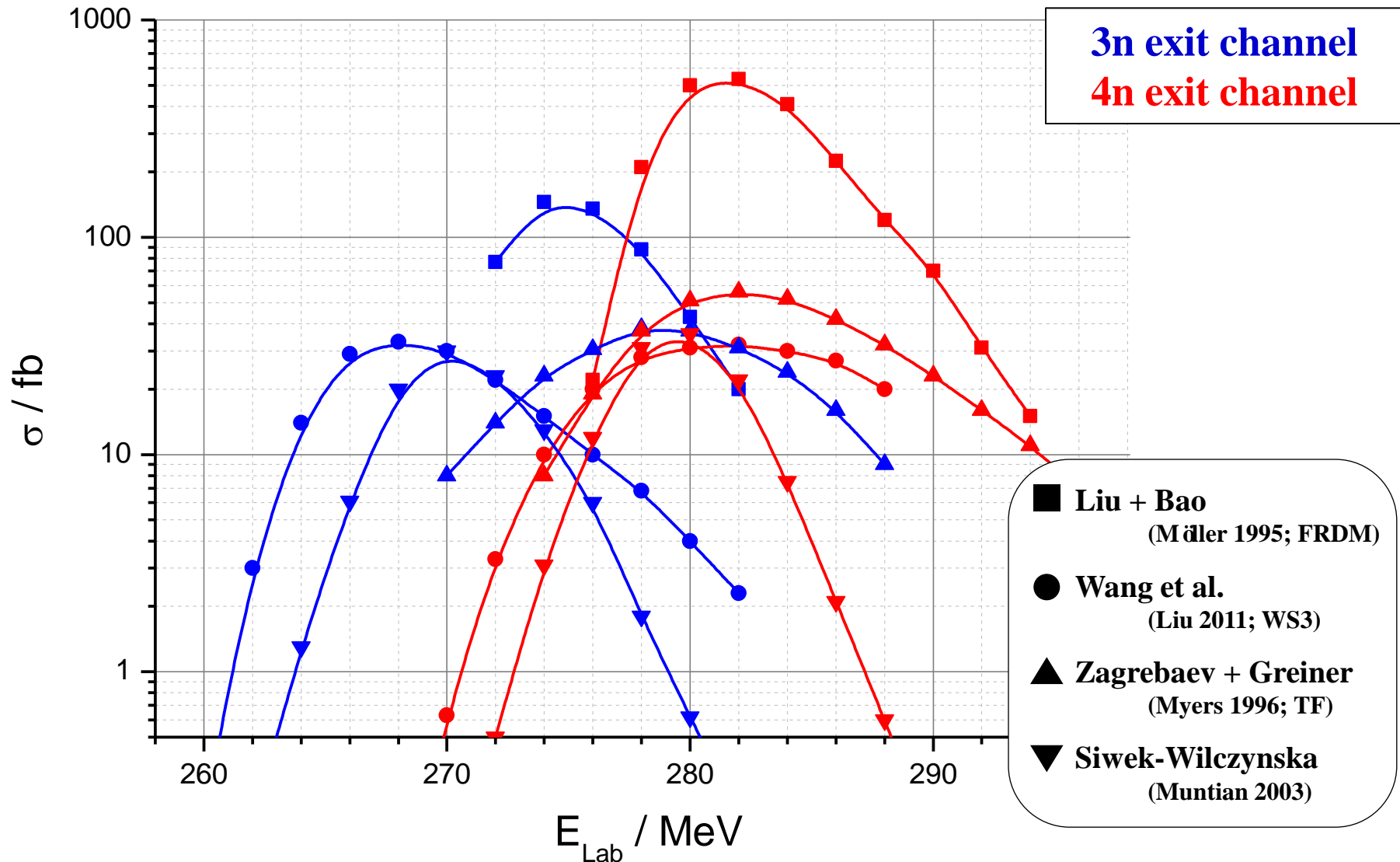
How to synthesize SHE w/ $Z > 118$?

Feng_Jin_Li_Scheid2007_PRC76-044606
 Nasirov_Giardina_Mandaglio_Manganaro_Hanappe_
 Heinz_Hofmann_Muminov_Scheid2009_PRC79-024606
 Adamian_Antonenko_Scheid2009_EPJA41-235
 Gan_Zhou_Huang_Feng_Li2011_SciChinaPMA54S1-61
 Nasirov_Mandaglio_Giardina_Sobiczewski_Muminov2011_PRC84-044612
 Liu_Bao2009_PRC80-054608
 Siwek-Wilczynska_Cap_Wilczynski2010_IJMPE19-500
 Liu_Bao2011_PRC83-044613
 Liu_Bao2011_PRC84-031602R
 Zagrebaev_Greiner2008_PRC78-034610
 Wang_Tian_Scheid2011_PRC84-061601R
 Siwek-Wilczynska_Cap_Kowal_Sobiczewski_
 Wilczynski 2012_PRC86-014611
 Liang_Zhu_Liu_Wang 2012_PRC86-037602
 Liu_Bao2013_PRC87-034616
 Zhang_Wang_Ren2013_NPA909-36
 Zhu_Xie_Zhang2014_PRC89-024615
 Bao_Gao_Li_Zhang2015_PRC91-011603R
 Bao_Gao_Li_Zhang2015_PRC92-034612
 Liu_Shen_Li_Tu_Wang_Wang2016_EPJA52-35
 Santhosh_Safoora2016_PRC94-024623
 Ghahramany_Ansari2016_EPJA52-287
 Hong_Adamian_Antonenko2016_EPJA52-305
 Santhosh_Safoora2017_PRC96-034610
 Adamian_Antonenko_Lenske2018_NPA970-22



Wang_Zhao_Scheid_SGZ 2012 PRC85-041601R

A big issue in synthesizing SHE w/ $Z > 118$



From Ch. E. Dullman (FUSHE2012)

Lectures 3 & 4

- Challenges in synthesizing SHN

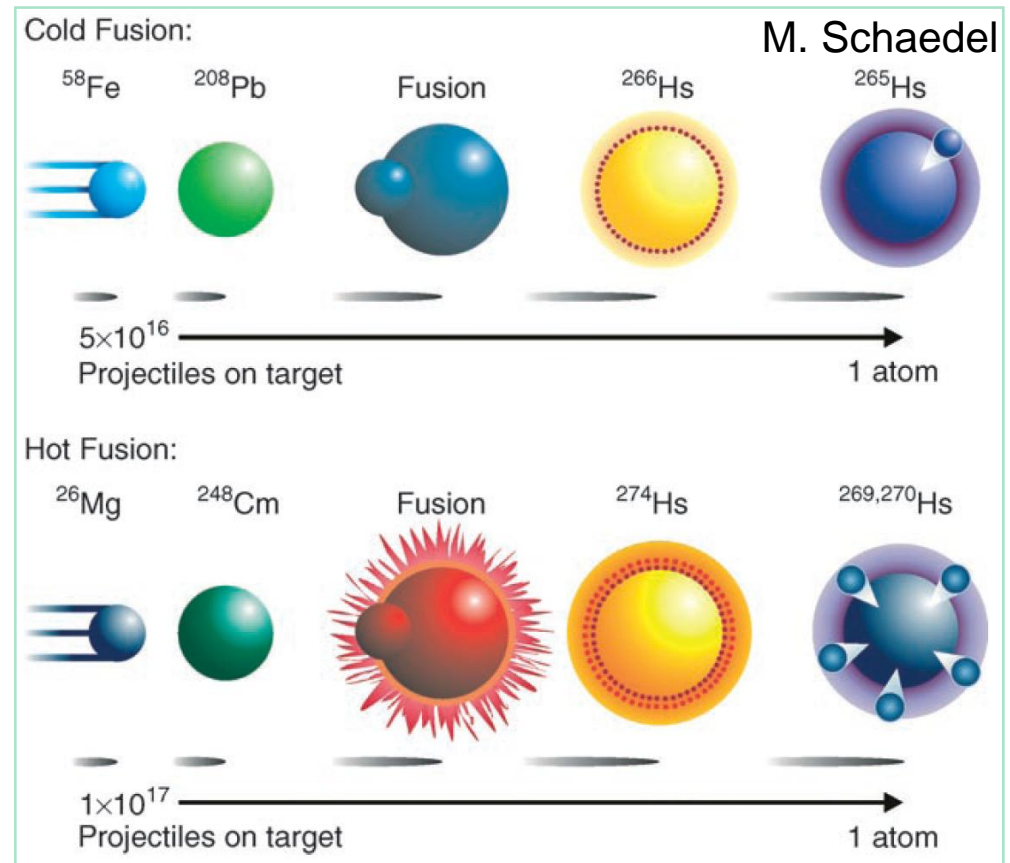
- Synthesis mechanism of SHN
 - Large uncertainties in predicted Xsections
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 - Capture
 - Fusion
 - Survival against fission
 - Multi-nucleon transfer reactions
 - Reactions w/ radioactive ion beams

Three steps to a SHN via heavy-ion EvR reaction

□ Capture

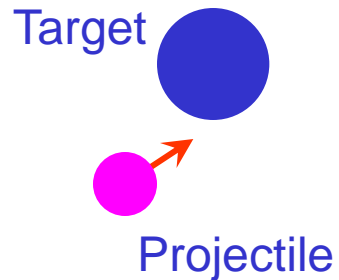
□ Formation of CN

□ Deexcitation of CN

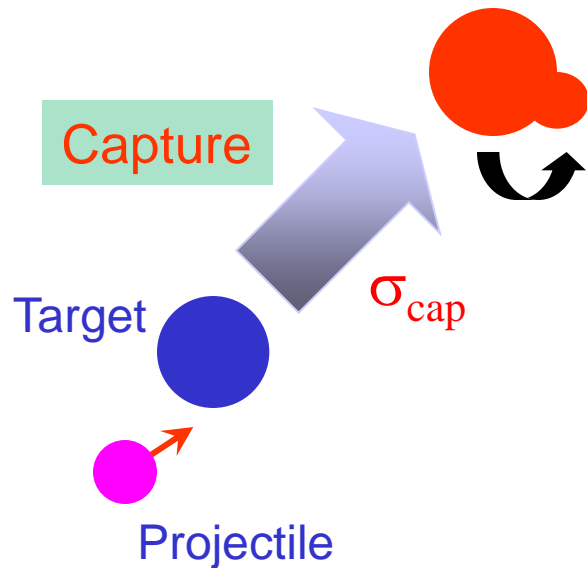


$$\sigma_{\text{ER}}(E_{\text{cm}}) = \sum_J \sigma_{\text{cap}}(E_{\text{cm}}, J) P_{\text{CN}}(E_{\text{cm}}, J) W_{\text{sur}}(E_{\text{cm}}, J)$$

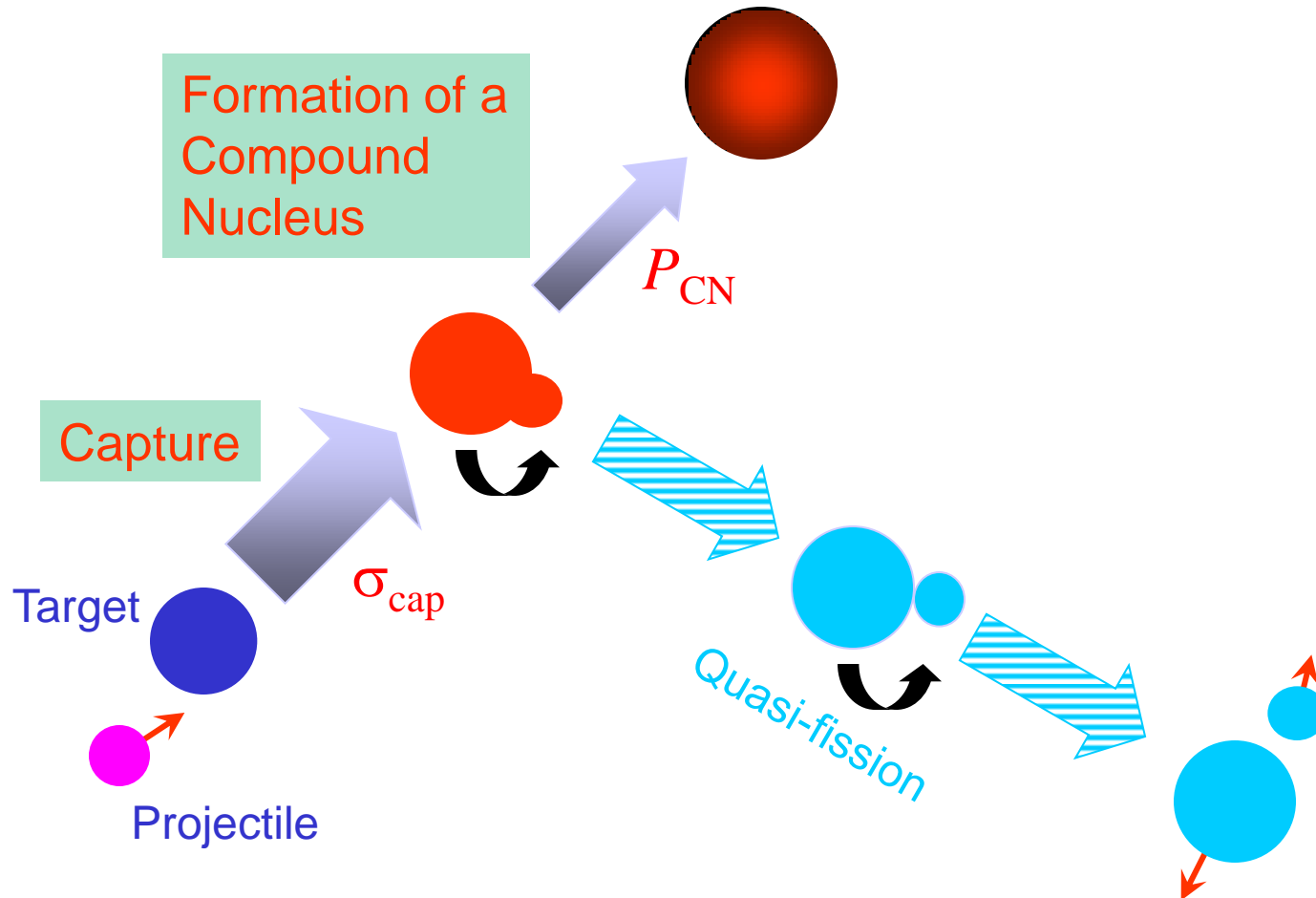
Three steps to a SHN via heavy-ion EvR reaction



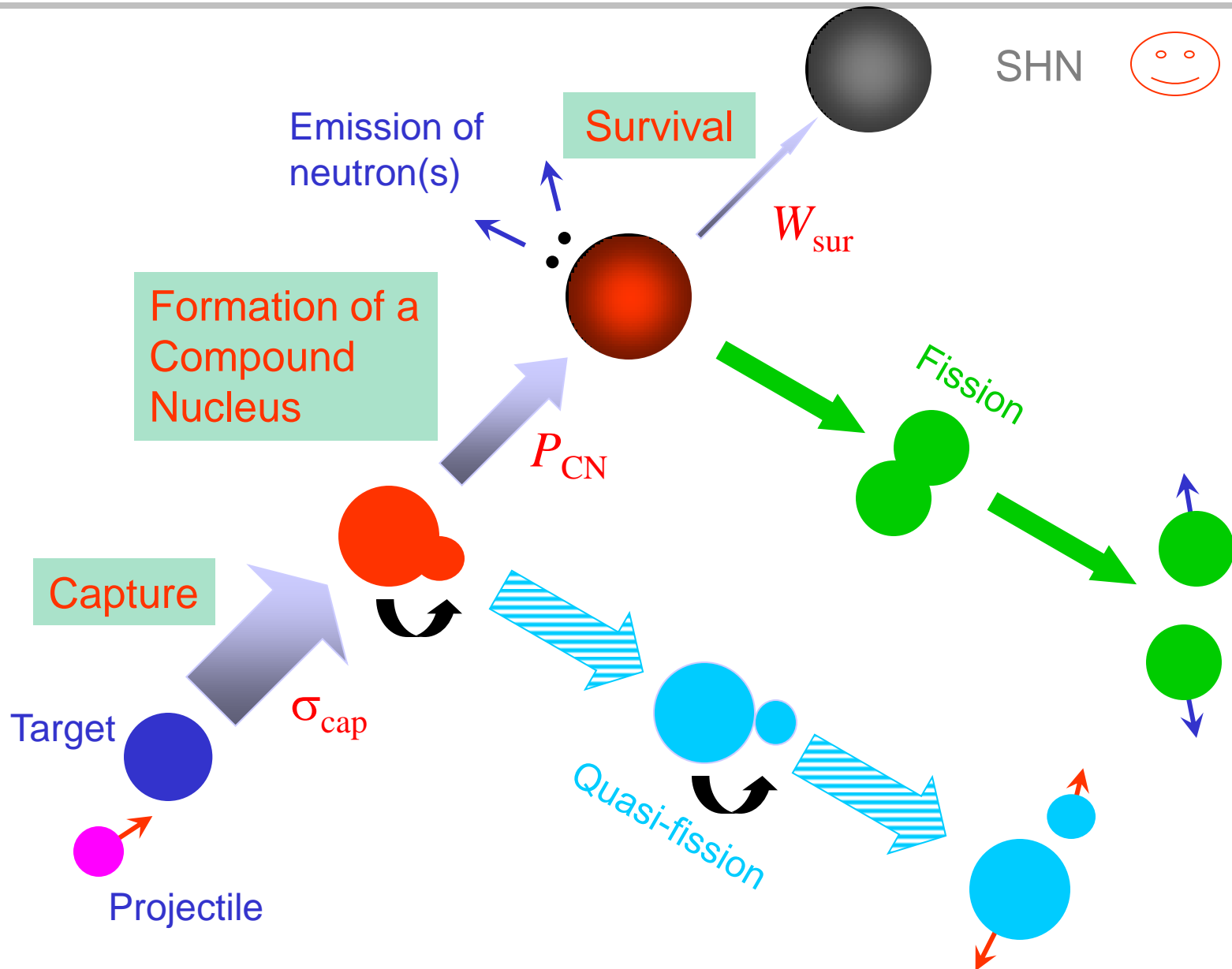
Three steps to a SHN via heavy-ion EvR reaction



Three steps to a SHN via heavy-ion EvR reaction



Three steps to a SHN via heavy-ion EvR reaction



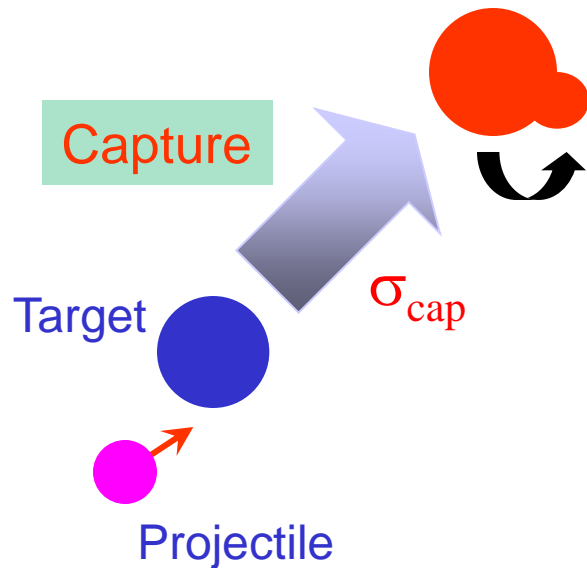
Lectures 3 & 4

- Challenges in synthesizing SHN

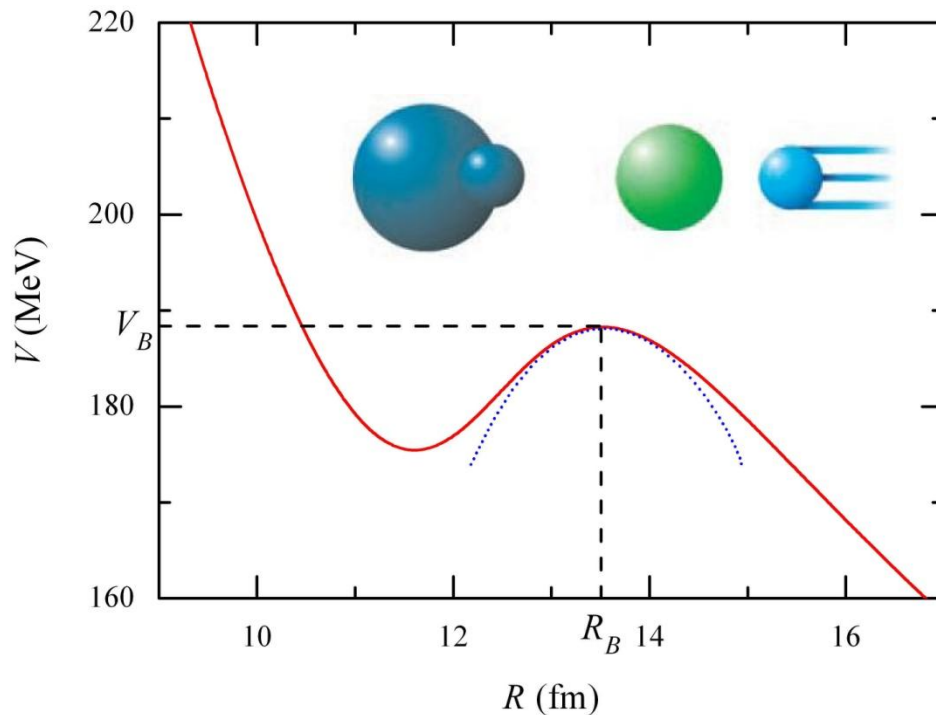
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Three steps to a SHN via heavy-ion EvR reaction

σ_{cap}



The capture process



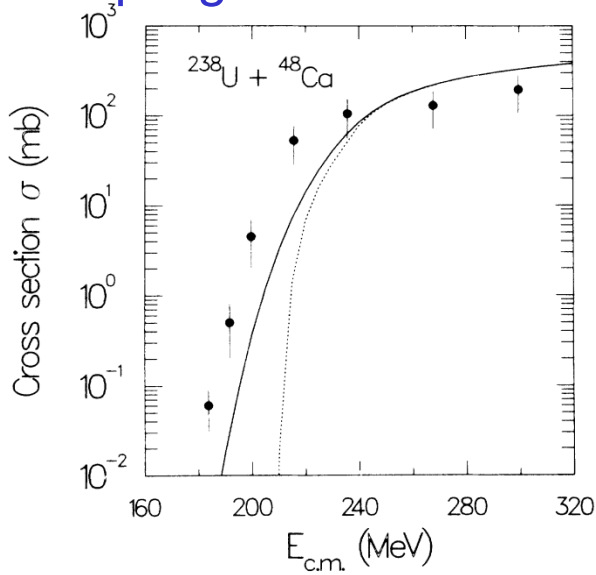
- Path integral method
- WKB approximation
- Hill-Wheeler formula
- New formula by Li et al.
- ...

$$\sigma_{\text{cap}}(E_{\text{cm}}, J) = \frac{\pi \hbar^2}{2\mu E_{\text{cm}}} (2J + 1) T(E_{\text{cm}}, J)$$

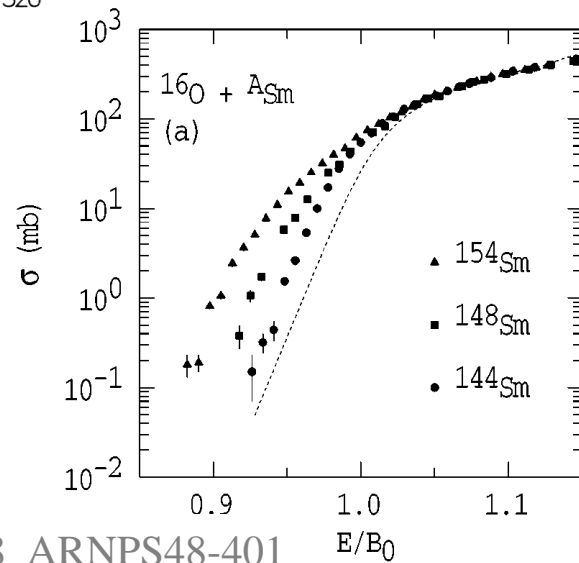
$$T(E_{\text{cm}}, J) = \left(1 + \exp \left[-\frac{2\pi}{\hbar\omega} (E_{\text{cm}} - E_{\text{rot}} - V_B) \right] \right)^{-1}$$

Channel coupling effects

Coupling effects due to rotation, vibration, nucleon transfer, ...



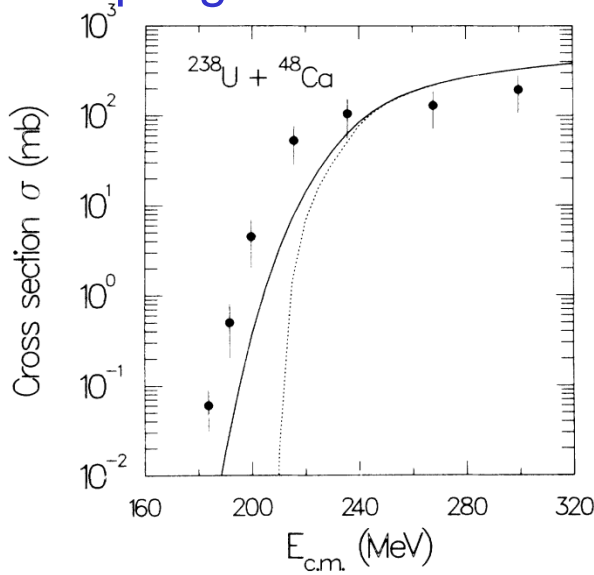
Shen+1987_PRC36-115



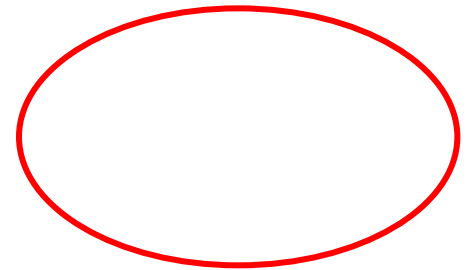
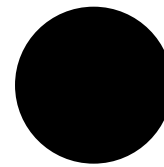
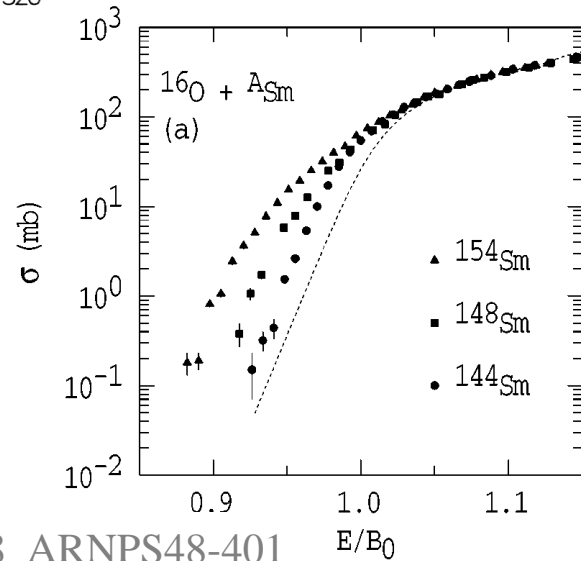
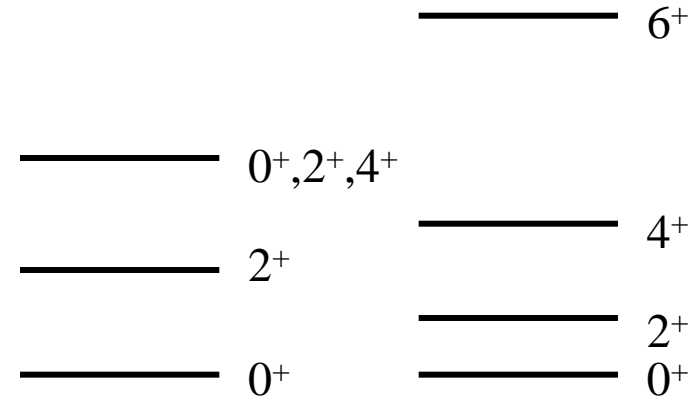
Dasgupta+1998_ARNPS48-401

Channel coupling effects

Coupling effects due to rotation, vibration, nucleon transfer, ...



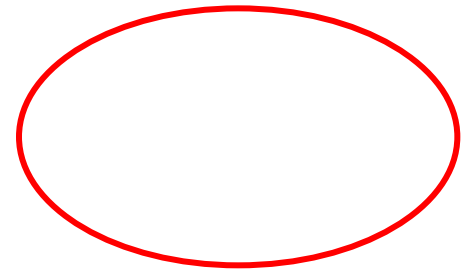
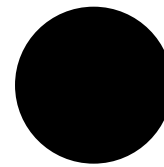
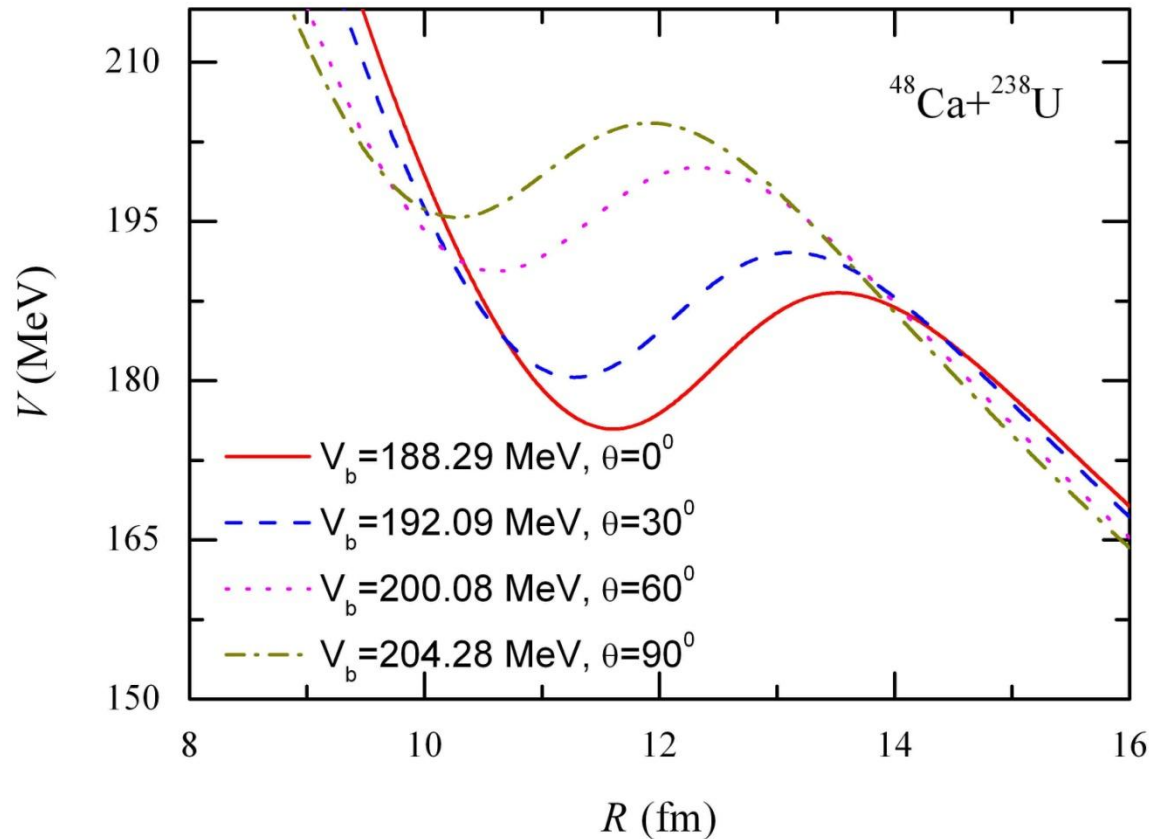
Shen+1987_PRC36-115



Dasgupta+1998_ARNPS48-401

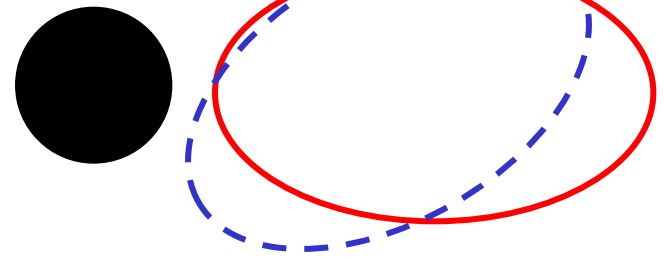
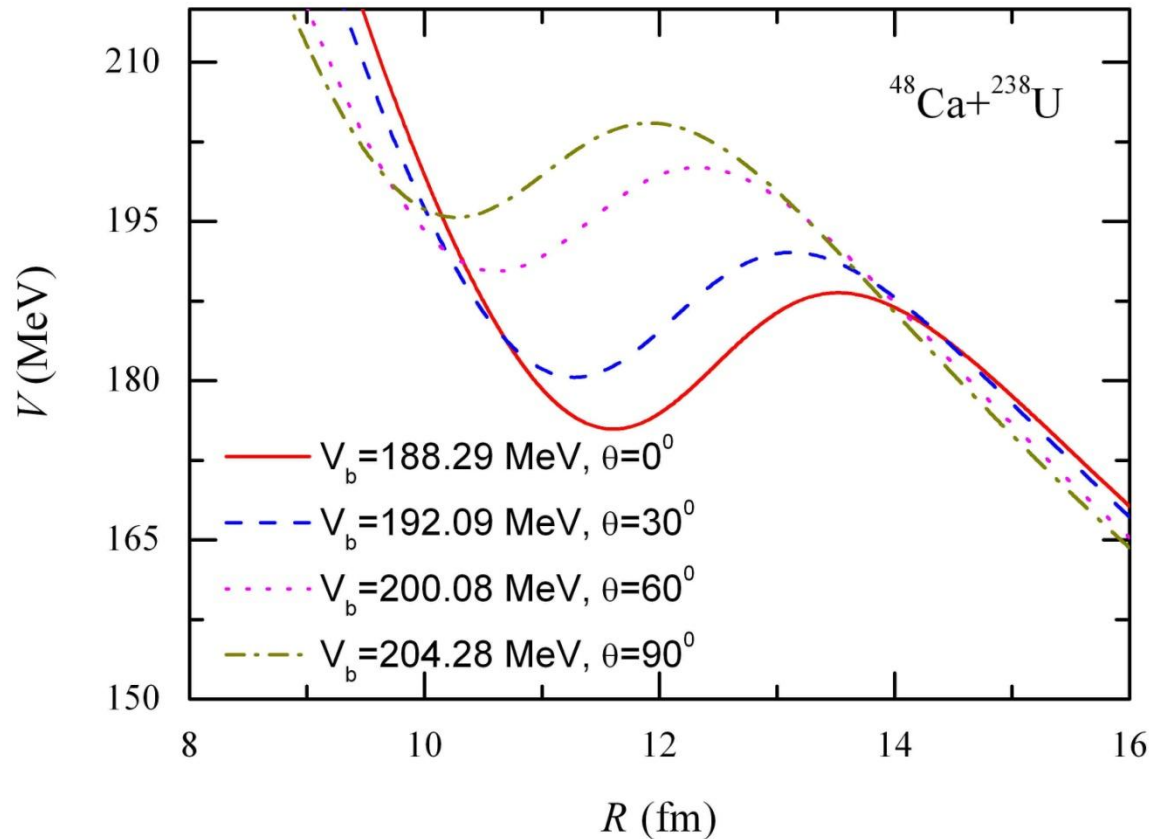
Barrier distribution

Coupling effects due to rotation, vibration, nucleon transfer, ...
are taken into account **empirically** by introducing a barrier distribution



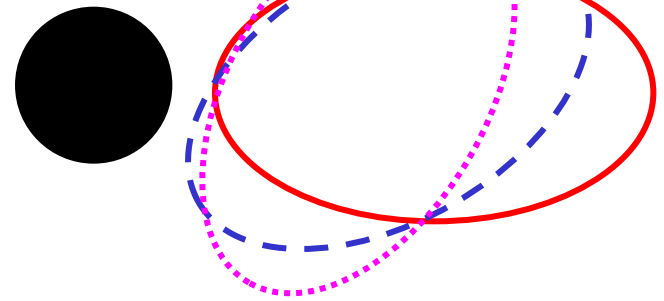
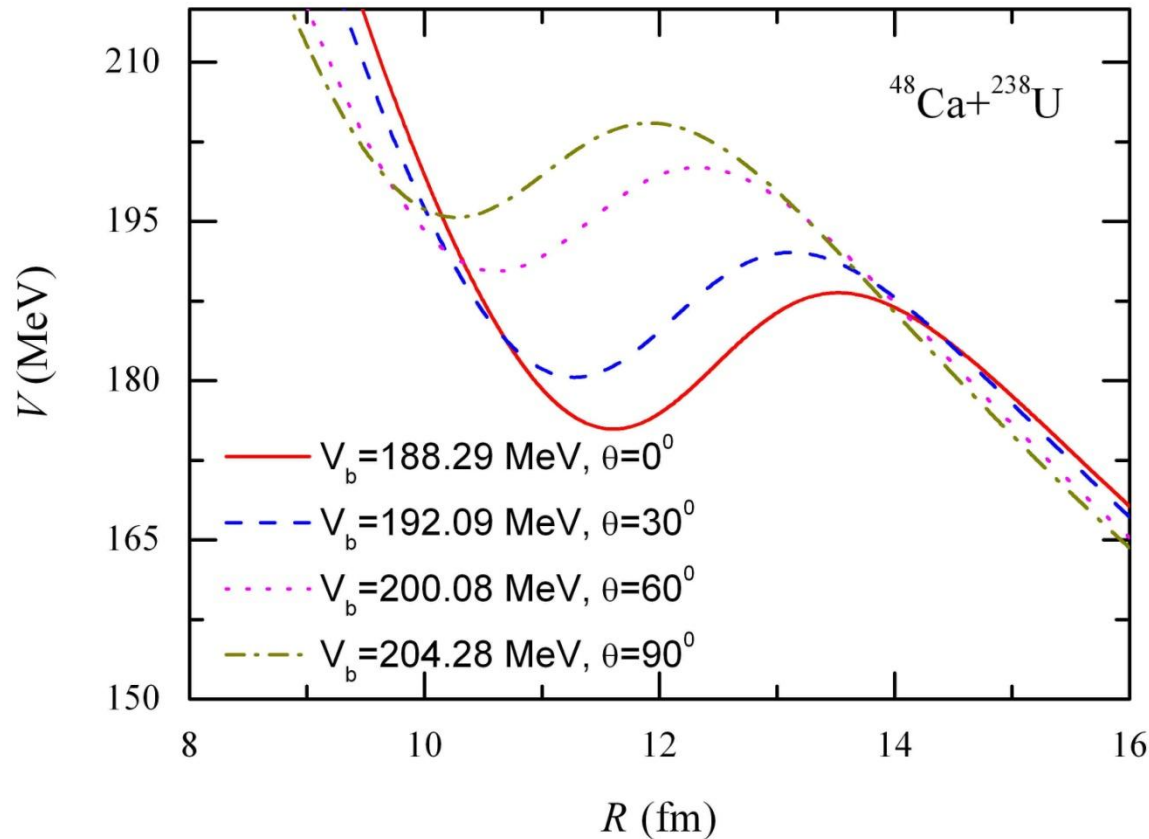
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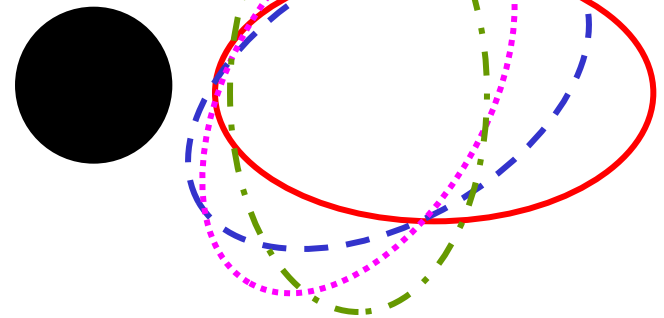
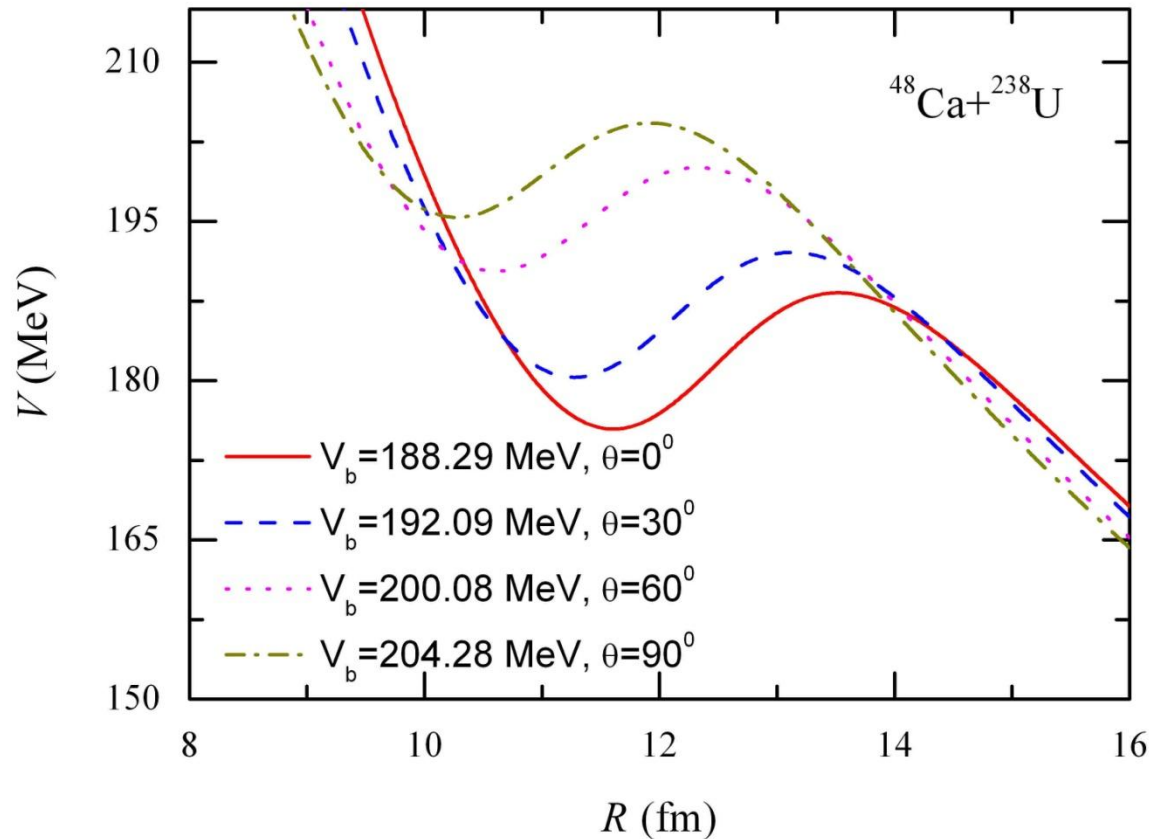
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Barrier distribution

Coupling effects due to rotation, vibration, nucleon transfer, ...
are taken into account **empirically** by introducing a barrier distribution



$$T(E_{\text{cm}}, J) = \int dB f(B) T(E_{\text{cm}}, B, J)$$

The empirical CC (ECC) model

□ Capture cross section

$$\sigma_{\text{capture}}(E) = \frac{\pi \hbar^2}{2\mu E} \sum_J^{J_{\text{max}}} (2J + 1) T(E, J)$$

$$T(E, J) = \int f(B) T_{\text{HW}}(E, J, B) dB$$

□ Barrier distribution

$$f(B) = \begin{cases} \frac{1}{N} \exp \left[- \left(\frac{B - B_m}{\Delta_1} \right)^2 \right], & B < B_m \\ \frac{1}{N} \exp \left[- \left(\frac{B - B_m}{\Delta_2} \right)^2 \right], & B > B_m \end{cases}$$

$$B_m = aV_B^{\text{Sp}} + (1 - a)V_B^{\text{S}},$$

$$\Delta_1 = bE_D,$$

$$\Delta_2 = cE_D,$$

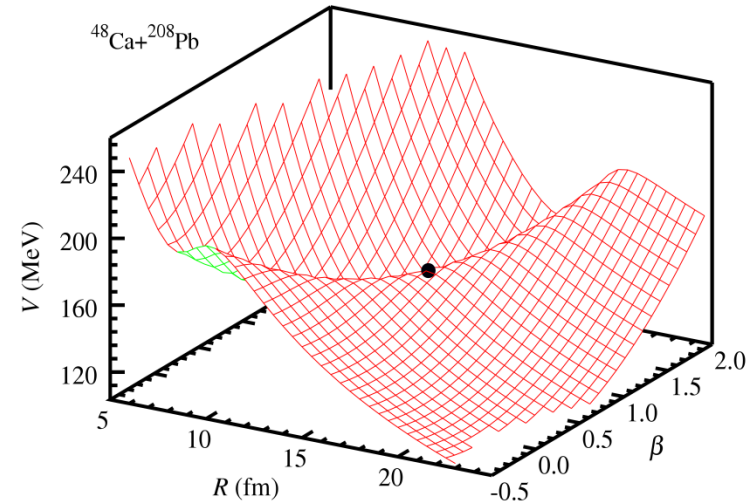
$$E_D = \frac{1}{2} \sum_{i=\text{P,T}} C_i (\beta_i^{\text{S}})^2$$

For reactions w/ positive Q-values for 2n transfer

$$\Delta_1 \rightarrow gQ(2n) + \Delta_1$$

$$\Delta_2 \rightarrow gQ(2n) + \Delta_2$$

Wang, Wen, Zhao, Zhao & SGZ
At. Data & Nucl. Data Tables 114 (2017) 281



Model parameters

□ Nucleus-nucleus potential

Wang, Wen, Zhao, Zhao & SGZ
At. Data & Nucl. Data Tables 114 (2017) 281

$$V_N(R, \beta_P^0, \beta_T^0, \theta_P, \theta_T) = \frac{-V_0}{1 + \exp[(R - R_P - R_T)/a]},$$

with

$$\begin{aligned} R_P &= R_{0P}[1 + (5/4\pi)^{1/2} \beta_P^0 P_2(\cos \theta_P)], & V_0 &= 80 \text{ MeV}, \\ R_T &= R_{0T}[1 + (5/4\pi)^{1/2} \beta_T^0 P_2(\cos \theta_T)], & r_0 &= 1.16 \text{ fm}, \\ R_{0i} &= r_0 A_i^{1/3}, \quad i = P, T, & a &= \left\{ 1.17 \left[1 + 0.53 \left(A_P^{-1/3} + A_T^{-1/3} \right) \right] \right\}^{-1} \text{ fm}. \end{aligned}$$

Deformation parameters: [Möller_Nix_Myers_Swiatecki_ADNDT59-185](#)

□ Barrier distribution

Reactions w/ spherical T & P

$$a = \begin{cases} 0.26, & Z_P Z_T < 1150, \\ 0.5, & Z_P Z_T \geq 1150, \end{cases}$$

$$b = 0.32,$$

$$c = 0.93.$$

Reactions w/ deformed T or P

$$a = \begin{cases} 0.23, & Z_P Z_T < 1150, \\ 0.37, & Z_P Z_T \geq 1150, \end{cases}$$

$$b = 0.12,$$

$$c = 1.12.$$

$$g = 0.32$$

Fusion reactions w/ well bound projectiles

□ 220 reactions with $182 \leq Z_p Z_T \leq 1640$

Wang, Wen, Zhao, Zhao & SGZ
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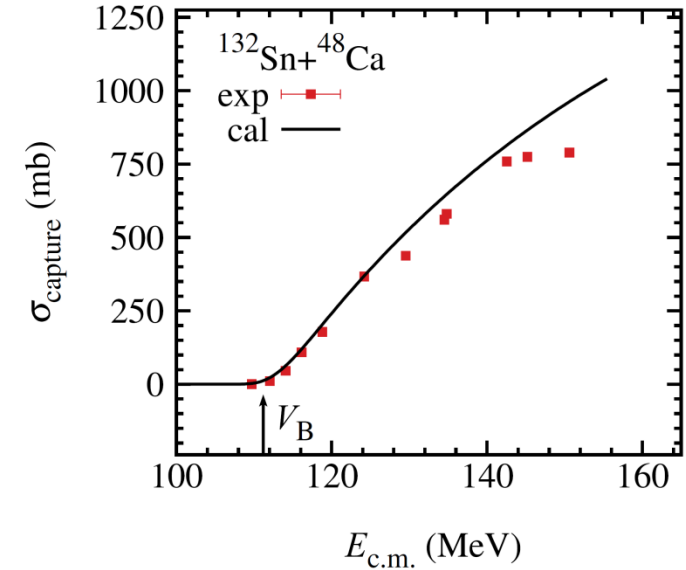
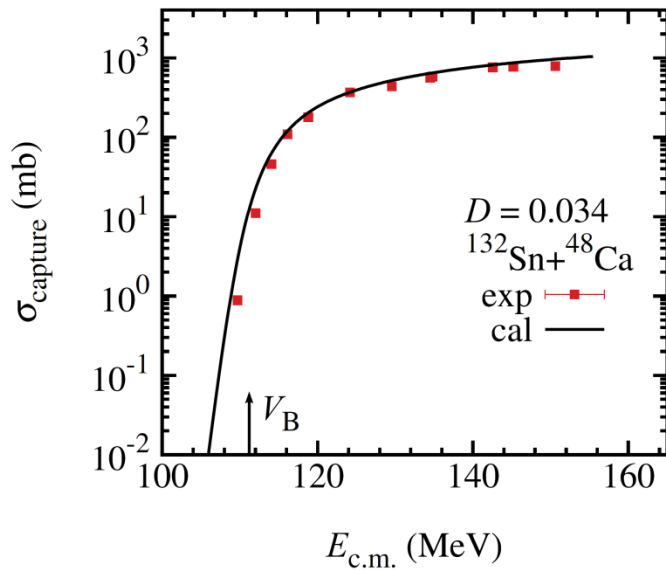
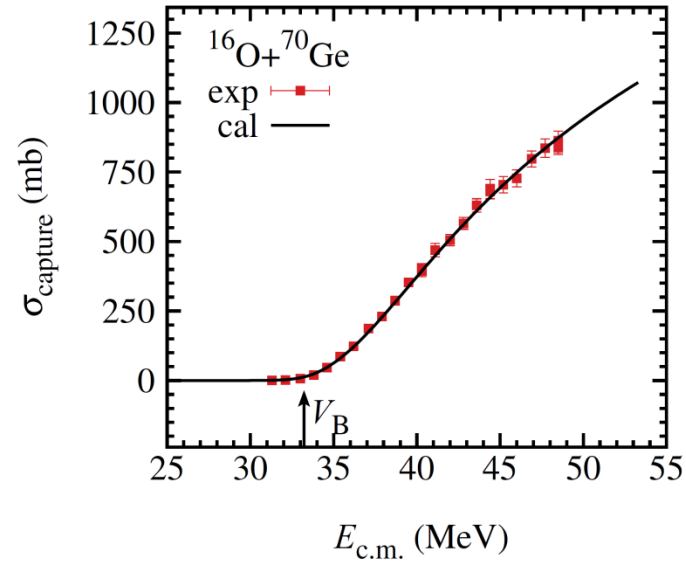
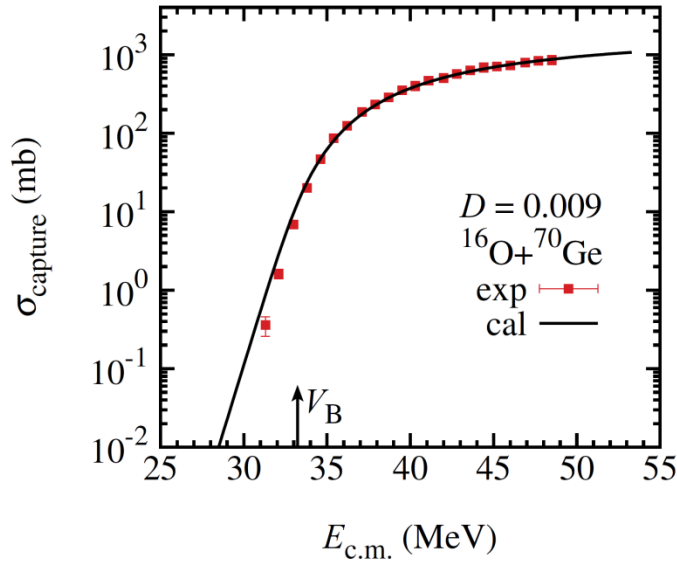
Table 2

The experimental and calculated excitation functions for a total of 217 reaction systems. See page 56 for Explanation of Tables.

Reaction	Detected particles	Data obtained	E_{lab} (MeV)	$E_{c.m.}$ (MeV)	σ (mb)	$\delta\sigma$ (mb)	σ^{cal} (mb)
$^{12}\text{C}+^{89}\text{Y}$	EvR	authors' graph [103] E_{lab}	30.682	27.037	0.033	—	0.195
			32.403	28.554	0.668	—	1.697
			33.789	29.774	4.757	—	9.135
			35.387	31.183	29.865	—	47.223
			36.912	32.526	96.919	—	132.929
			39.174	34.520	287.692	—	302.095
			41.019	36.145	369.296	—	434.894
			43.207	38.074	597.752	—	575.375
			46.281	40.783	665.270	—	744.850
			$^{12}\text{C}+^{92}\text{Zr}$	EvR	authors' table [104] $E_{c.m.}$	31.88	28.20
32.88	29.09	1.45				0.2	1.613
33.89	29.98	2.91				0.2	5.554
34.89	30.86	13.2				0.5	17.351
35.89	31.75	38.6				1.1	45.991
36.89	32.63	83.6				1.2	95.738
37.89	33.52	136				3	163.147
38.89	34.40	197				2	237.855
39.90	35.30	253				3	315.353
40.90	36.18	308				3	388.645
41.91	37.07	366				3	459.032
42.90	37.95	421				3	524.753
43.89	38.83	476				4	586.786
45.90	40.60	570				5	701.389
47.91	42.38	664				5	804.554
49.91	44.15	731	9	896.774			
$^{12}\text{C}+^{144}\text{Sm}$	EvR	authors' graph [105] E_{lab}	46.464	42.890	0.377	—	0.815
			46.968	43.355	0.975	—	1.489
			47.407	43.760	1.458	—	2.508
			47.807	44.110	2.888	—	4.484

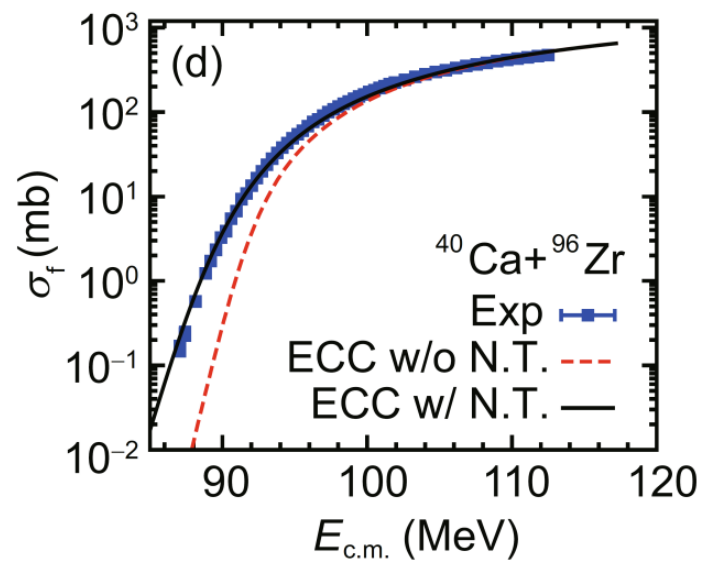
“Good” examples

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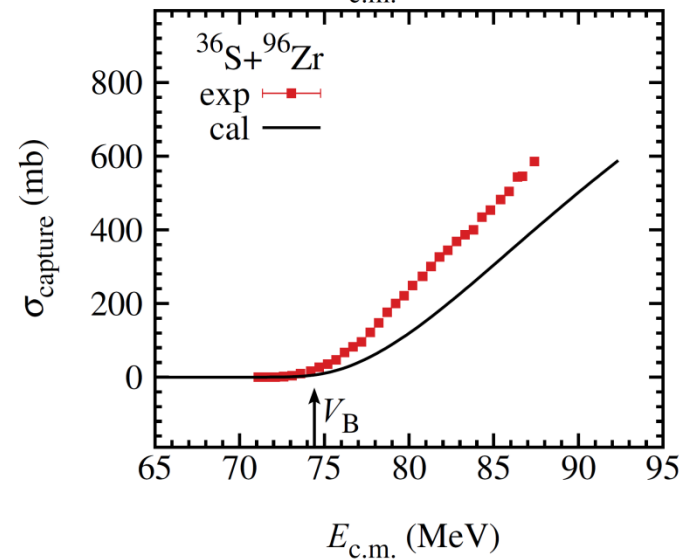
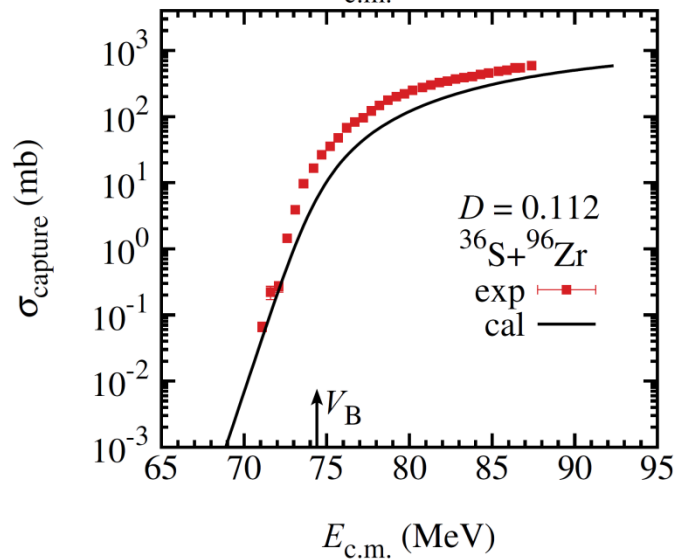
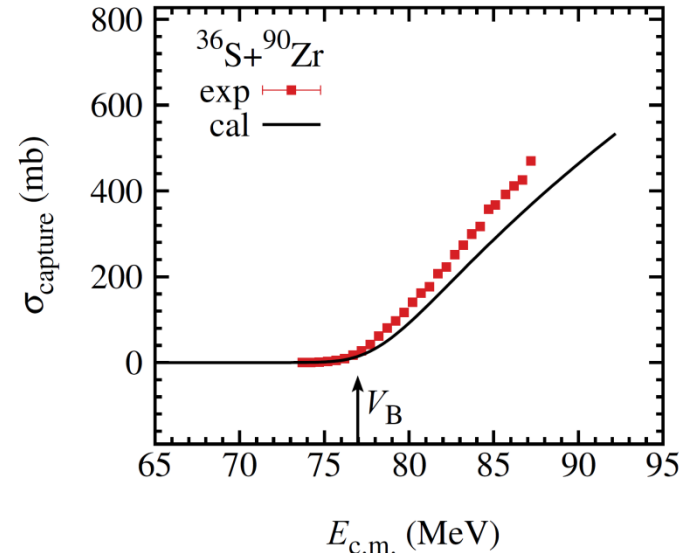
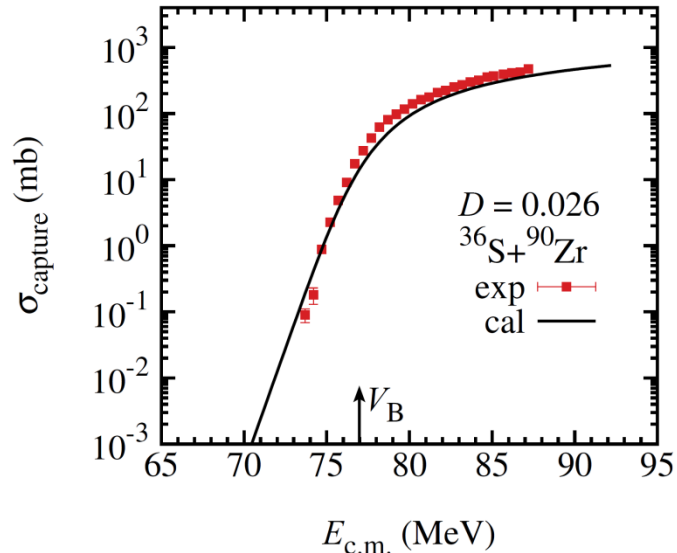
$^{32}\text{S}(^{40}\text{Ca})+^{96}\text{Zr}$

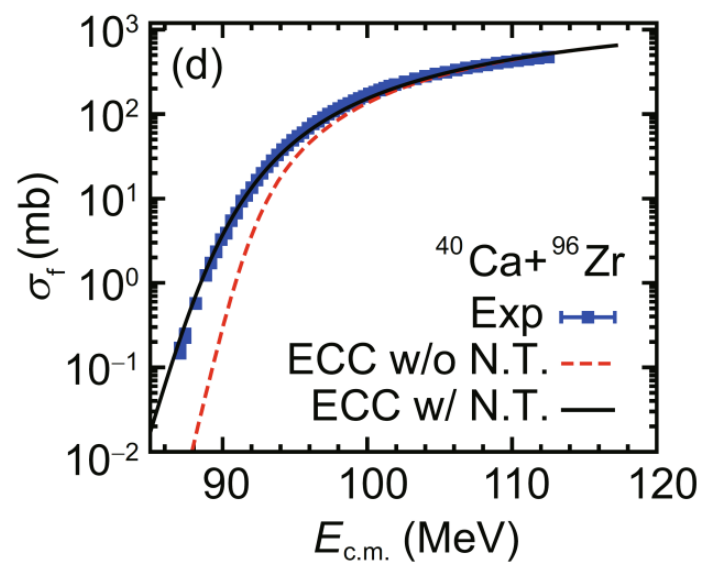
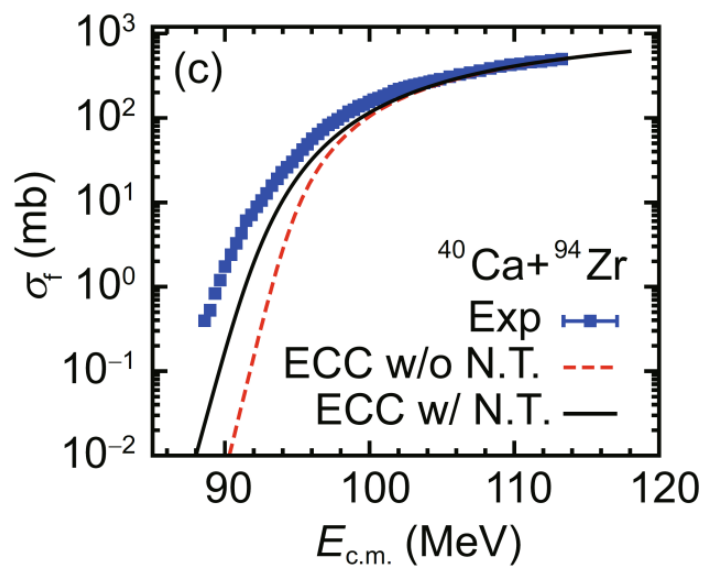
Wang, Zhao, Zhao & SGZ
Sci. China-Phys. Mech. & Astron. 59 (2016) 642002



“Bad” examples

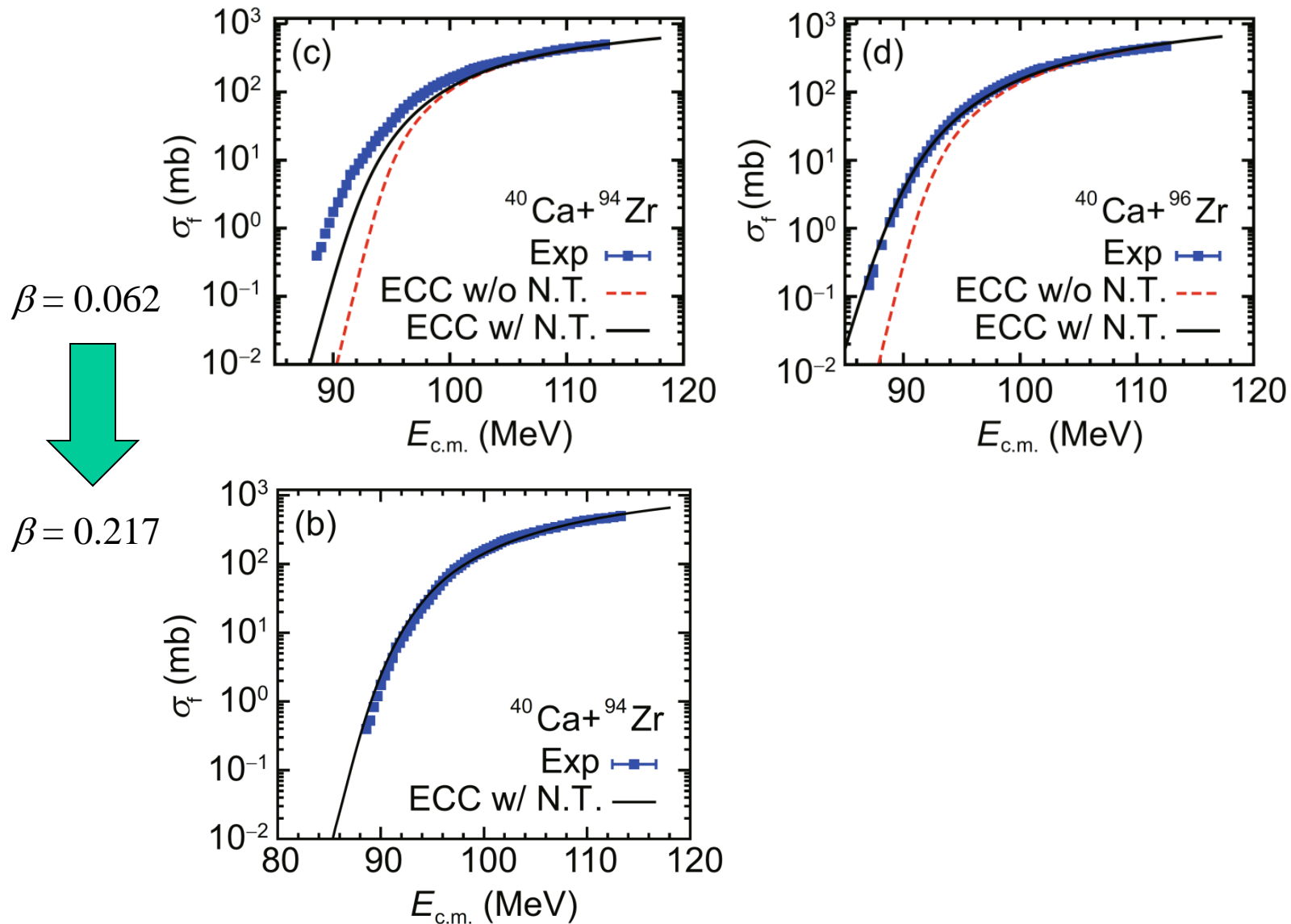
Wang, Wen, Zhao, Zhao & SGZ
At. Data & Nucl. Data Tables 114 (2017) 281

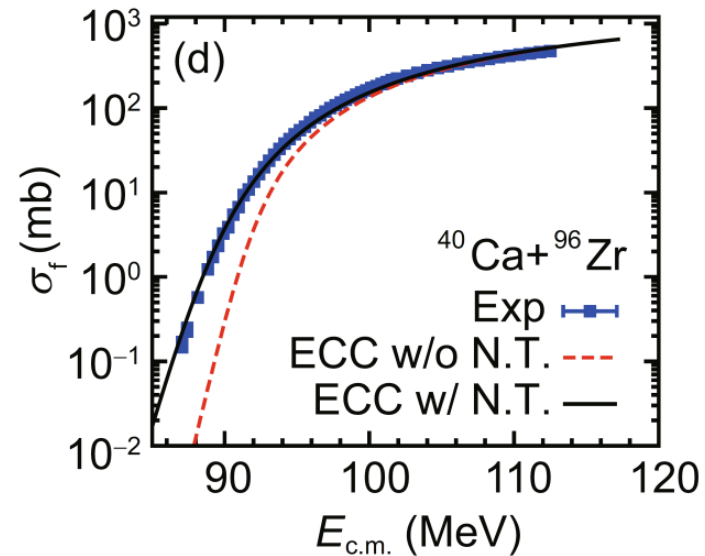
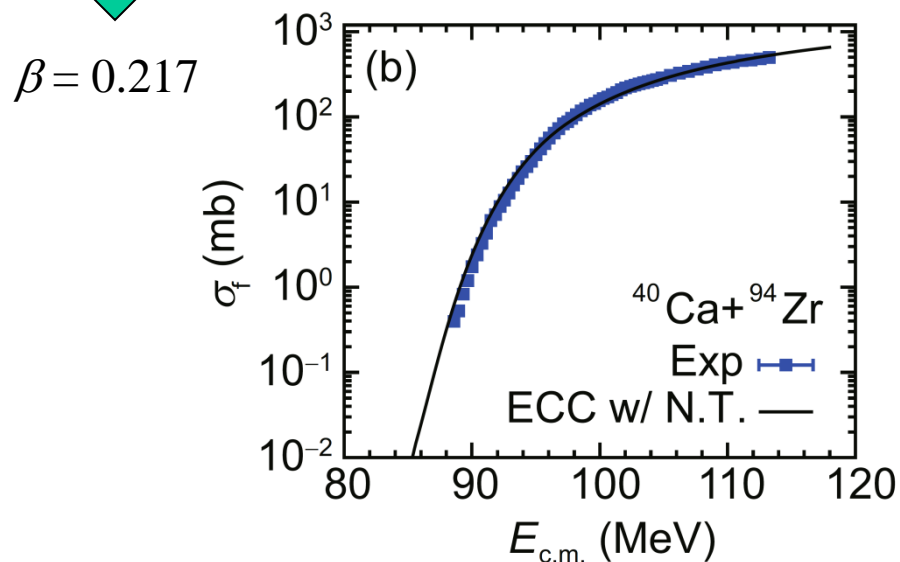
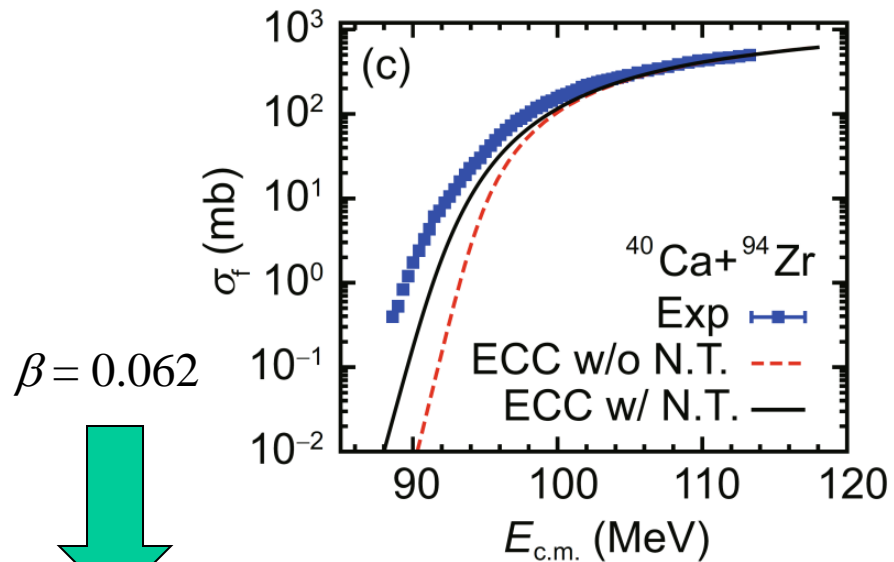




$^{40}\text{Ca}+^{94,96}\text{Zr}$

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Sci. China-Phys. Mech. & Astron. 59 (2016) 642002

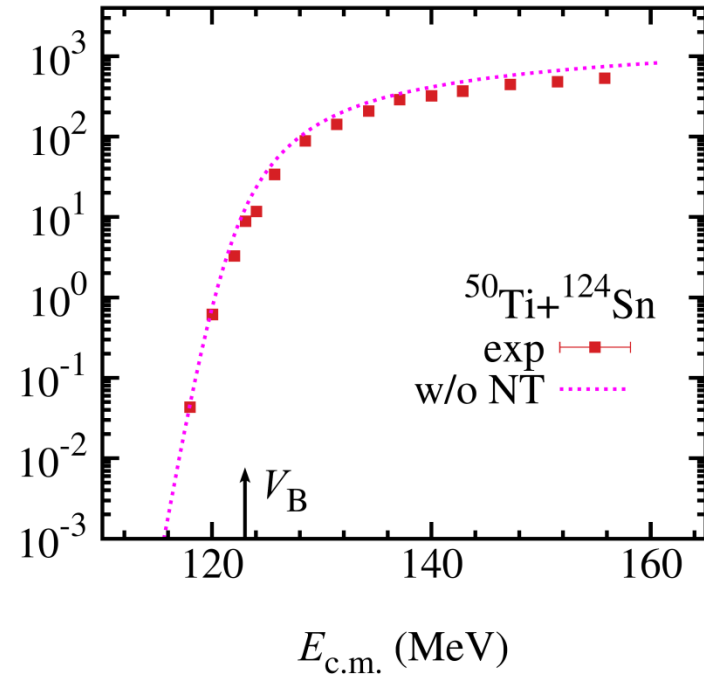
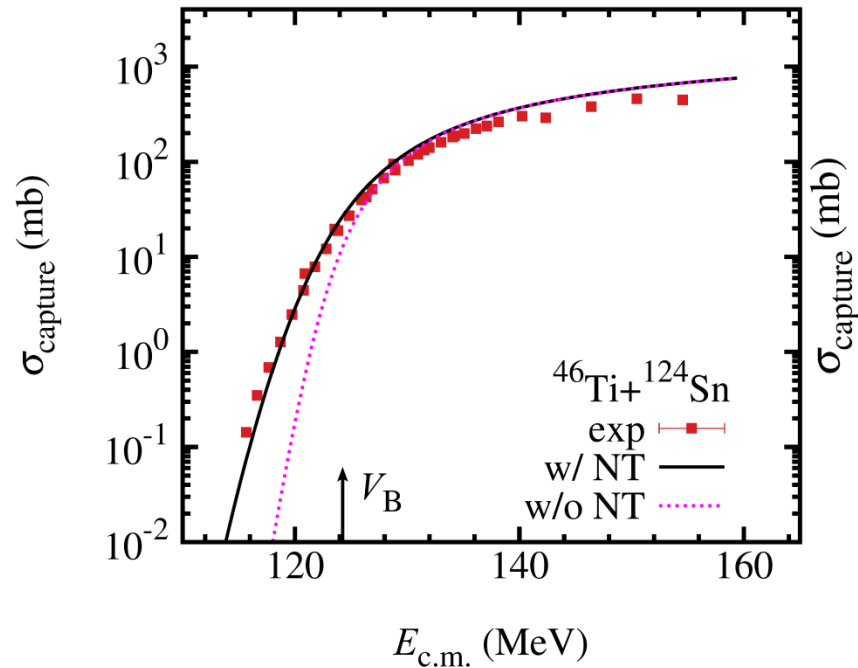




Problems of this model

- (Too) simple
- Depth of nucleus-nucleus potl. fixed
- Higher-order def.: β_3, β_4, \dots
- Mainly around Coulomb barrier
- ...

Predictions: $^{46,50}\text{Ti}+^{124}\text{Sn}$



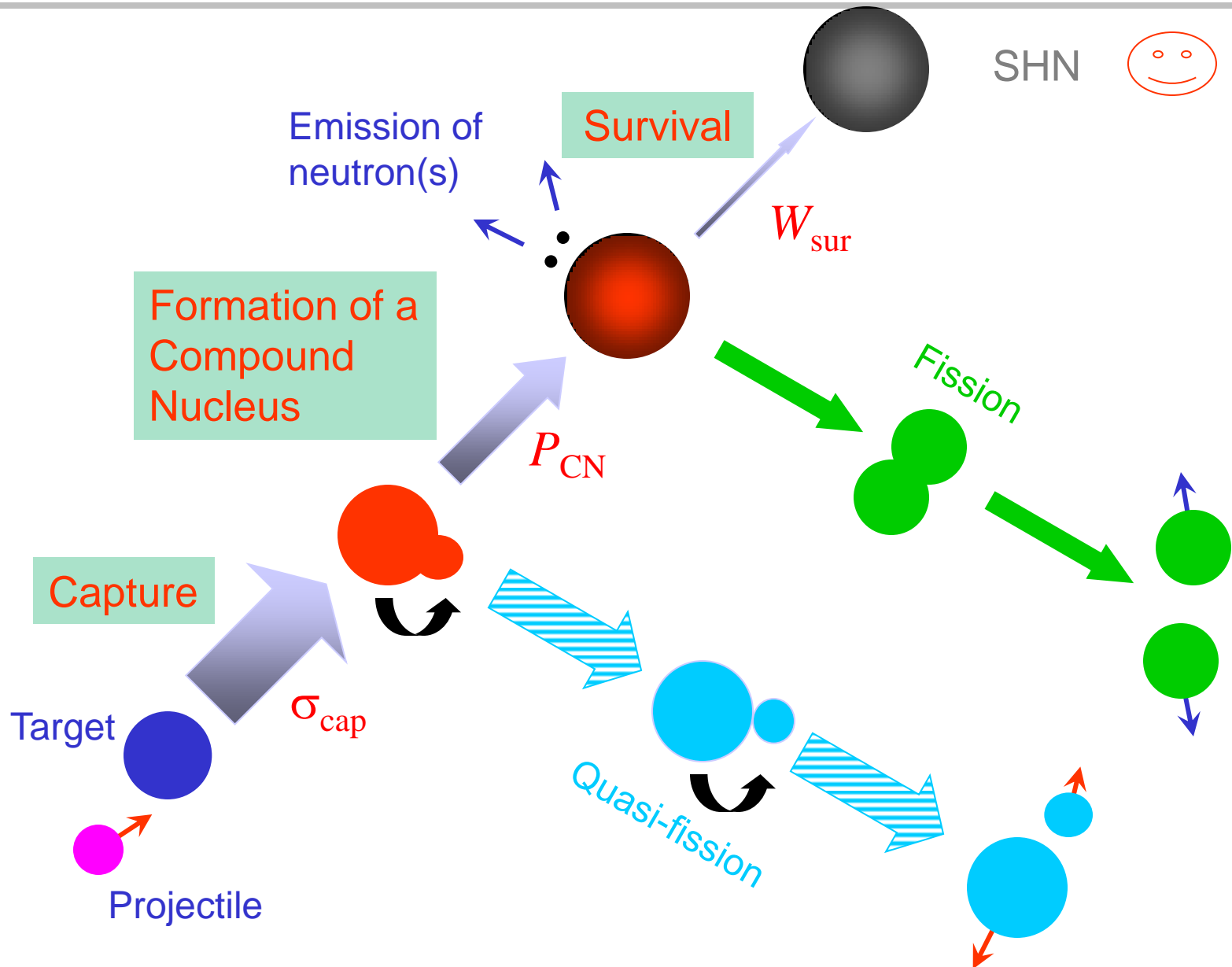
Expt: Liang+2016_PRC94-024616

Lectures 3 & 4

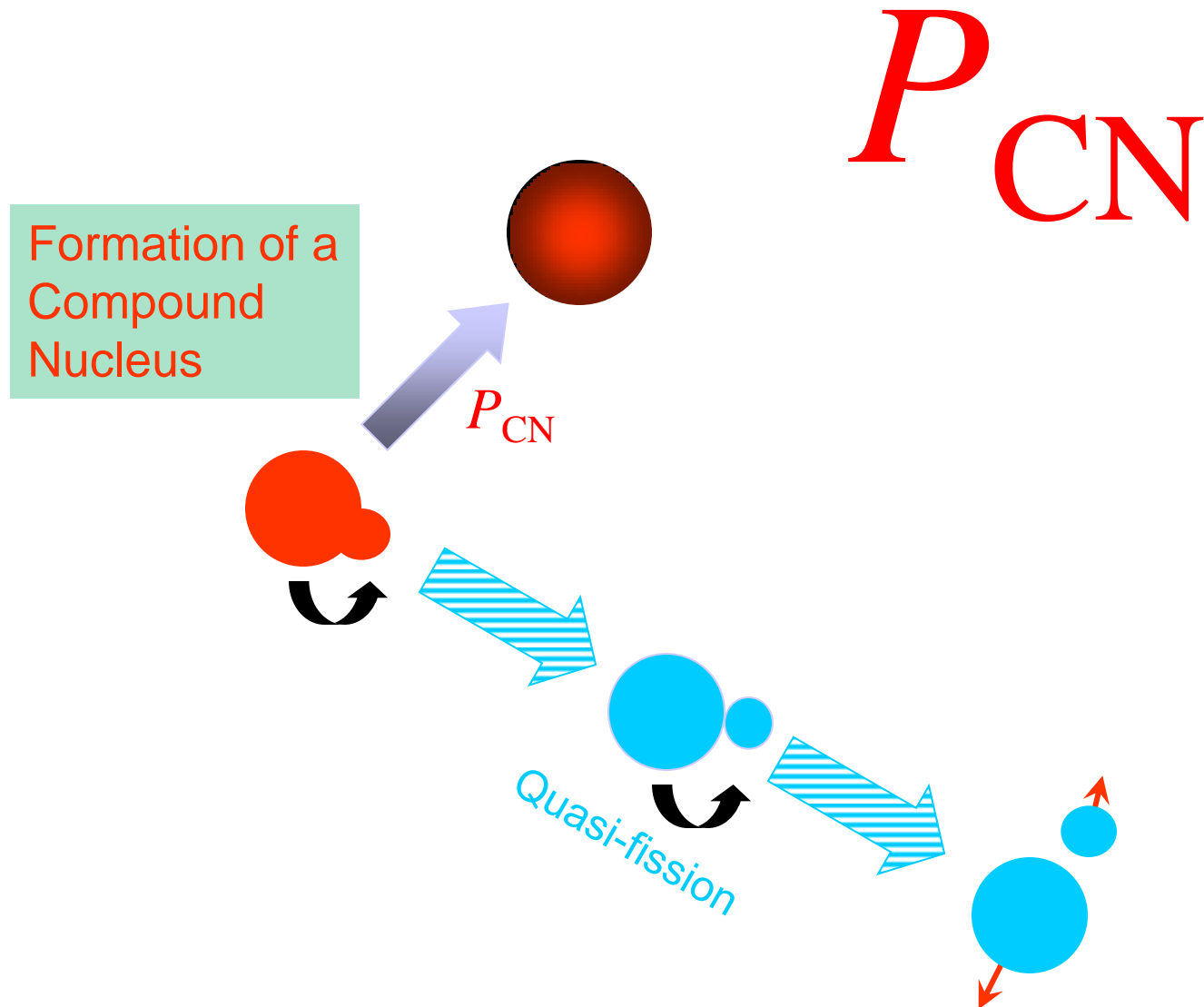
- Challenges in synthesizing SHN

- Synthesis mechanism of SHN
 - Large uncertainties in predicted Xsections
 - Heavy ion fusion reactions
 - Capture
 - Fusion
 - Survival against fission
 - Multi-nucleon transfer reactions
 - Reactions w/ radioactive ion beams

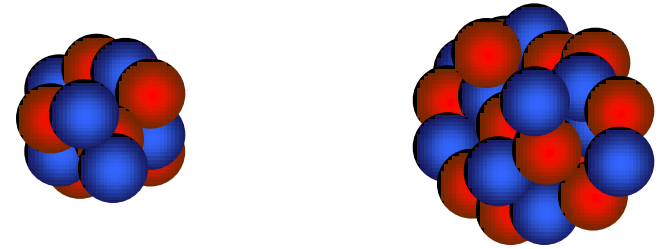
Three steps to a SHN via heavy-ion EvR reaction



Three steps to a SHN via heavy-ion EvR reaction



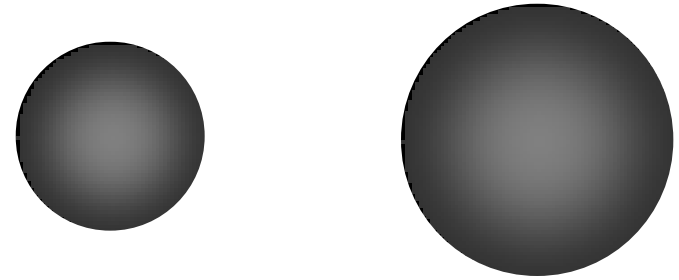
Fusion of two many-body systems



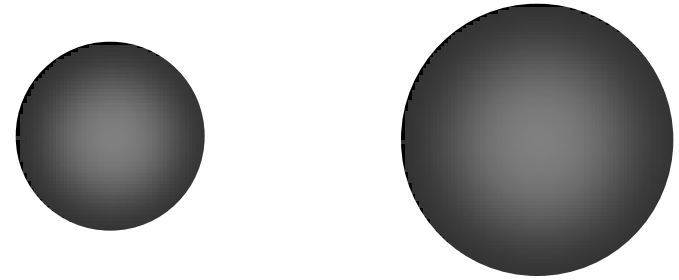
Fusion of two charged liquid drops



Fusion of two charged liquid drops



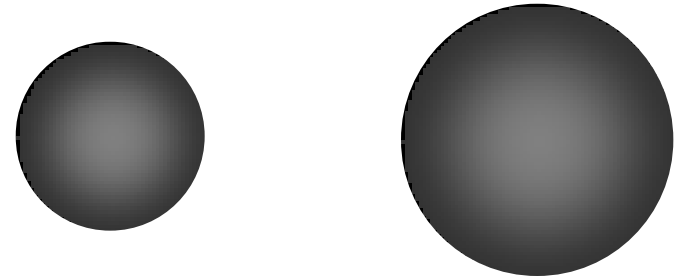
Fusion of two charged liquid drops



Mac. DoFs

- Relative distance
- Mass & charge
- Shape & orientation
- Neck formation
- Transfer of nucleons

Fusion of two charged liquid drops



Mac. DoFs

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- Mass & charge
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Fusion of two charged liquid drops

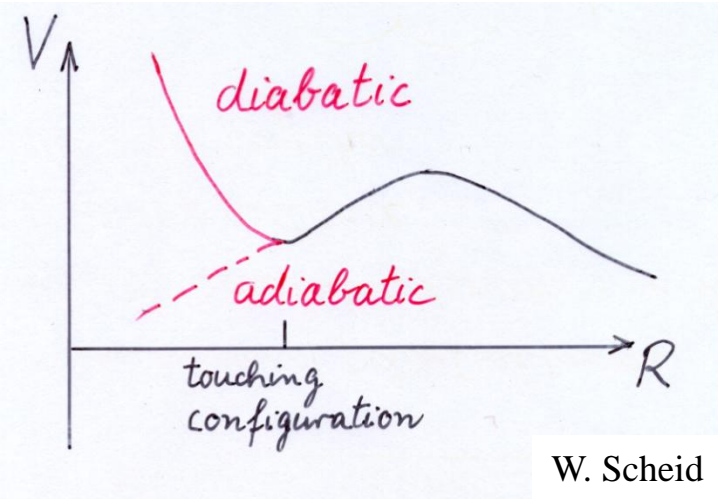
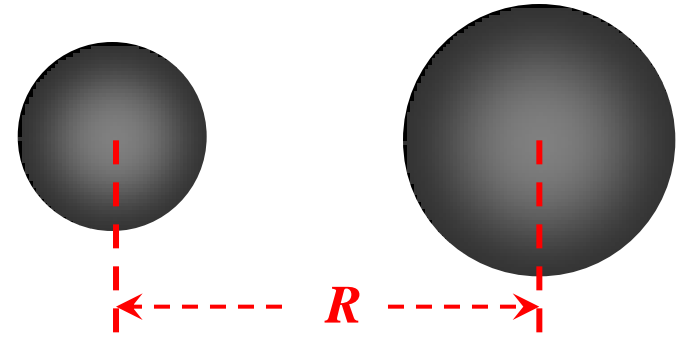
□ Langevin dynamics

$$\left\{ \begin{array}{l} \frac{du(t)}{dt} = - \int_{-\infty}^t \gamma(t-t')u(t')dt' + \frac{1}{\mu}\delta F(t) - \frac{1}{\mu} \frac{dV(R)}{dR} \\ u(t) = \frac{dR(t)}{dt} \end{array} \right.$$

- $R(t)$: Rel. distance
- $u(t)$: Rel. velocity
- $V(R)$: Interaction potl.
- $dF(t)$: Random force
- $g(t-t')$: Friction force

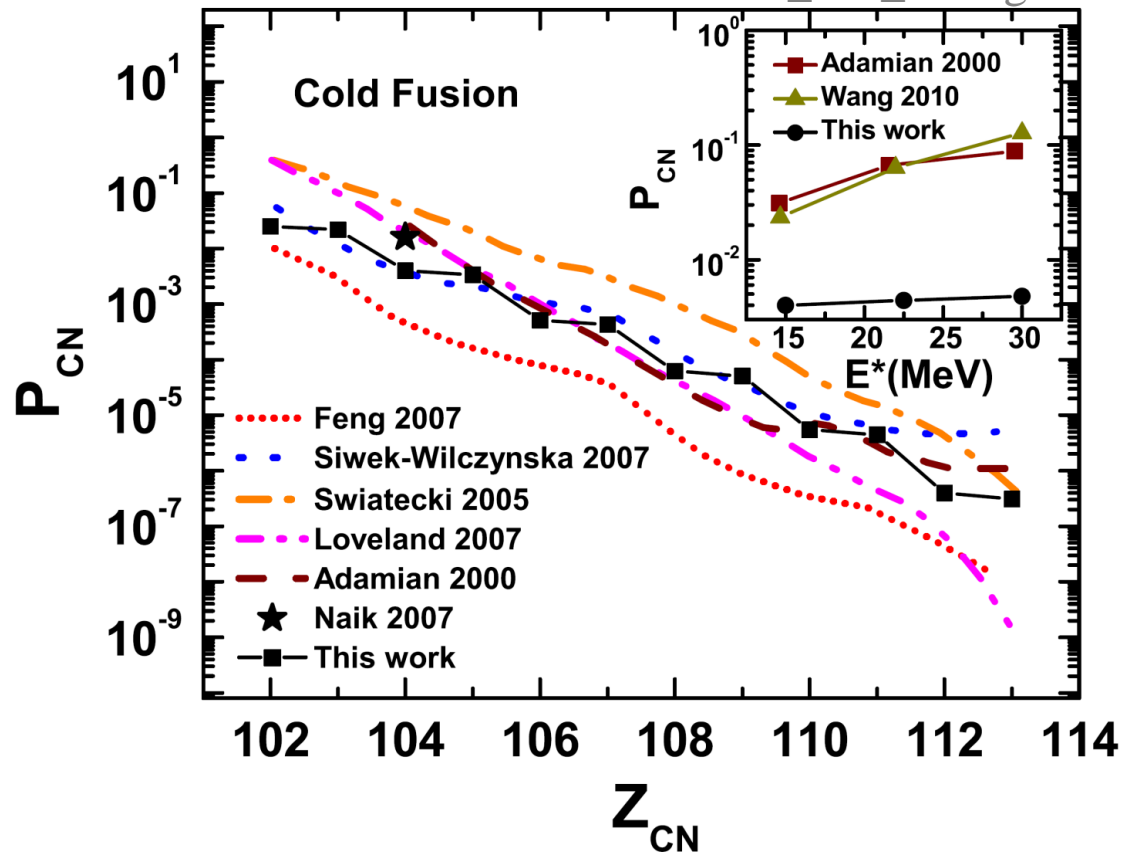
□ Dinuclear system (DNS) model

- Projectile & target keep staying in the potl. Pocket and individuality
- Transfer of nucleons betw. Projectile & target may lead to fusion



Large uncertainties in P_{CN} from various models

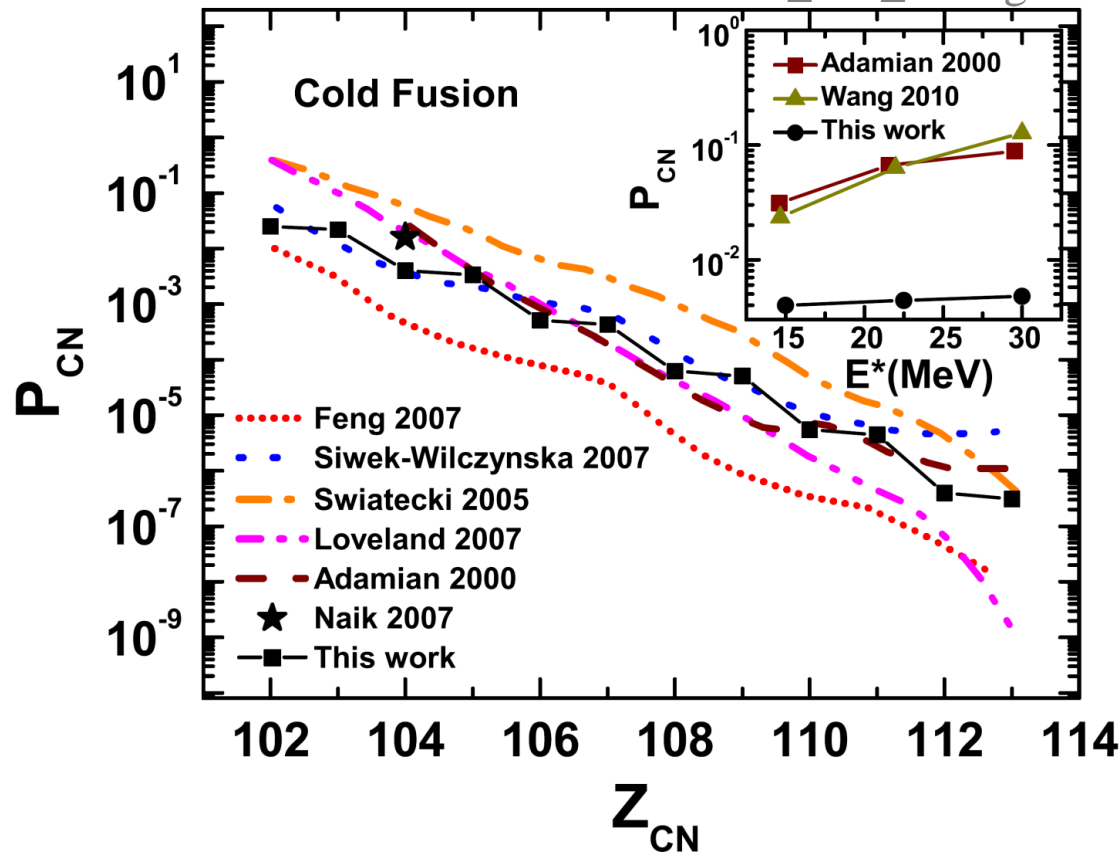
Zhu_Xie_Zhang2014_PRC89-024615



Large uncertainties in P_{CN} from various models

- Experimental study? Systematics?
- Microscopic study of fusion mechanism

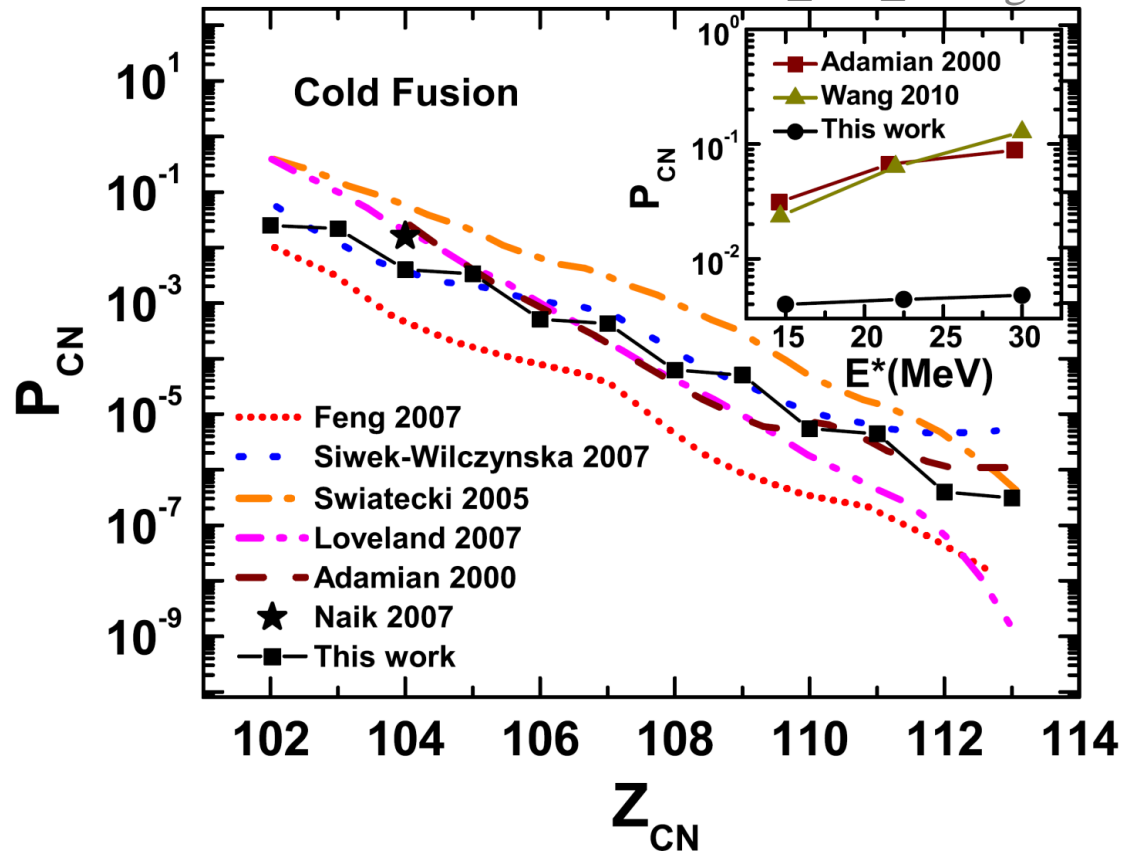
Zhu_Xie_Zhang2014_PRC89-024615



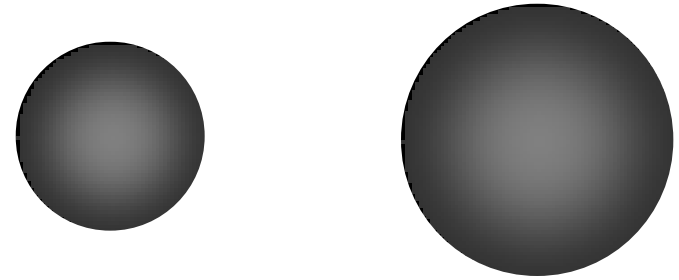
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Zhu_Xie_Zhang2014_PRC89-024615



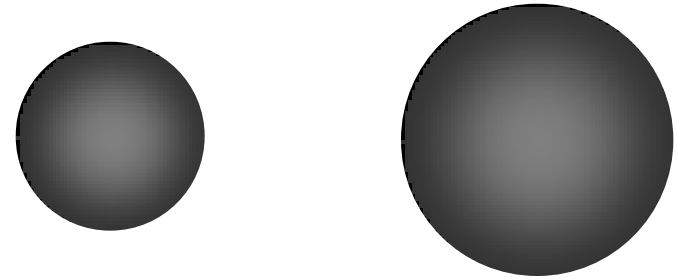
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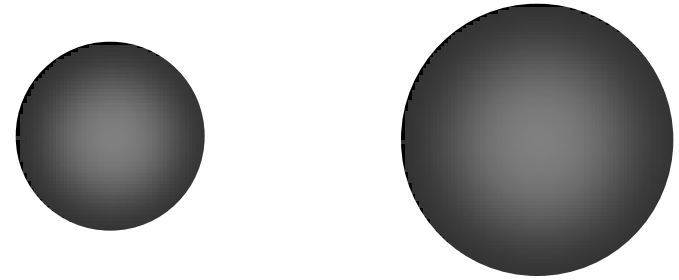
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Mac. DoFs

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-

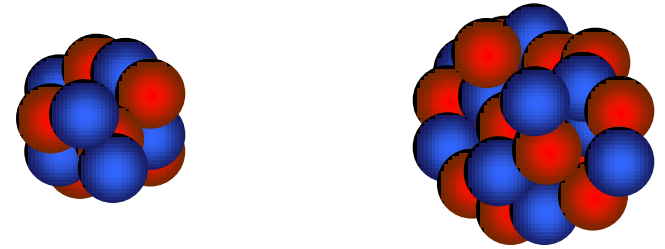
Fusion of two charged liquid drops



Fusion of two charged liquid drops



Fusion of two many-body systems



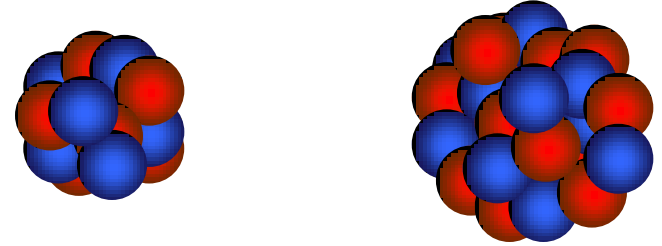
Fusion of two many-body systems

Time dependent Hartree-Fock theory

$$\mathcal{S} = \int_{t_1}^{t_2} dt \langle \Phi(t) | H - i\hbar \partial_t | \Phi(t) \rangle$$

$$H = \sum_i^A t_i + \sum_{i<j}^A v_{ij}$$

$$\Phi(r_1, r_2, \dots, r_A, t) = \frac{1}{\sqrt{A!}} \det |\phi_\lambda(r_i, t)|$$



$$i\hbar \frac{\partial \phi_\lambda}{\partial t} = h \phi_\lambda$$

□ Advantages

- Microscopic
- Successful in structure & reaction

□ Disadvantages

- Only one-body dissipation (collisions with “wall” of mean field)

Negele 1982 Rev. Mod. Phys. 54-913

Guo...2007 Phys. Rev. C76-014601

Guo...2008 Phys. Rev. C77-041301(R)

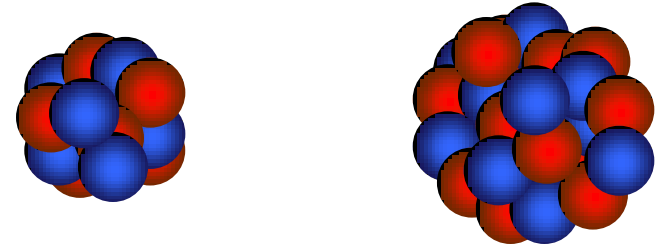
...

Simenel 2012 Euro. Phys. J. A 48-152

Fusion of two many-body systems

Quantum Molecular Dynamics (QMD)

$$\phi_i(\mathbf{r}) = \frac{1}{(2\pi\sigma_r^2)^{3/4}} \exp \left[-\frac{(\mathbf{r} - \mathbf{r}_i)^2}{4\sigma_r^2} + \frac{i}{\hbar} \mathbf{r} \cdot \mathbf{p}_i \right]$$
$$\dot{\mathbf{r}}_i = \frac{\partial H}{\partial \mathbf{p}_i}, \quad \dot{\mathbf{p}}_i = -\frac{\partial H}{\partial \mathbf{r}_i}$$



Aichelin 1991 Phys. Rep. 202-233

...

Wang_Li_Wu 2002 PRC65-064608

Wang...2004 PRC69-034608

□ Advantages

- Microscopic
- Both mean field & collision terms included

□ More attentions should be paid on

- Pauli exclusive principle
- Shell effects

ImQMD & TDHF: DIC & multi-nucleon transfer leading to n-rich isotopes

Physics Letters B 760 (2016) 236–241

Wang&Guo2016_PLB760-236

Contents lists available at [ScienceDirect](http://www.sciencedirect.com)

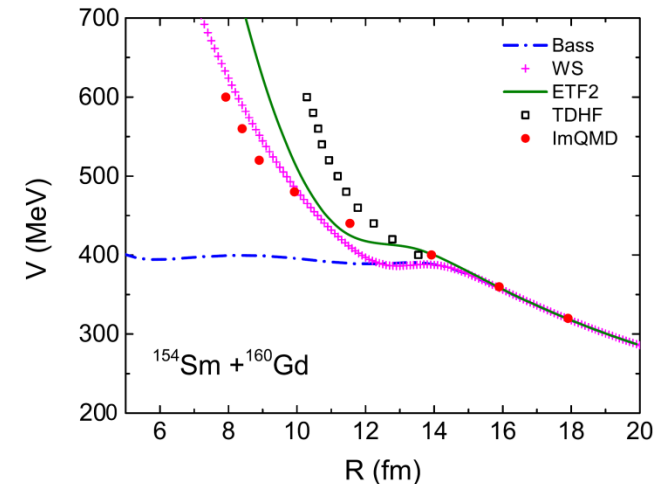
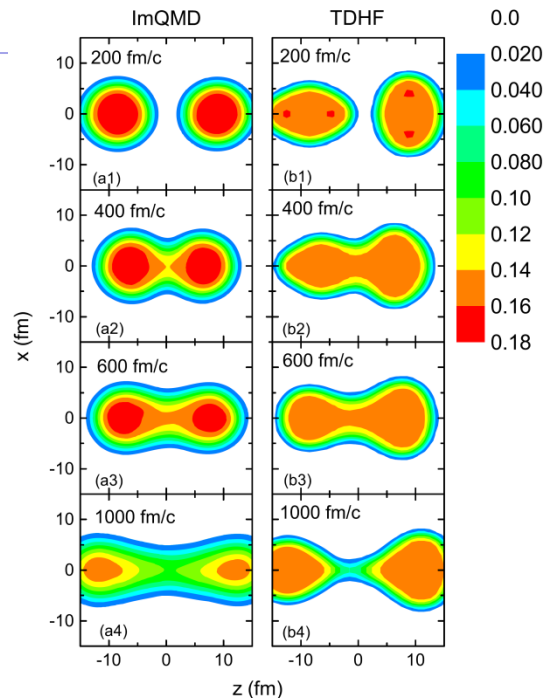
Physics Letters B

www.elsevier.com/locate/physletb



New neutron-rich isotope production in $^{154}\text{Sm} + ^{160}\text{Gd}$

Ning Wang^{a,*}, Lu Guo^{b,c}



TDHF: quasifission dynamics

SCIENCE CHINA
Physics, Mechanics & Astronomy

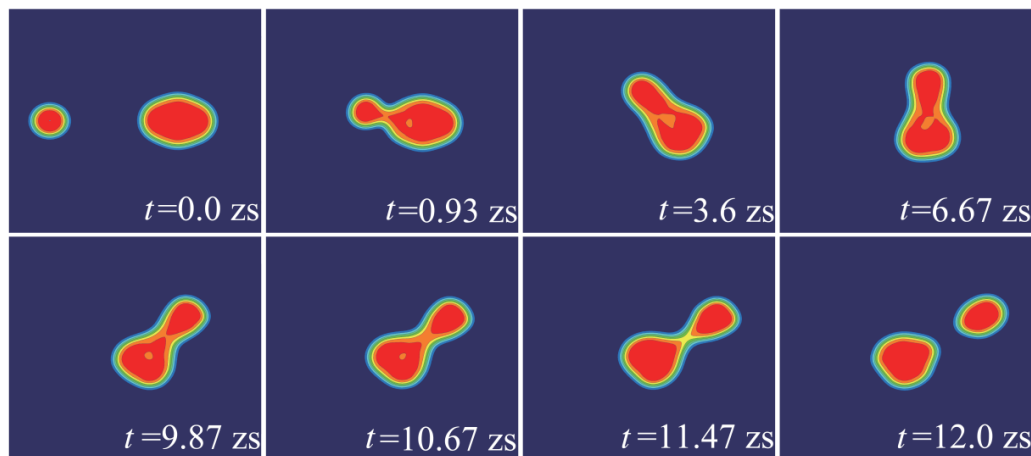


• Article •

September 2017 Vol. 60 No. 9: 092011
doi: [10.1007/s11433-017-9063-3](https://doi.org/10.1007/s11433-017-9063-3)

Angular momentum dependence of quasifission dynamics in the reaction $^{48}\text{Ca} + ^{244}\text{Pu}$

Chong Yu, and Lu Guo*



- ❑ Q.f. dynamics is sensitive to the angular momentum
- ❑ Contact time of q.f. decreases w/ the angular momentum
- ❑ Q.f. is accompanied by MNT

TDHF: fusion dynamics

Physics Letters B 782 (2018) 401–405

Contents lists available at [ScienceDirect](#)



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Physics Letters B

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The role of tensor force in heavy-ion fusion dynamics

Lu Guo^{a,*}, Cédric Simenel^b, Long Shi^a, Chong Yu^a

^a School of Physics, University of Chinese Academy of Sciences, Beijing 100049, China

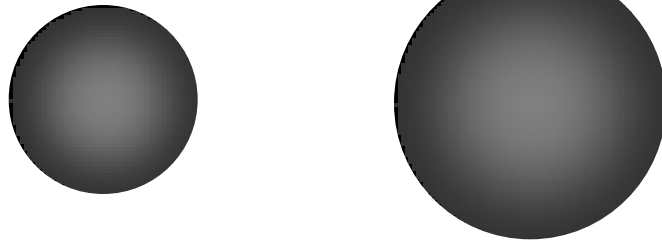
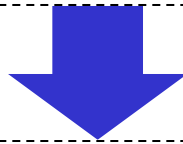
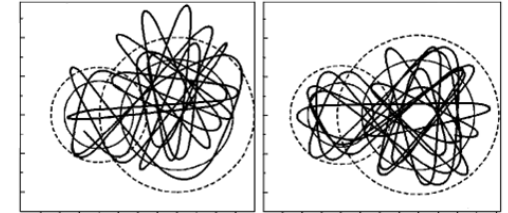
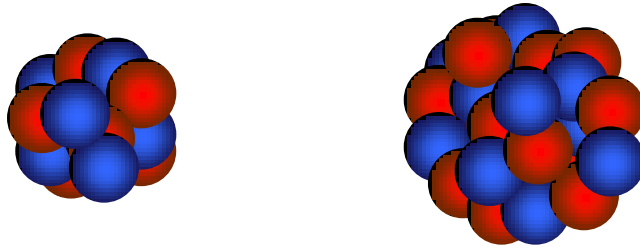
^b Department of Nuclear Physics, RSPE, Australian National University, Canberra, Australia



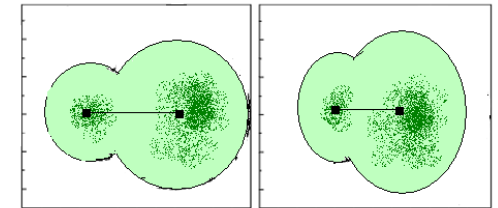
- ❑ Notable effects on fusion barriers & Xsections are observed by the inclusion of tensor force
- ❑ The effects are essentially attributed to the shift of low-lying vibration states of colliding partners & nucleon transfer in the asymmetric reactions

Fusion: Bridging microscopic & macroscopic description

Improved QMD simulations

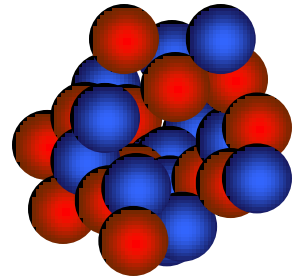
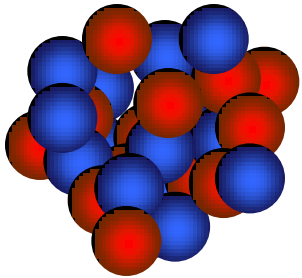


$$\begin{cases} \frac{du(t)}{dt} = - \int_{-\infty}^t \gamma(t-t')u(t')dt' + \frac{1}{\mu}\delta F(t) - \frac{1}{\mu} \frac{dV(R)}{dR} \\ u(t) = \frac{dR(t)}{dt} \end{cases}$$



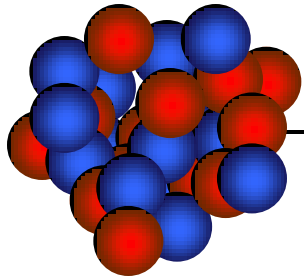
ImQMD simulations

□ $^{90}\text{Zr}+^{90}\text{Zr}$: Monte Carlo sampling; 10,000 events



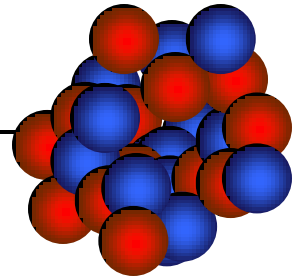
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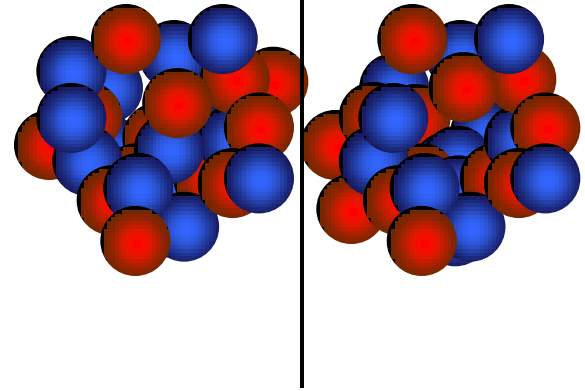
Central collision, $E_{\text{cm}} = 195 \text{ MeV}$

$R_{\infty} = 30 \text{ fm}$



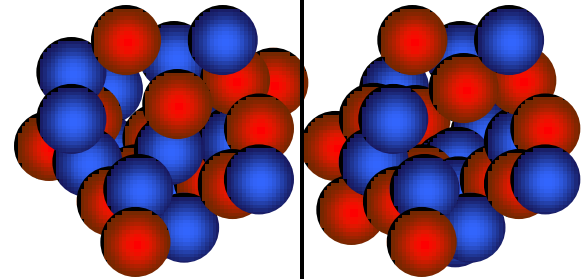
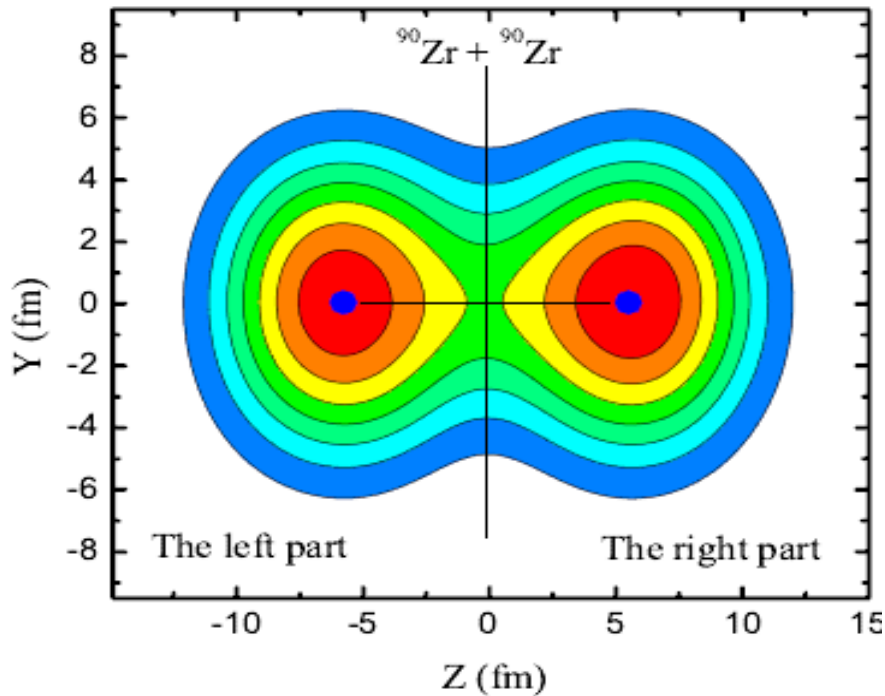
ImQMD simulations

□ $^{90}\text{Zr}+^{90}\text{Zr}$: $b=0$, $E_{\text{c.m.}}=195$ MeV, $R_{\infty}=30$ fm, 10,000 events



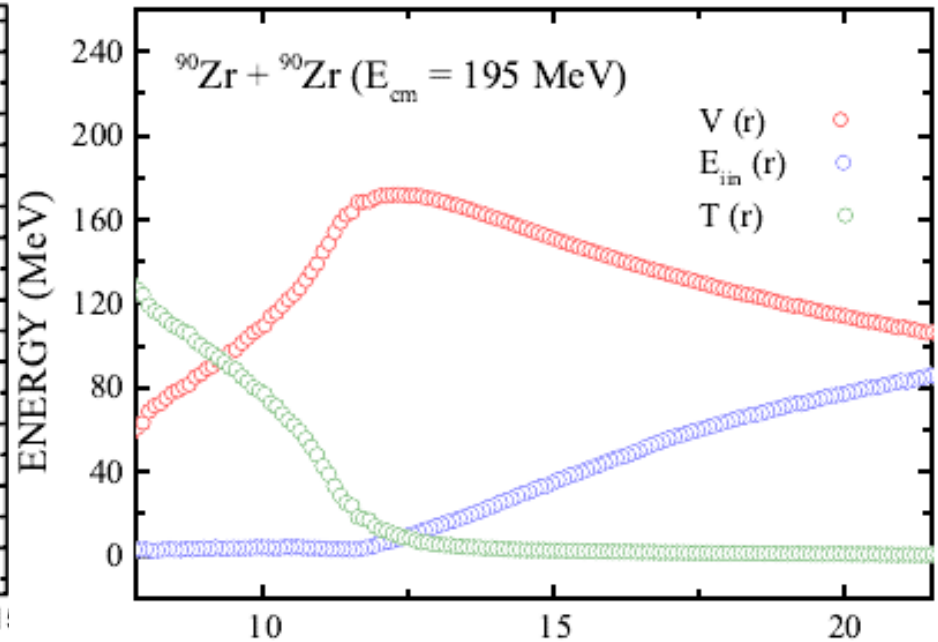
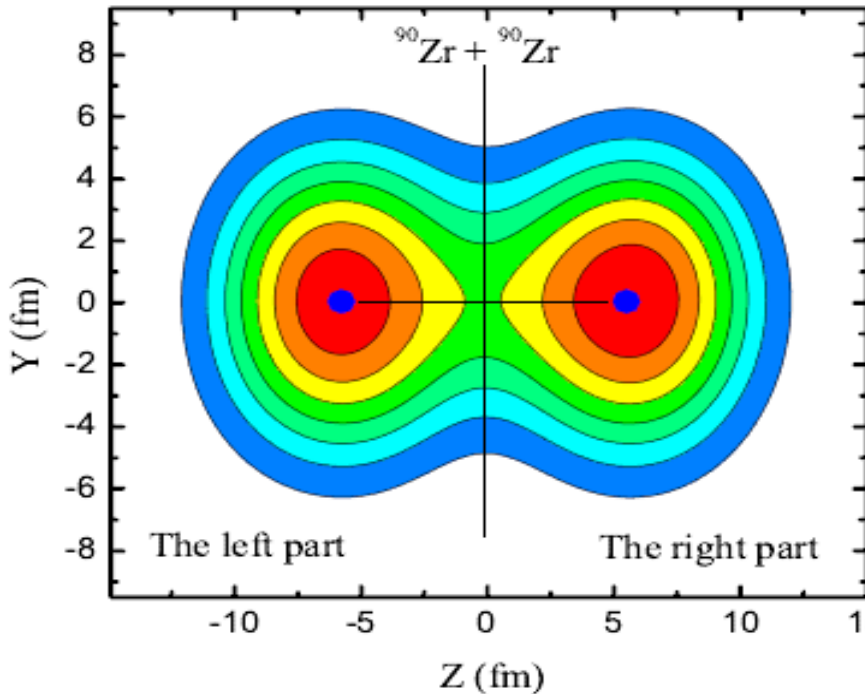
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ImQMD simulations

□ $^{90}\text{Zr} + ^{90}\text{Zr}$: $b=0$, $E_{\text{c.m.}}=195$ MeV, $R_{\infty}=30$ fm, 10,000 events



Potential:

$$V(R) = E_{\text{tot}}(R) - E_{\text{left}}(R) - E_{\text{right}}(R),$$

Collective Energy:

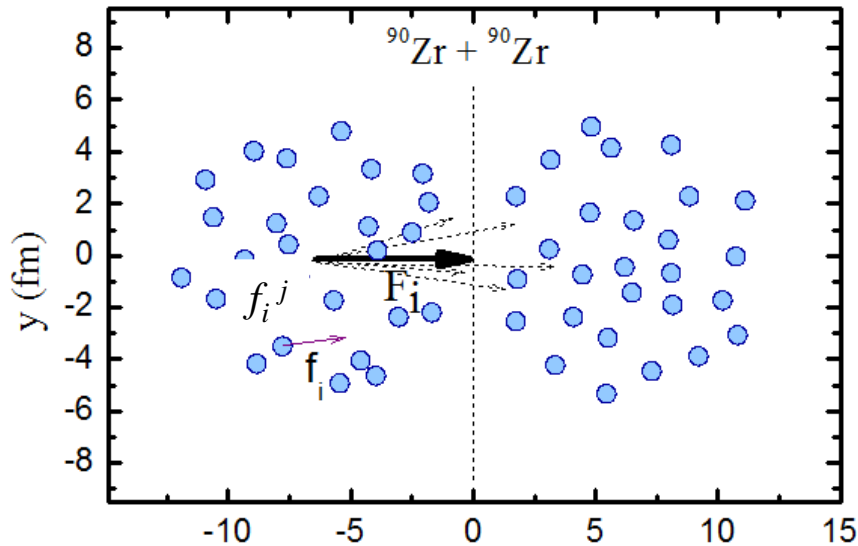
$$E_{\text{coll}}(R) = P^2/2\mu + V(R).$$

Intrinsic Energy:

$$E_{\text{intr}}(R) \equiv E_{\text{tot}}(R) - E_{\text{coll}}(R)$$

Macroscopic reduction of random force

□ Fluctuations from initialization & collisions



For the i -th event: $F_i(x) \equiv \sum_{j=1}^A f_i^j(x)$

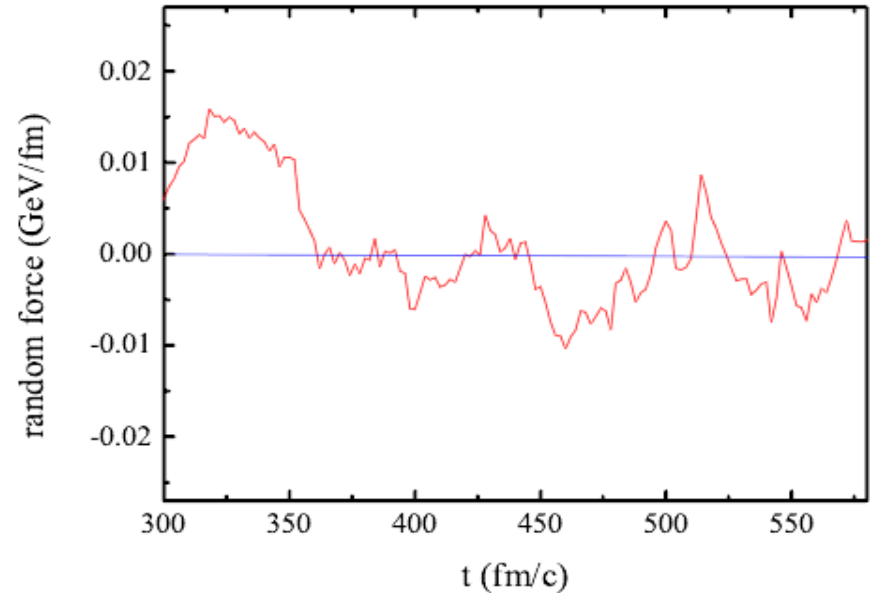
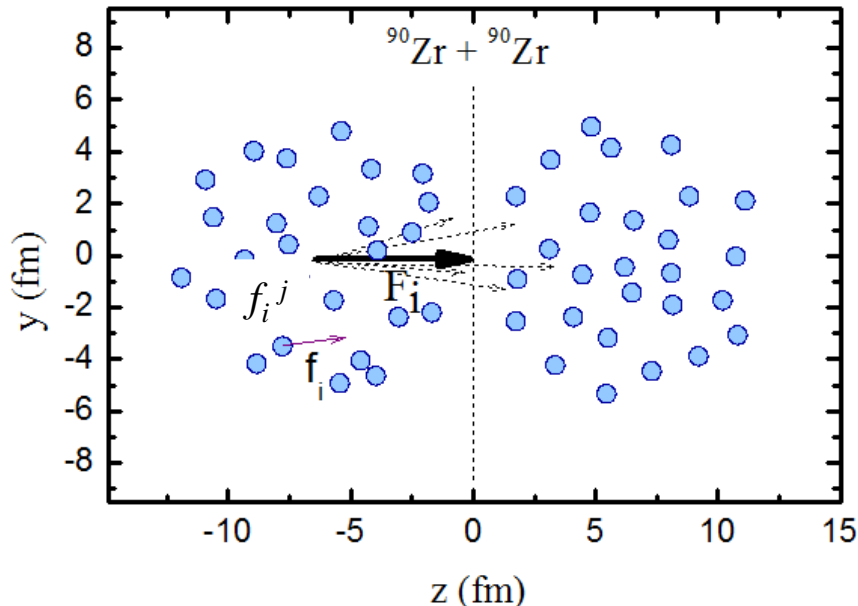
On average: $\langle F(x) \rangle \equiv \frac{1}{n} \sum_{i=1}^n F_i(x)$

Random force:

$$\delta F(x)_i \equiv F_i(x) - \langle F(x) \rangle, \quad x = t \text{ or } R,$$

Macroscopic reduction of random force

Fluctuations from initialization & collisions



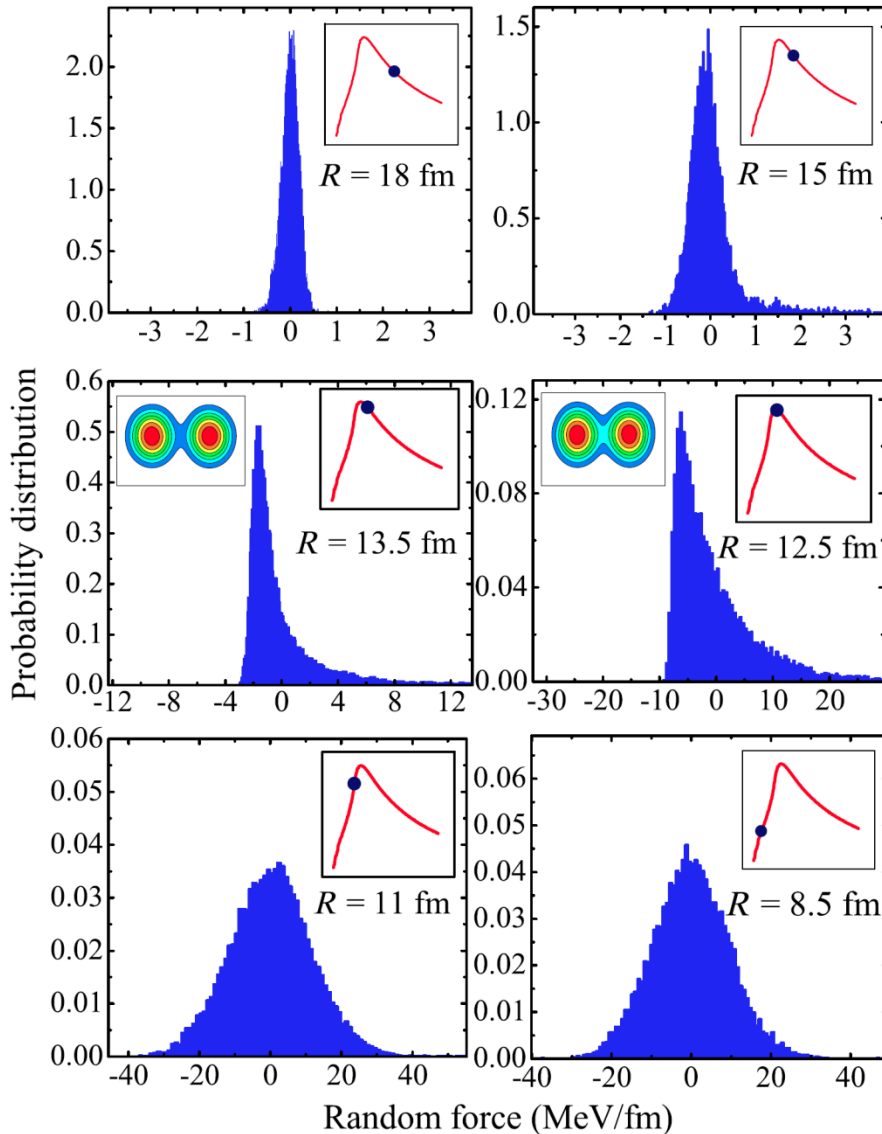
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Random force:

$$\delta F(x)_i \equiv F_i(x) - \langle F(x) \rangle, \quad x = t \text{ or } R,$$

Distribution of random force



Approaching stage

Touching & fusion stage

Post-fusion stage

Fusion: Bridging microscopic & macroscopic description

- ❑ Macroscopic parameters, including the random force and the friction coefficient, characterizing the Langevin type description of the nuclear fusion are extracted from the microscopic QMD dynamics around the Coulomb barrier
- ❑ The dissipation dynamics of the relative motion between two fusing nuclei is associated with non-Gaussian distributions of the random force
- ❑ A proper treatment of the non-Markovian (memory) effects in the Langevin dynamics is crucial for the dynamics of emergence in the nuclear dissipative fusion motion

[Wen_Sakata_Li_Wu_Zhang_SGZ 2013 Phys. Rev. Lett. 111, 012501](#)

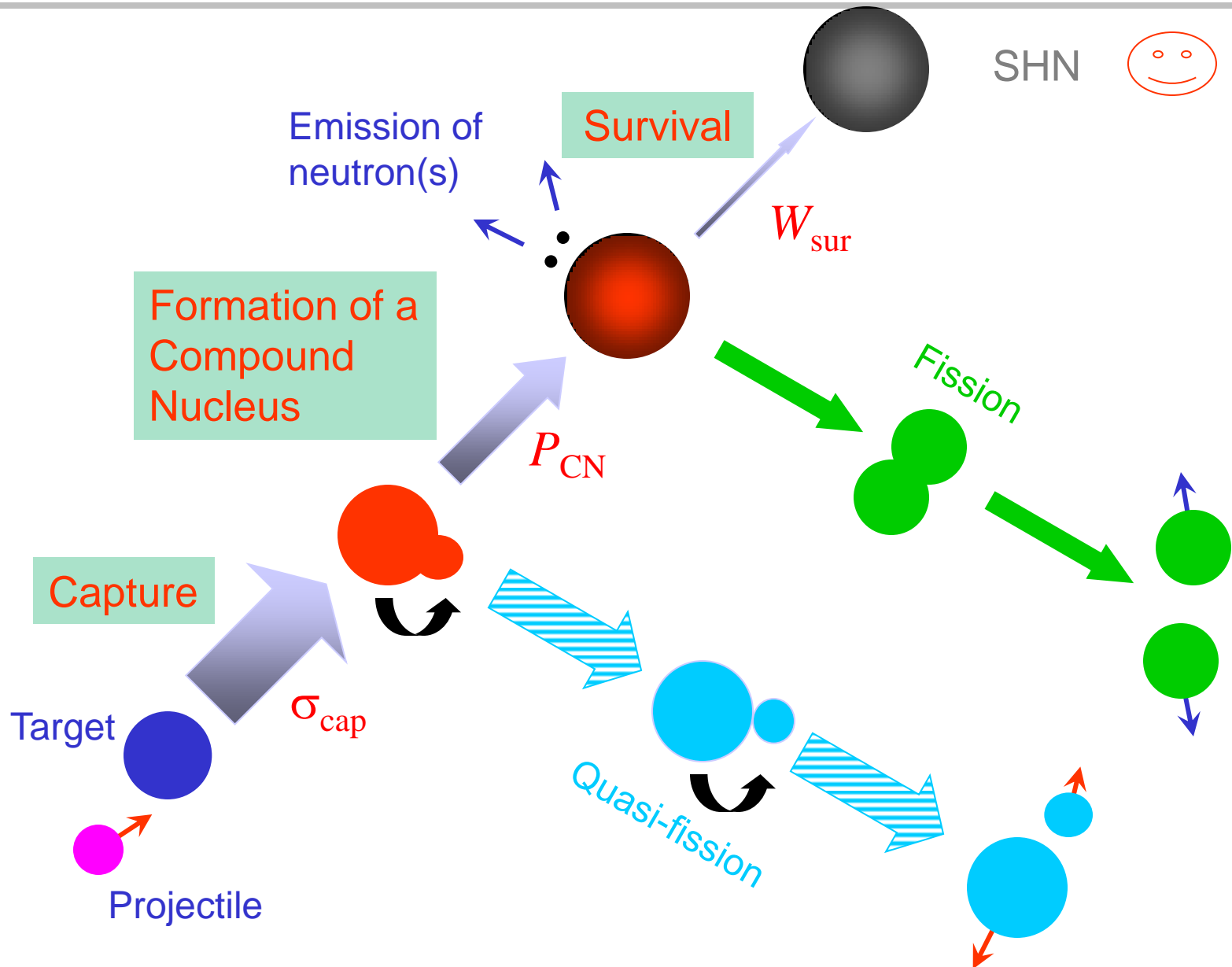
[Wen_Sakata_Li_Wu_Zhang_SGZ 2014 Phys. Rev. C90, 054613](#)

Lectures 3 & 4

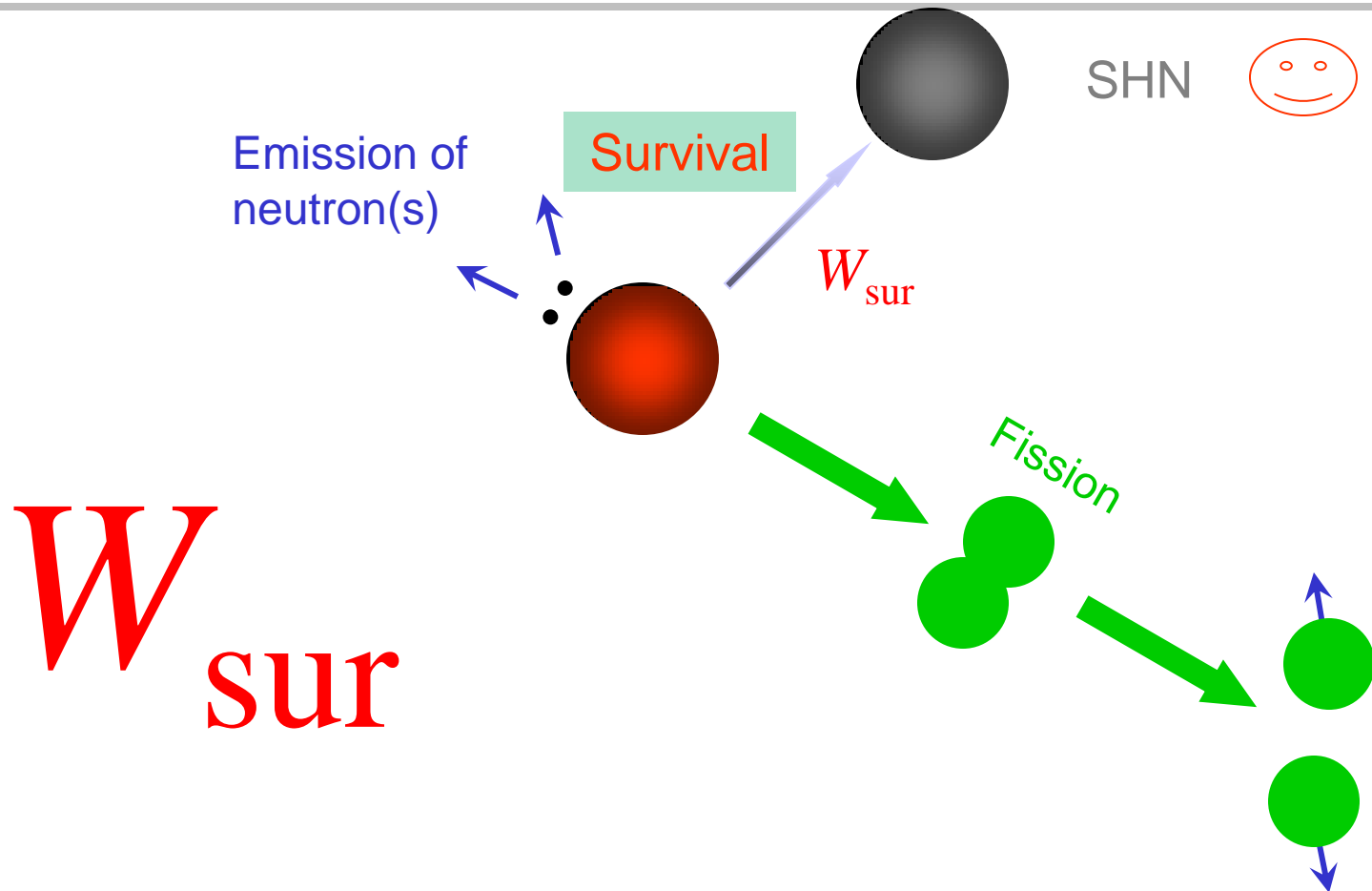
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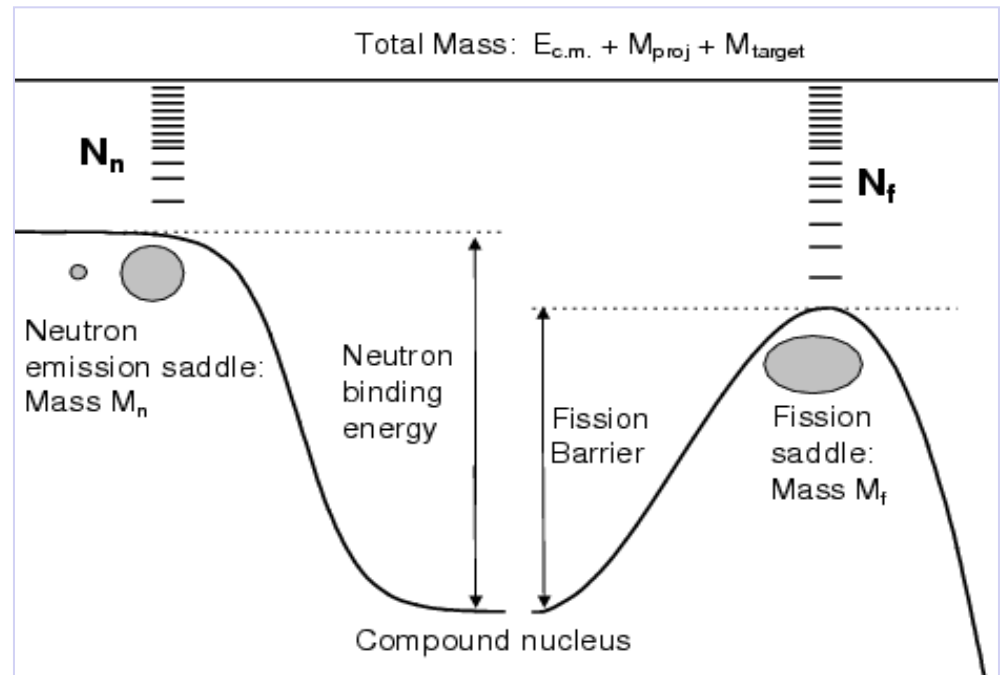
Fission of heavy & superheavy nuclei

De-excitation of excited CN

- Fission
- Neutron emission
- ...

Parameters needed

- Fission barrier
- Neutron separation energy
- ...



Swiatecki_Siwiek-Wilczynska_Wilczynski2005_PRC71-014602

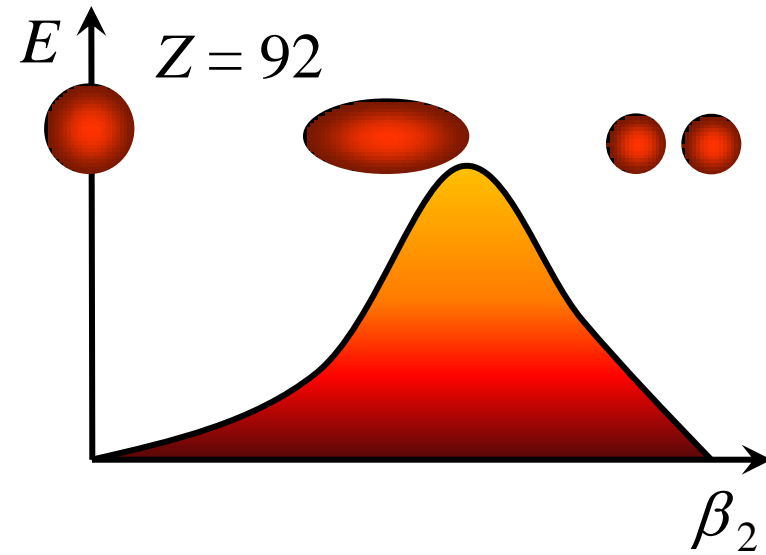
Theoretical approaches for calculating fission barriers:

- Macroscopic-microscopic approaches
- Self-consistent approaches

Multi-dim. constrained relativistic mean field models
($\beta_{20}, \beta_{22}, \beta_{30}, \dots$)

Nuclear fission

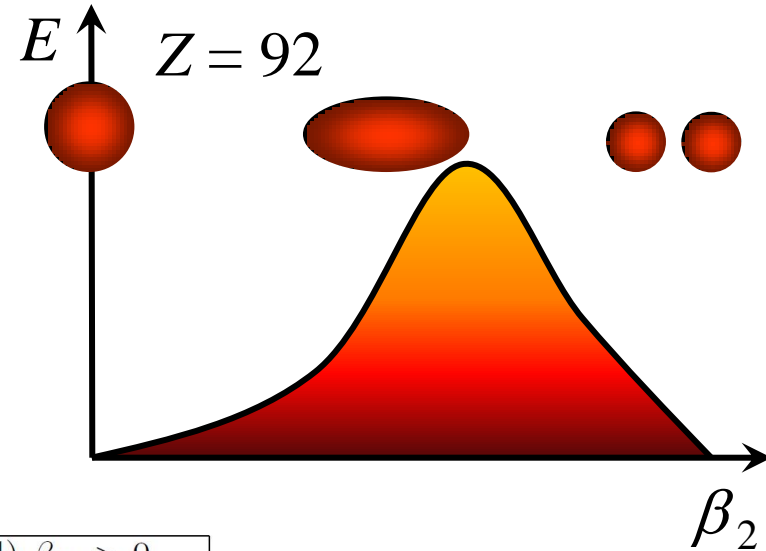
- Fission barrier is crucial for the description of fission

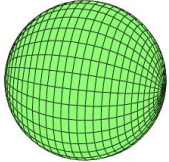
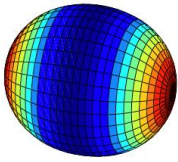
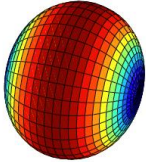
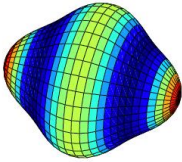
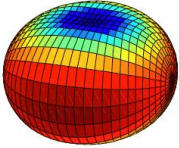
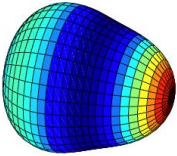
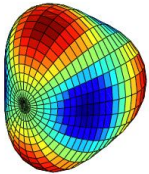
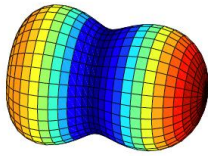


Nuclear fission

- ❑ Fission barrier is crucial for the description of fission
- ❑ Various shapes may appear during fission

$$R(\theta, \varphi) = R_0 \left[1 + \beta_{00} + \sum_{\lambda=1}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \beta_{\lambda\mu}^* Y_{\lambda\mu}(\theta, \varphi) \right]$$

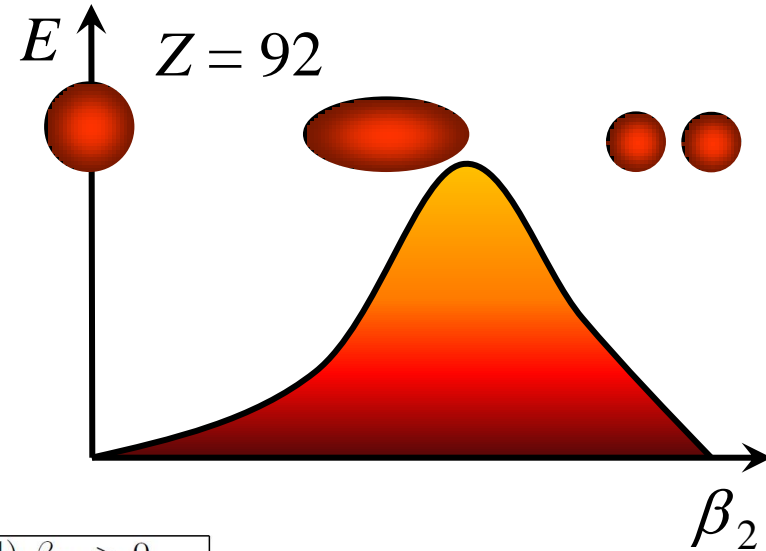


(a) $\beta_{\lambda\mu} = 0$	(b) $\beta_{20} > 0$	(c) $\beta_{20} < 0$	(d) $\beta_{40} > 0$
			
(e) $\beta_{22} \neq 0$	(f) $\beta_{30} \neq 0$	(g) $\beta_{32} \neq 0$	(h) $\beta_{20} \gg 0$
			

Nuclear fission

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$$R(\theta, \varphi) = R_0 \left[1 + \beta_{00} + \sum_{\lambda=1}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \beta_{\lambda\mu}^* Y_{\lambda\mu}(\theta, \varphi) \right]$$



(a) $\beta_{\lambda\mu} = 0$	(b) $\beta_{20} > 0$	(c) $\beta_{20} < 0$	(d) $\beta_{40} > 0$
(e) $\beta_{22} \neq 0$	(f) $\beta_{30} \neq 0$	(g) $\beta_{32} \neq 0$	(h) $\beta_{20} \gg 0$

To include as many shape degrees of freedom as possible

Covariant Density Functional Theory (CDFT)

$$\begin{aligned}
 \mathcal{L} = & \bar{\psi}_i (i\partial - M) \psi_i + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma - U(\sigma) - g_\sigma \bar{\psi}_i \sigma \psi_i \\
 & - \frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - g_\omega \bar{\psi}_i \psi \psi_i \\
 & - \frac{1}{4} \vec{R}_{\mu\nu} \vec{R}^{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \vec{\rho}^\mu - g_\rho \bar{\psi}_i \vec{\rho} \vec{\tau} \psi_i \\
 & - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - e \bar{\psi}_i \frac{1 - \tau_3}{2} A \psi_i,
 \end{aligned}$$

Serot_Walecka1986_ANP16-1

Reinhard1989 RPP52-439

Ring1996_PPNP37-193

Vretenar_Afanasjev_Lalazissis_Ring2005_PR409-101

Meng_Toki_SGZ_Zhang_Long_Geng2006_PPNP57-470

$$(\alpha \cdot \mathbf{p} + \beta(M + S(\mathbf{r})) + V(\mathbf{r})) \psi_i = \epsilon_i \psi_i$$

Liang_Meng_SGZ2015_PR570-1

$$(-\nabla^2 + m_\sigma^2) \sigma = -g_\sigma \rho_S - g_2 \sigma^2 - g_3 \sigma^3$$

Meng_SGZ2015_JPG42-093101

$$(-\nabla^2 + m_\omega^2) \omega = g_\omega \rho_V - c_3 \omega^3$$

$$(-\nabla^2 + m_\rho^2) \rho = g_\rho \rho_3$$

$$-\nabla^2 A = e \rho_C$$

MDC-CDFT ($\beta_{20}, \beta_{22}, \beta_{30}, \beta_{32}, \beta_{40}, \dots$)

ph channel	Non-linear	Density-dependent
Meson exchange	NL3, NL3*, PK1, ...	DD-ME1, DD-ME2, ...
Point Coupling	PC-F1, PC-PK1, ...	DD-PC1, ...

MDC-RMF

MDC-RHB

pp channel	BCS	Bogoliubov
Constant gap	√	
Constant strength	√	
Delta force	√	√
Separable force	√	√

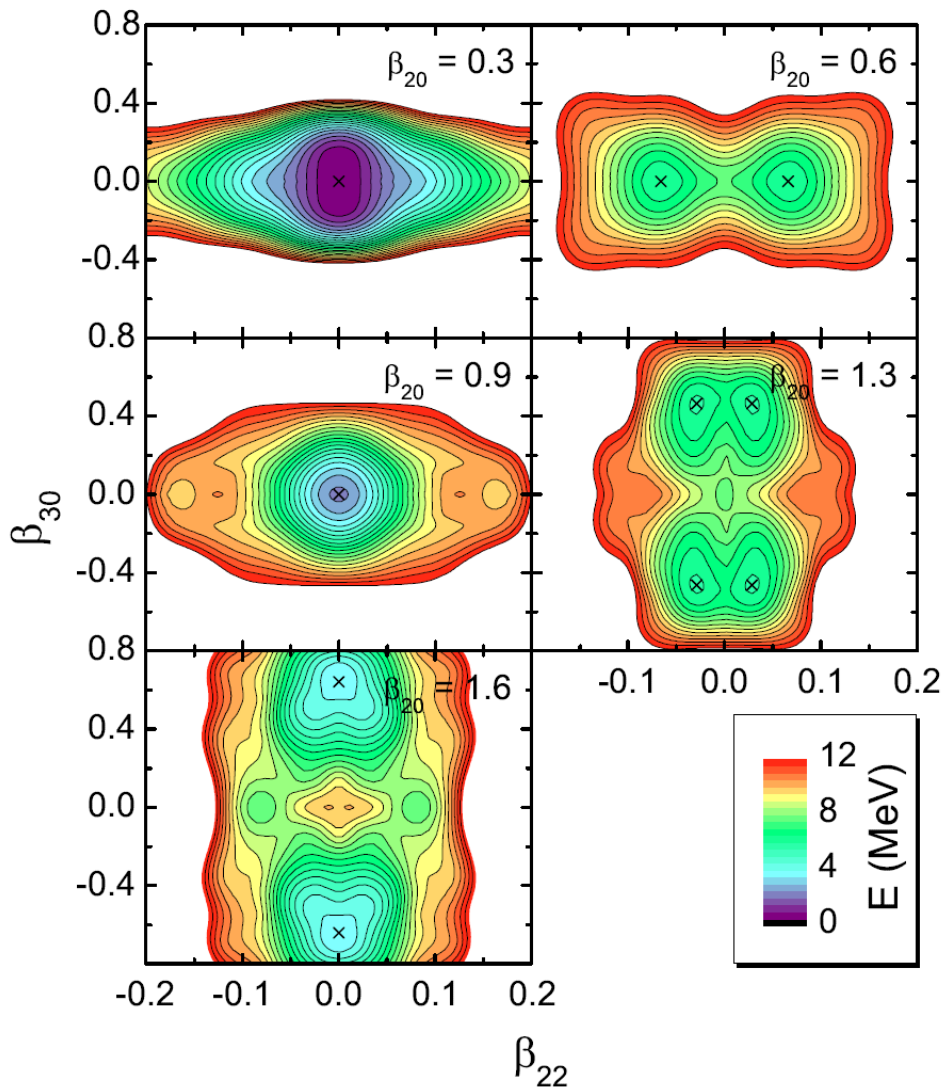
[Lu_Zhao_SGZ 2011_PRC84-014328](#)

[Lu_Zhao_SGZ 2012_PRC85-011301R](#)

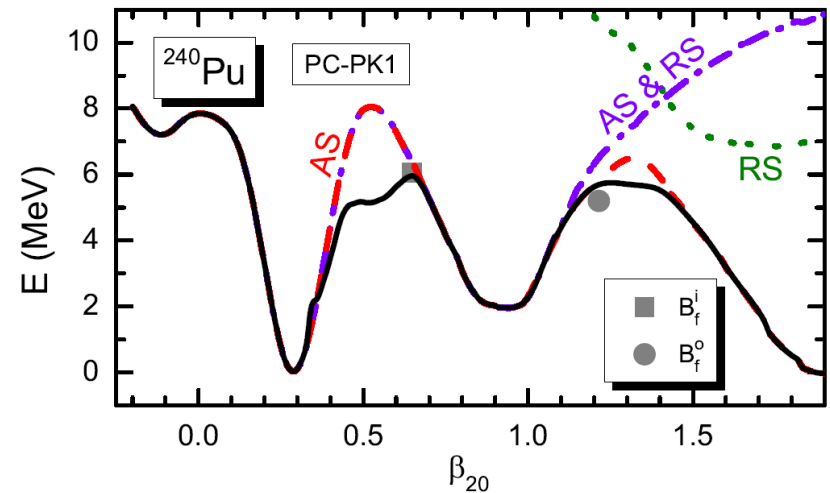
[Zhao_Lu_Zhao_SGZ 2012_PRC86-057304](#)

[Lu_Zhao_Zhao_SGZ 2014_PRC89-014323](#)

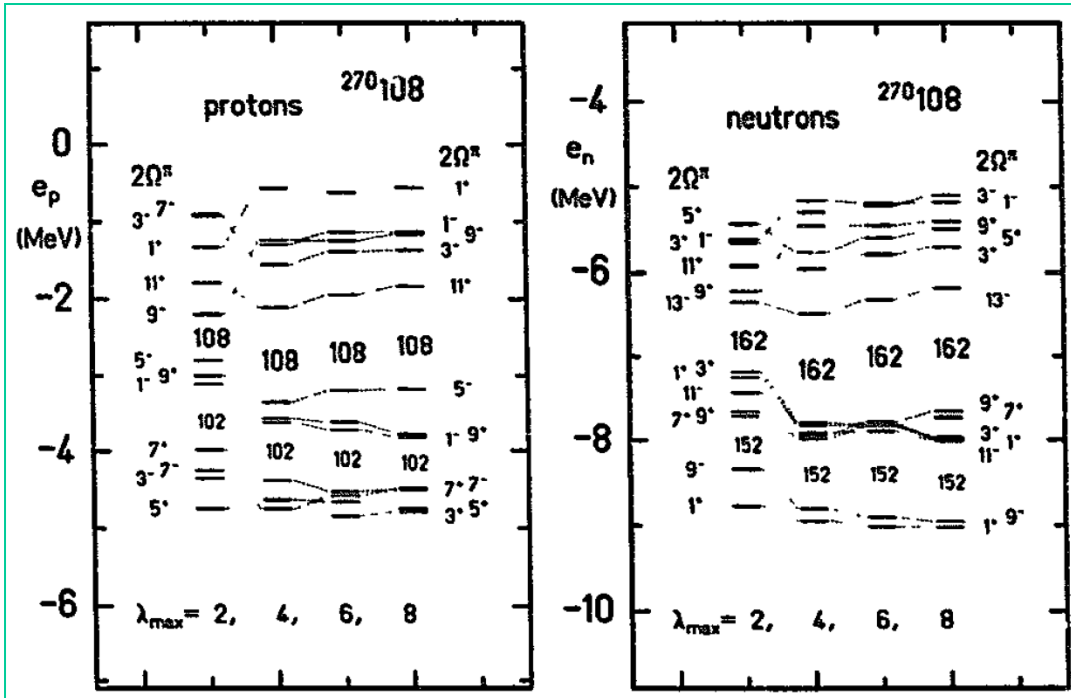
^{240}Pu : 3-dim. PES (β_{20} , β_{22} , β_{30})



- AS & RS for g.s. & isomer, the latter is stiffer
- Triaxial & octupole shape around the outer barrier
- Triaxial deformation crucial around barriers

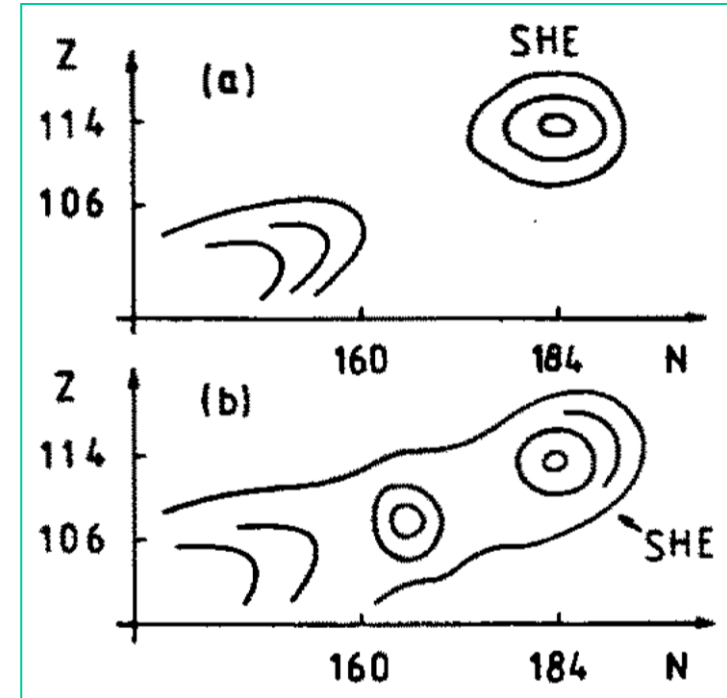
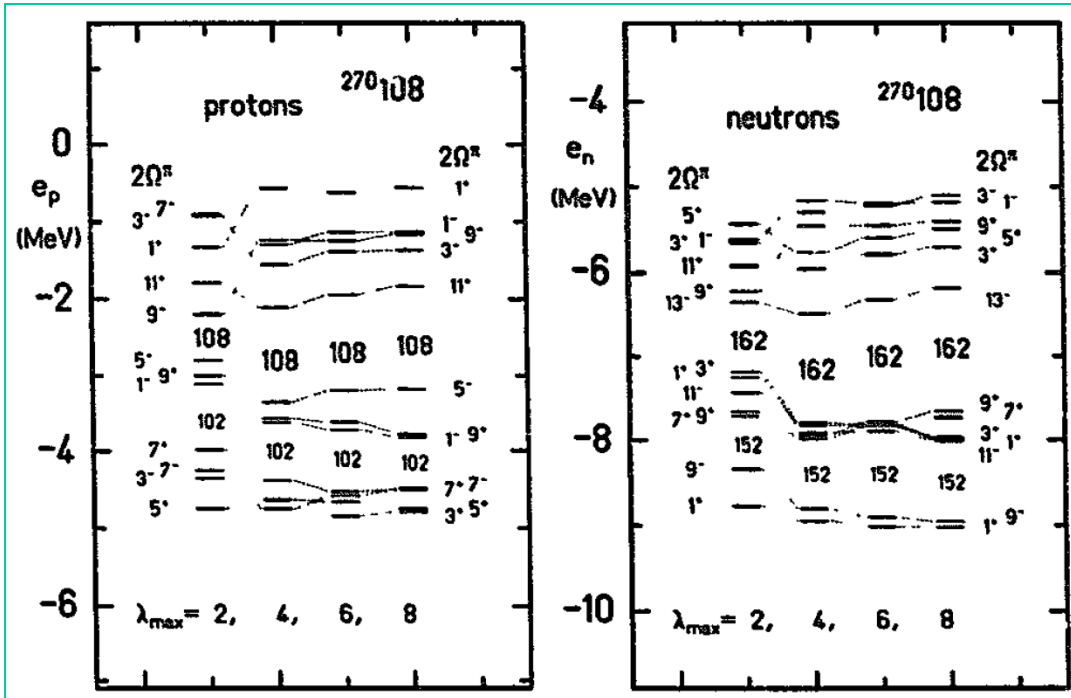


^{270}Hs : A doubly magic deformed SHN



Patyk_Sobiczewski1991_NPA533-132

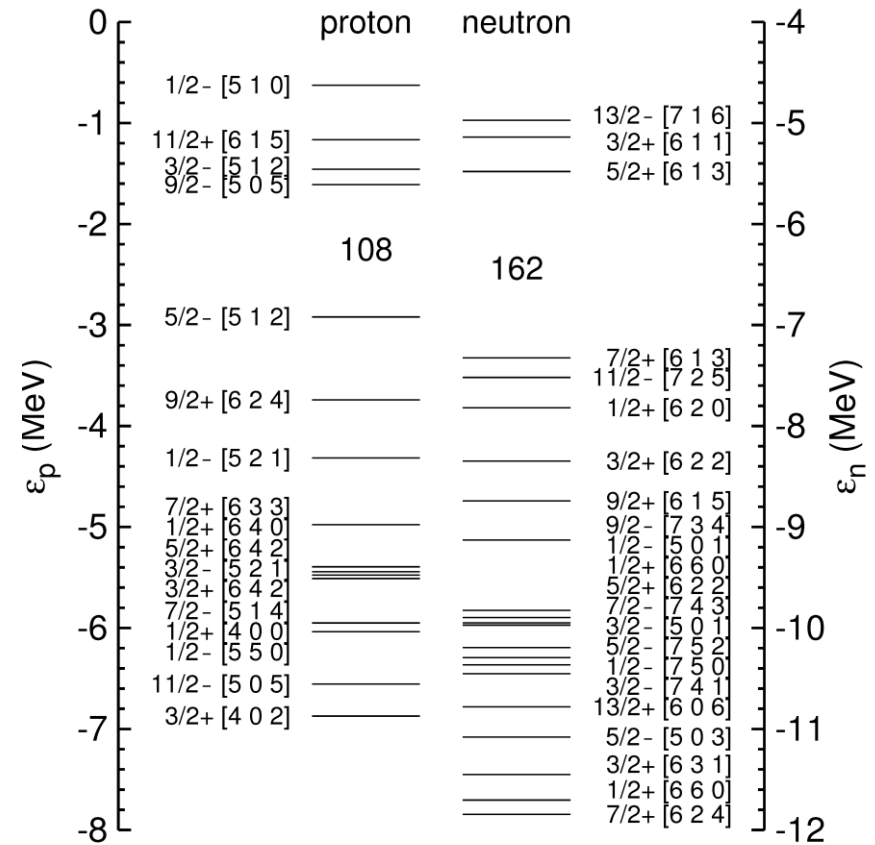
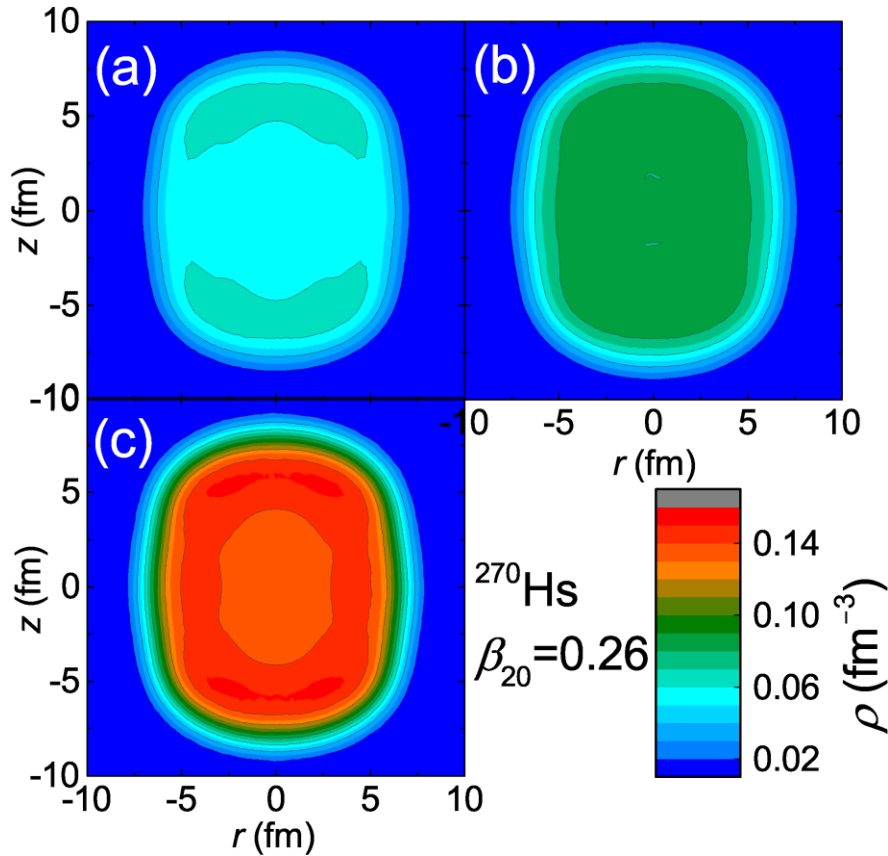
^{270}Hs : A doubly magic deformed SHN



Patyk_Sobiczewski1991_NPA533-132

Patyk_Skalski_Sobiczewski_Cwiok
1989_NPA502-591c

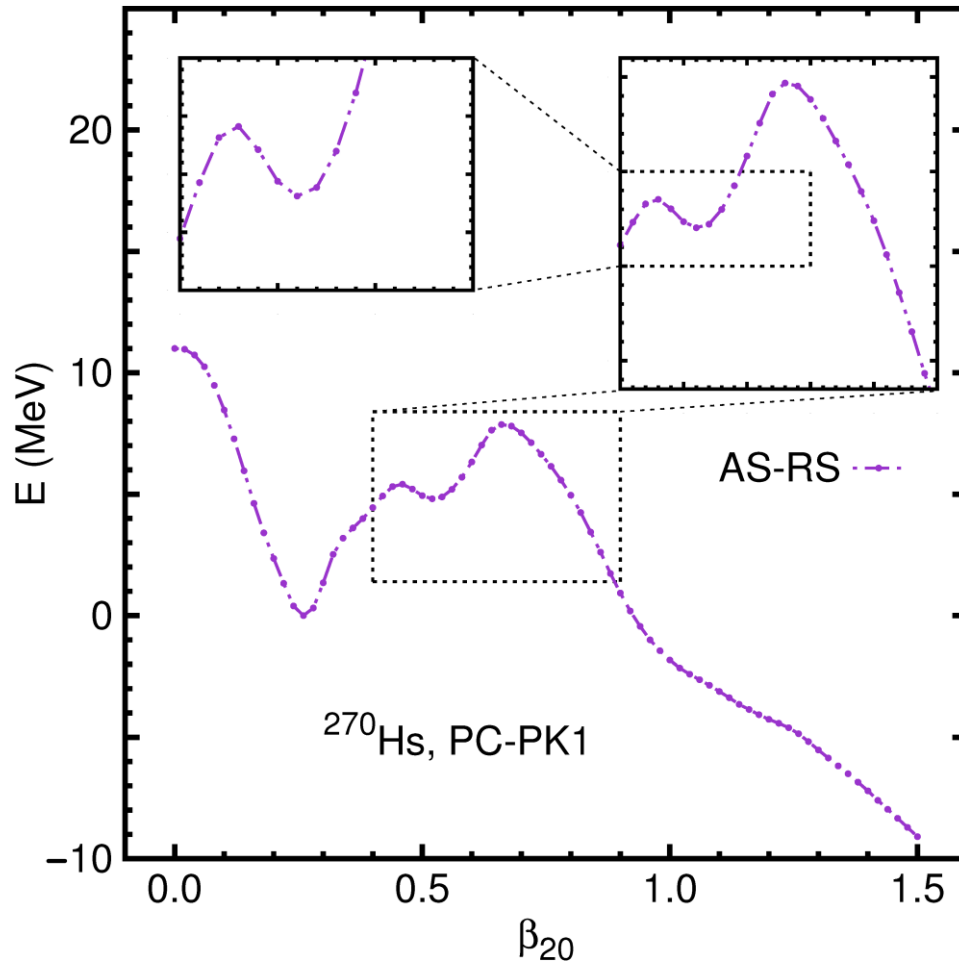
^{270}Hs : ground state from MDC-RMF calc.



Xu Meng (孟旭), PhD thesis

Meng_Lu_Zhou2020
Sci. China-Phys. Mech. Astron. 63, 212011

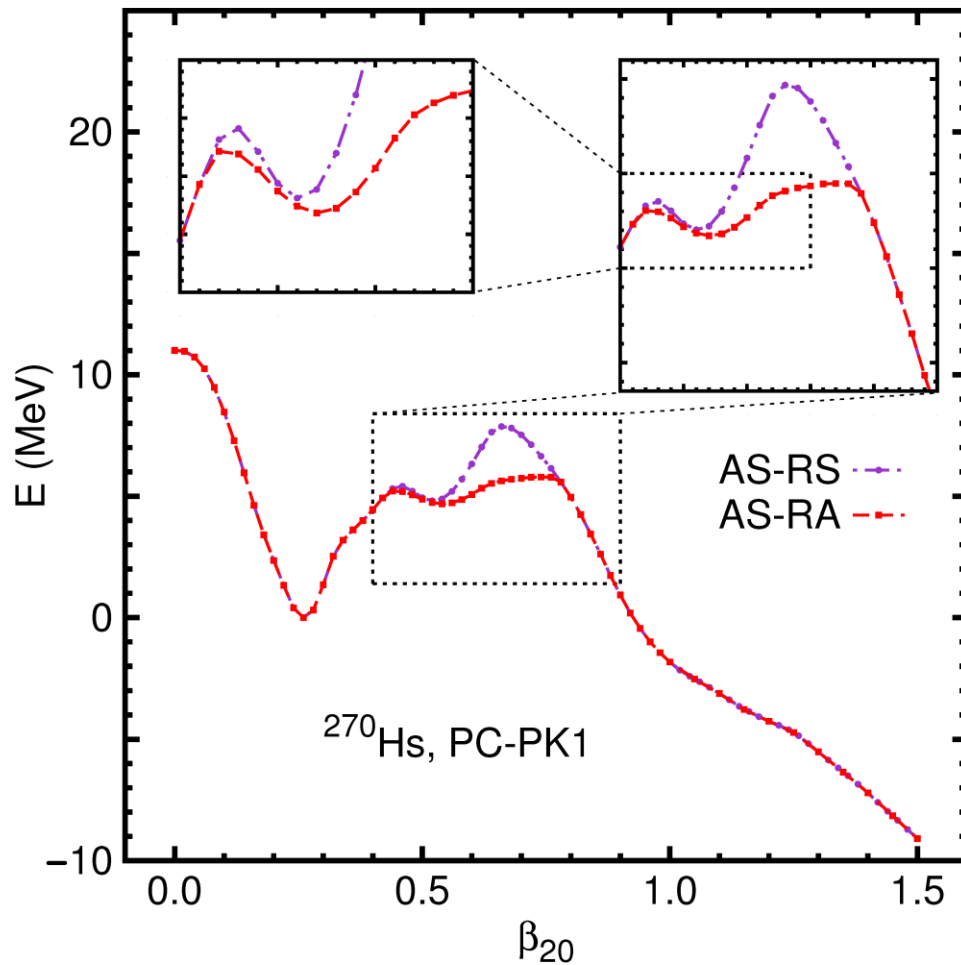
^{270}Hs : 1D PEC from MDC-RMF calc.



AS-RS: Axially-Symmetric &
Reflection Symmetric

Courtesy of Xu Meng (孟旭)

^{270}Hs : 1D PEC from MDC-RMF calc.

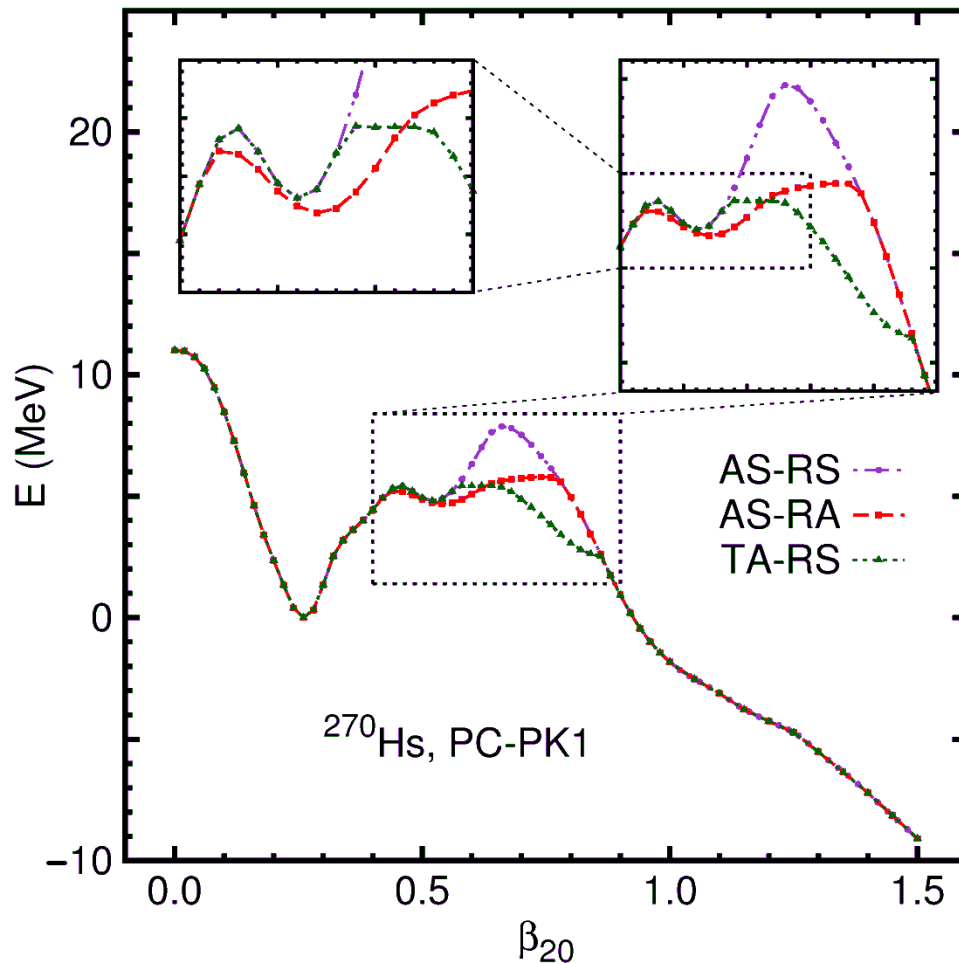


AS-RS: Axially-Symmetric &
Reflection Symmetric

AS-RA: Axially-Symmetric &
Reflection Asymmetric

Courtesy of Xu Meng (孟旭)

^{270}Hs : 1D PEC from MDC-RMF calc.



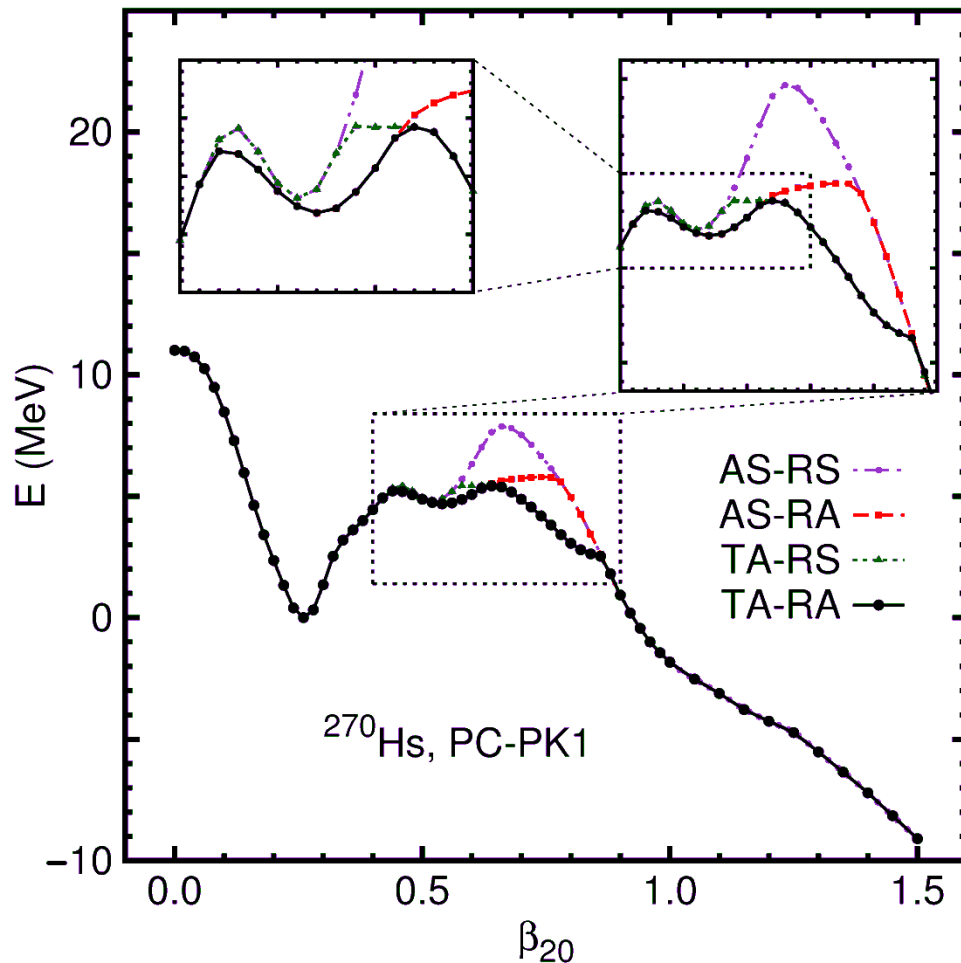
AS-RS: Axially-Symmetric &
Reflection Symmetric

AS-RA: Axially-Symmetric &
Reflection Asymmetric

TA-RS: TriAxial &
Reflection Symmetric

Courtesy of Xu Meng (孟旭)

^{270}Hs : 1D PEC from MDC-RMF calc.



AS-RS: Axially-Symmetric &
Reflection Symmetric

AS-RA: Axially-Symmetric &
Reflection Asymmetric

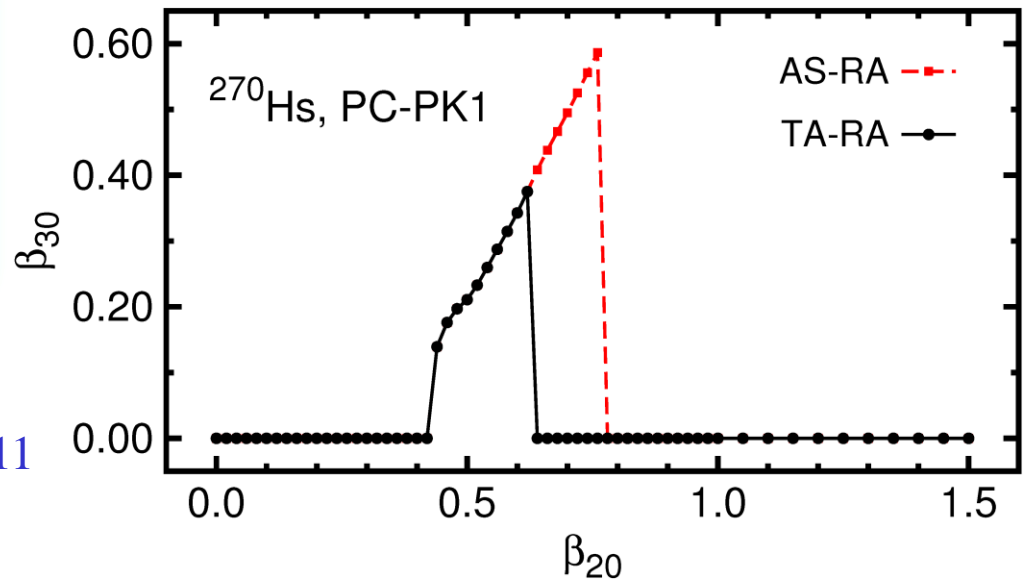
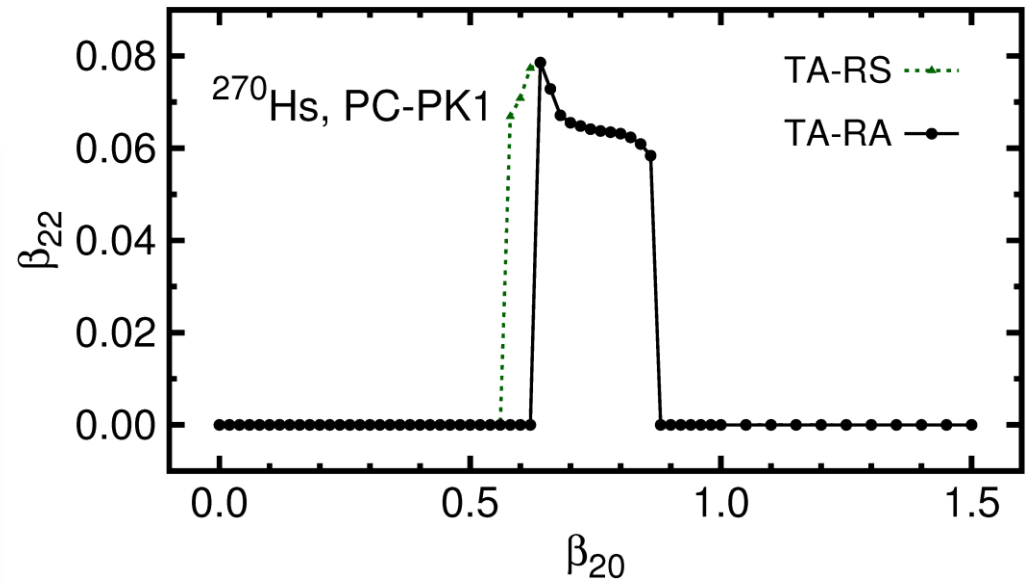
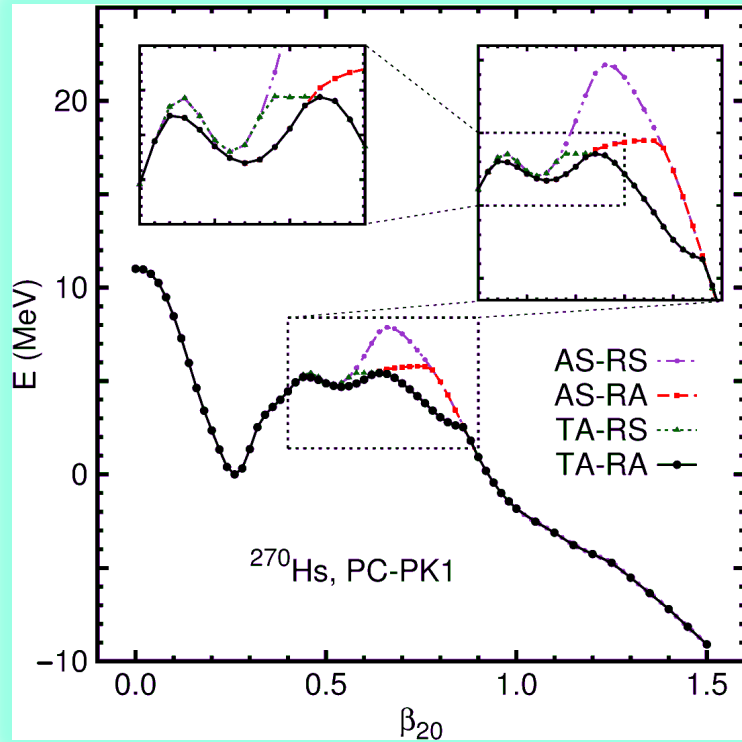
TA-RS: TriAxial &
Reflection Symmetric

TA-RA: TriAxial &
Reflection Asymmetric

Meng_Lu_Zhou2020

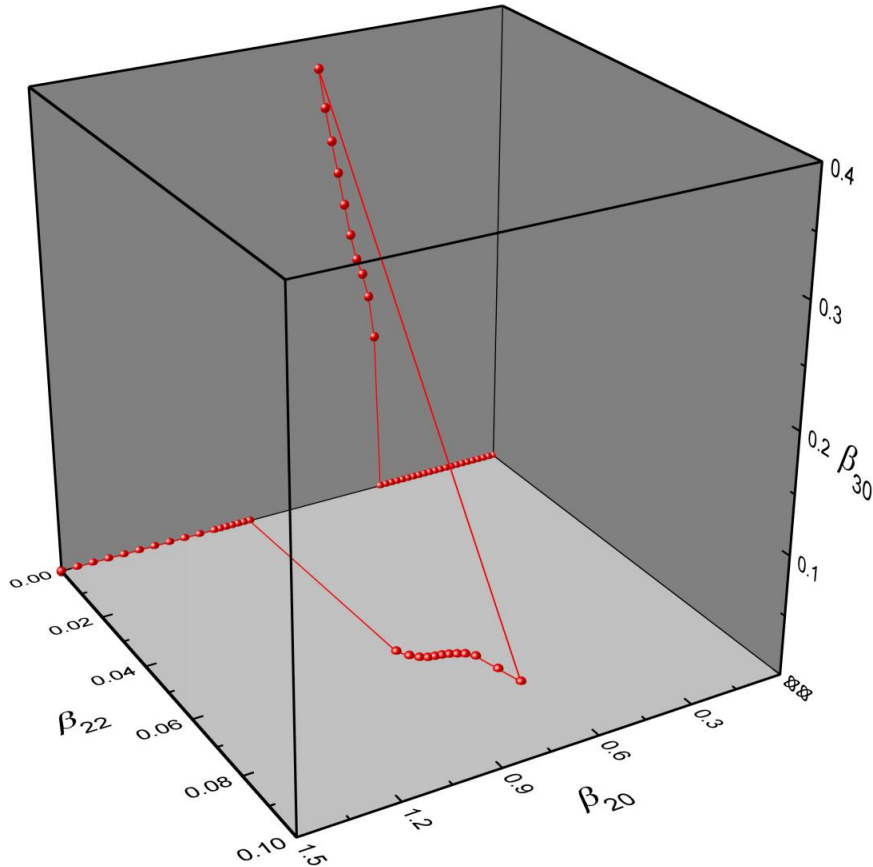
Sci. China-Phys. Mech. Astron. 63, 212011

Discontinuities in 1D PECs



^{270}Hs : 1D TA-RA “PEC” viewed in 3D def. space

$(\beta_{20}, \beta_{22}, \beta_{30})$ deformation space



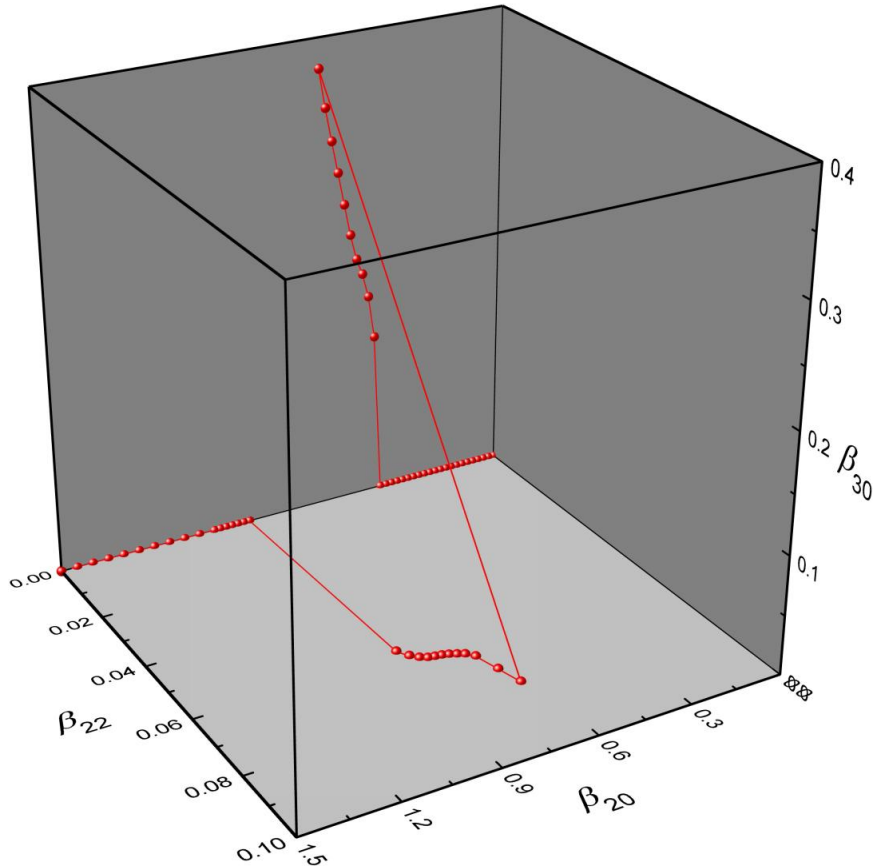
TA-RA “PEC” consists of 4 segments:

- $\beta_{20} = (0, 0.42)$, $\beta_{22} = 0$, $\beta_{30} = 0$
- $\beta_{20} = (0.44, 0.62)$, $\beta_{22} = 0$, $\beta_{30} \neq 0$
- $\beta_{20} = (0.64, 0.86)$, $\beta_{22} \neq 0$, $\beta_{30} = 0$
- $\beta_{20} = (0.88, 1.50)$, $\beta_{22} = 0$, $\beta_{30} = 0$

Courtesy of Xu Meng (孟旭)

^{270}Hs : 1D TA-RA “PEC” viewed in 3D def. space

$(\beta_{20}, \beta_{22}, \beta_{30})$ deformation space



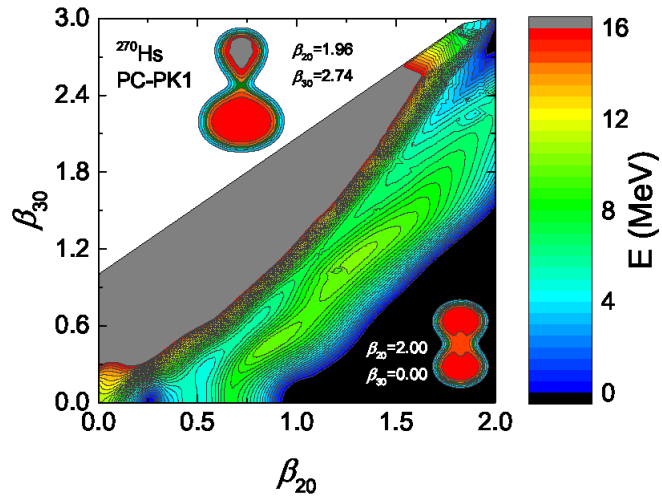
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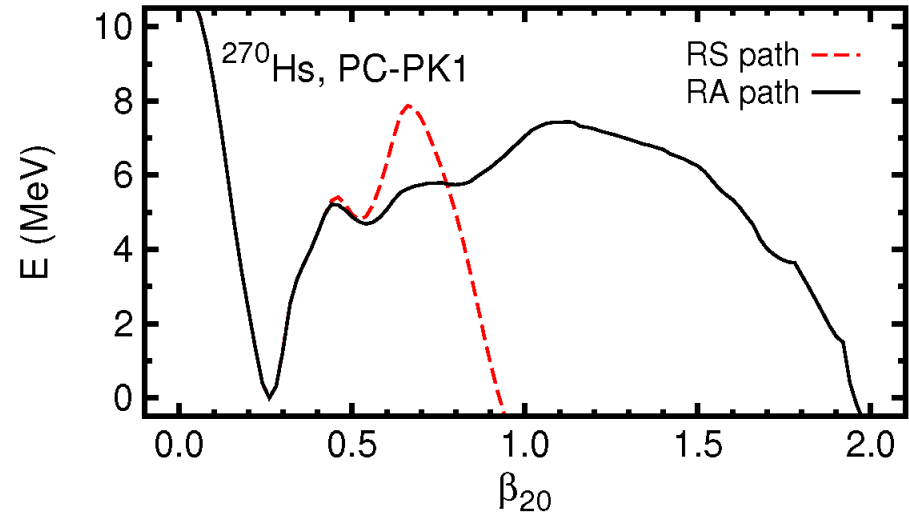
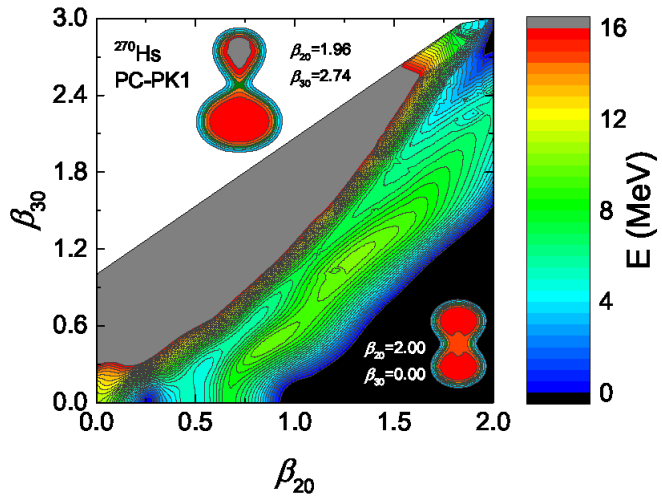
How can a nucleus change its shape so suddenly ?

Courtesy of Xu Meng (孟旭)

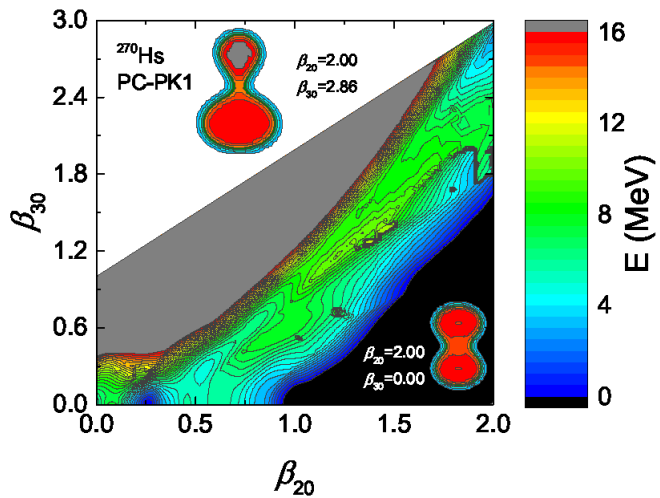
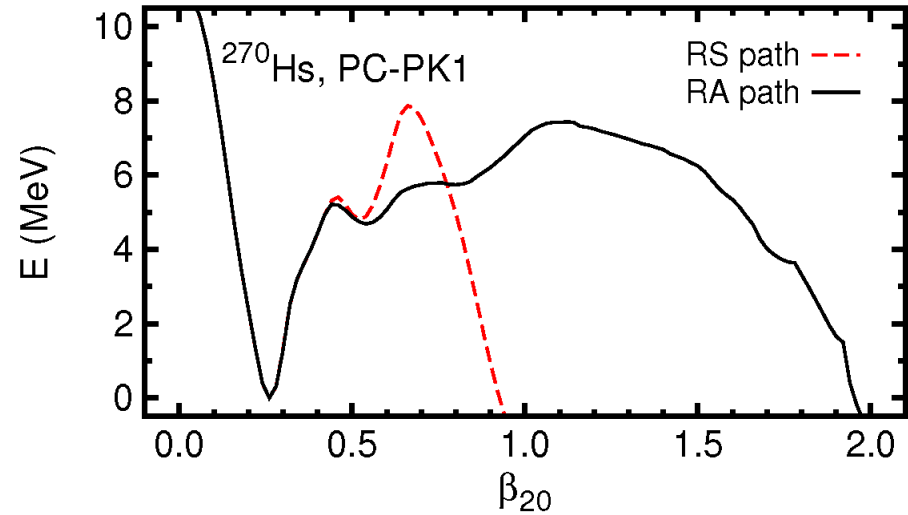
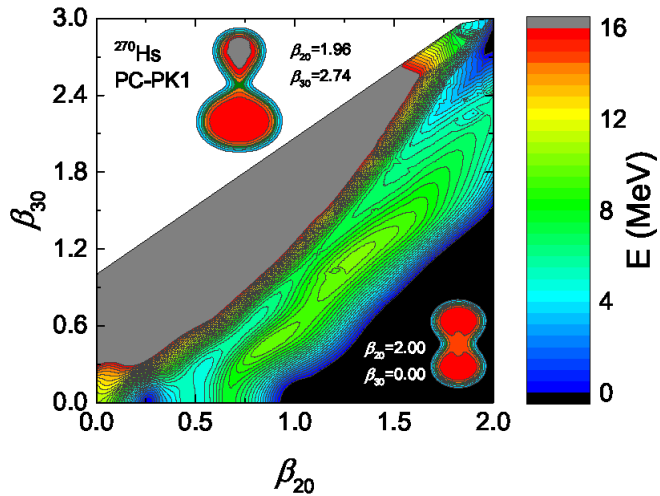
^{270}Hs : 2D PES & optimal static fission path



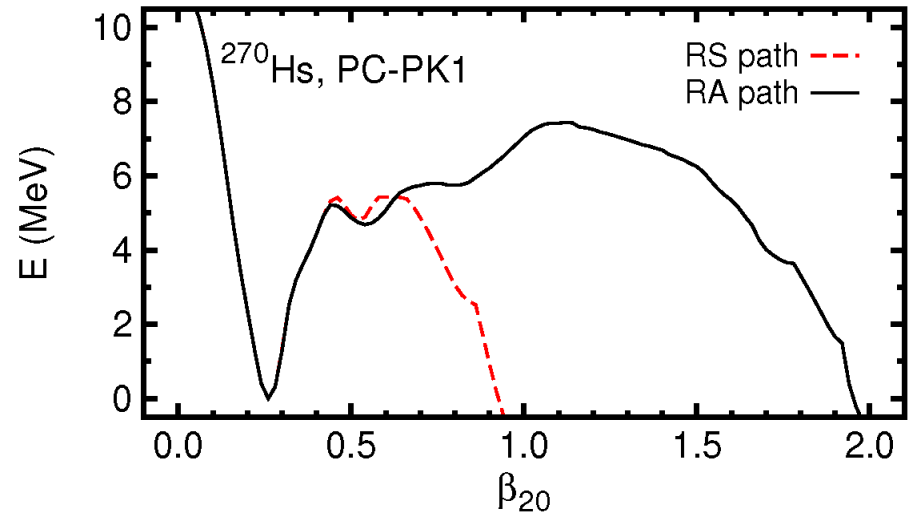
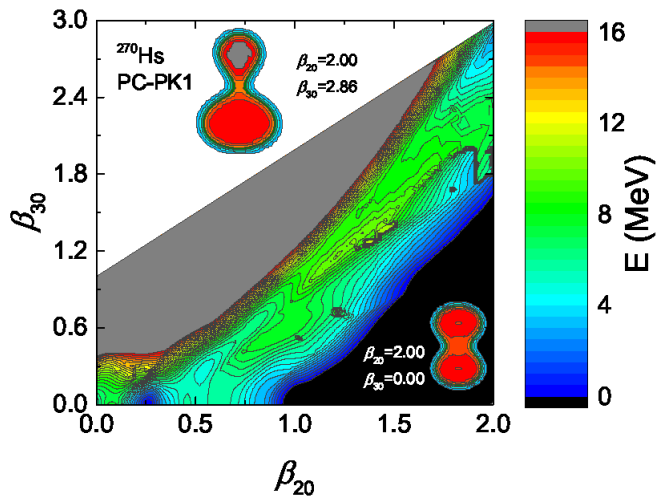
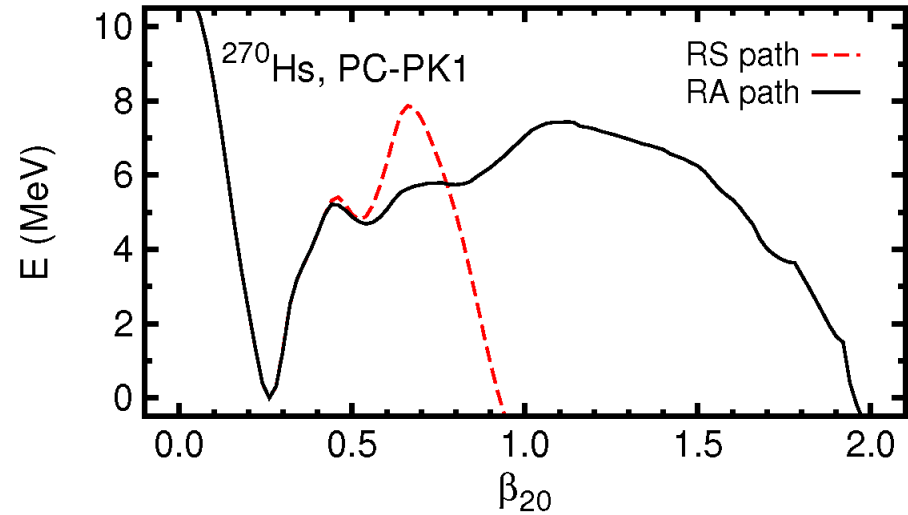
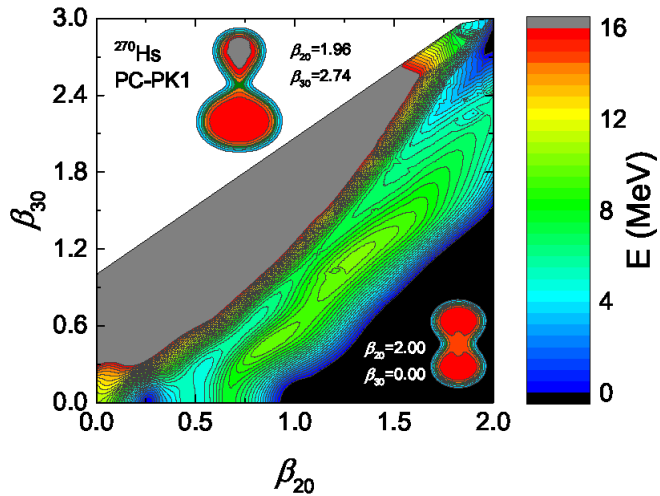
^{270}Hs : 2D PES & optimal static fission path



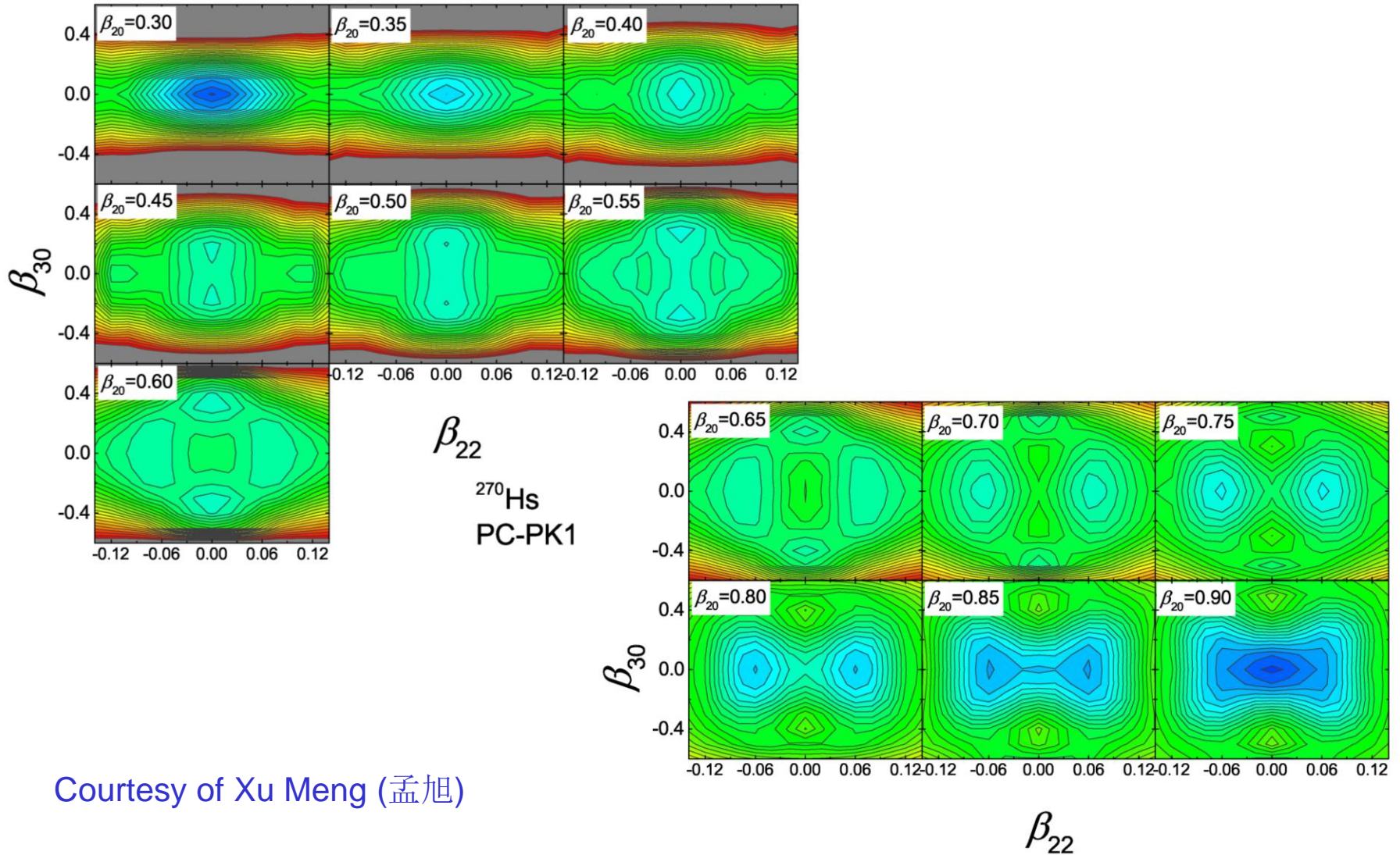
^{270}Hs : 2D PES & optimal static fission path



^{270}Hs : 2D PES & optimal static fission path



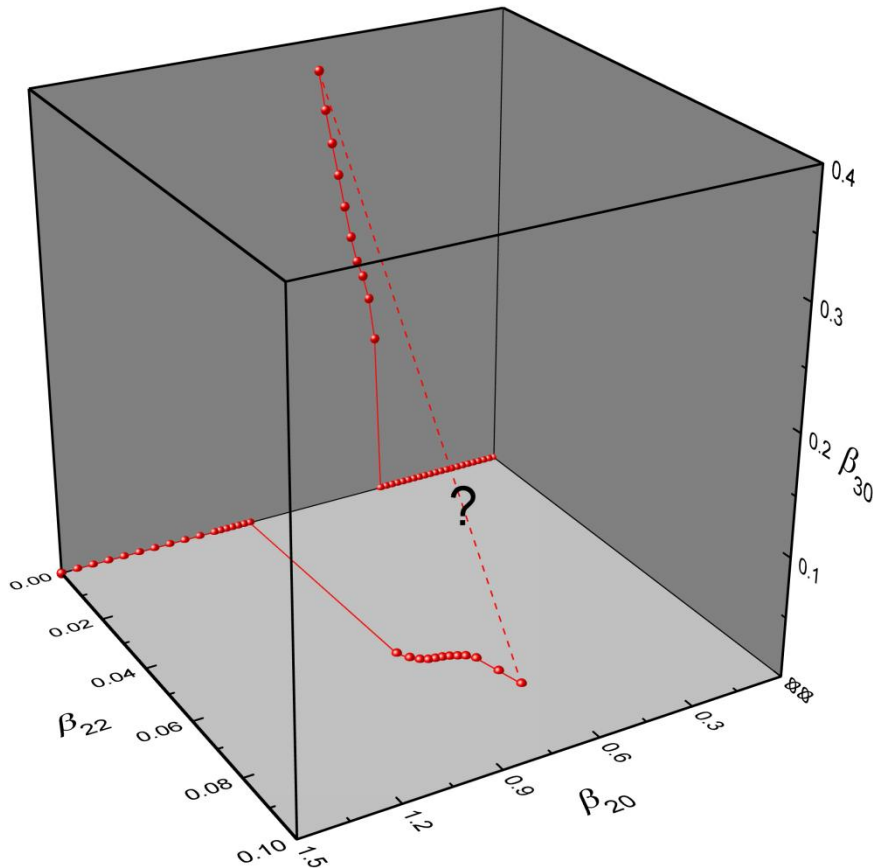
^{270}Hs : 3D PES



Courtesy of Xu Meng (孟旭)

^{270}Hs : 1D TA-RA “PEC” viewed in 3D def. space

$(\beta_{20}, \beta_{22}, \beta_{30})$ deformation space



Courtesy of Xu Meng (孟旭)

TA-RA “PEC” consists of 4 segments:

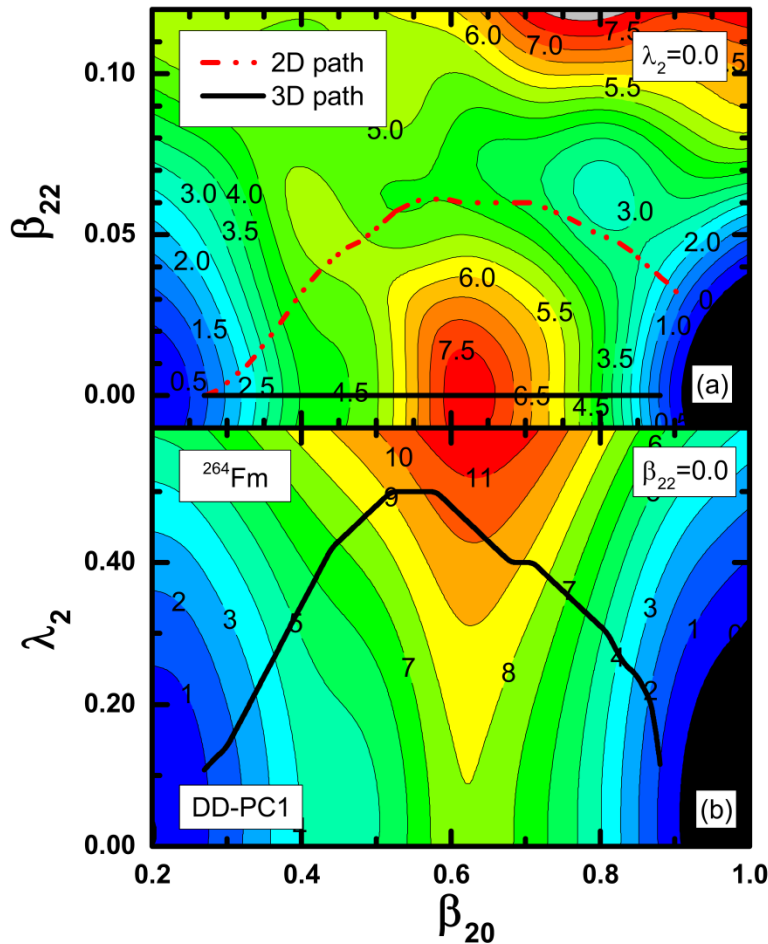
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- $\beta_{20} = (0.44, 0.62), \beta_{22} = 0, \beta_{30} \neq 0$
- $\beta_{20} = (0.64, 0.86), \beta_{22} \neq 0, \beta_{30} = 0$
- $\beta_{20} = (0.88, 1.50), \beta_{22} = 0, \beta_{30} = 0$

How can a nucleus change its shape so suddenly ?

It may do it through when shape DoFs are considered explicitly !

MDC-RMF study of spontaneous fission: Coupling between shape and pairing

WKB w/ action in coll. spaces based on PES from MDC-RMF



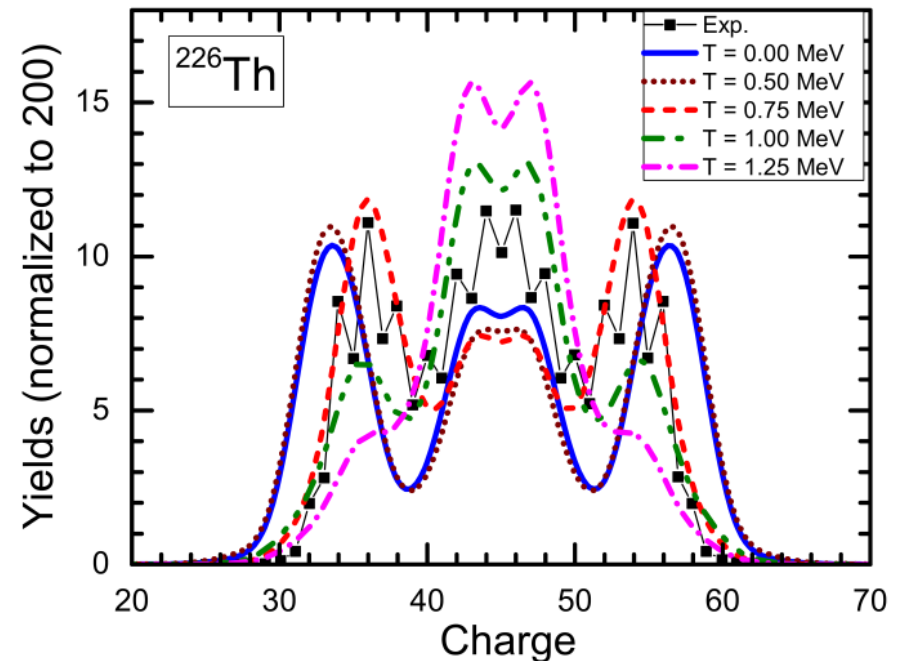
- Pairing favors axially symmetric shapes along the dynamic path of the fissioning system, amplifies pairing as the path traverses the fission barriers, significantly reduces the action integral, and shortens the corresponding SF half-life.

Zhao_Lu_Niksic_Vretenar_SGZ
2016_PRC93-044315

Micro. self-consistent description of induced fission dyn.: finite temperature (FT) effects

TDGCM + GOA based on PES from FT MDC-RMF

- The model can qualitatively reproduce the empirical triple-humped structure of the fission charge and mass distributions at $T = 0$
- But the precise expt. position of the asymmetric peaks & the symmetric-fission yield can only be accurately reproduced when FT effects included



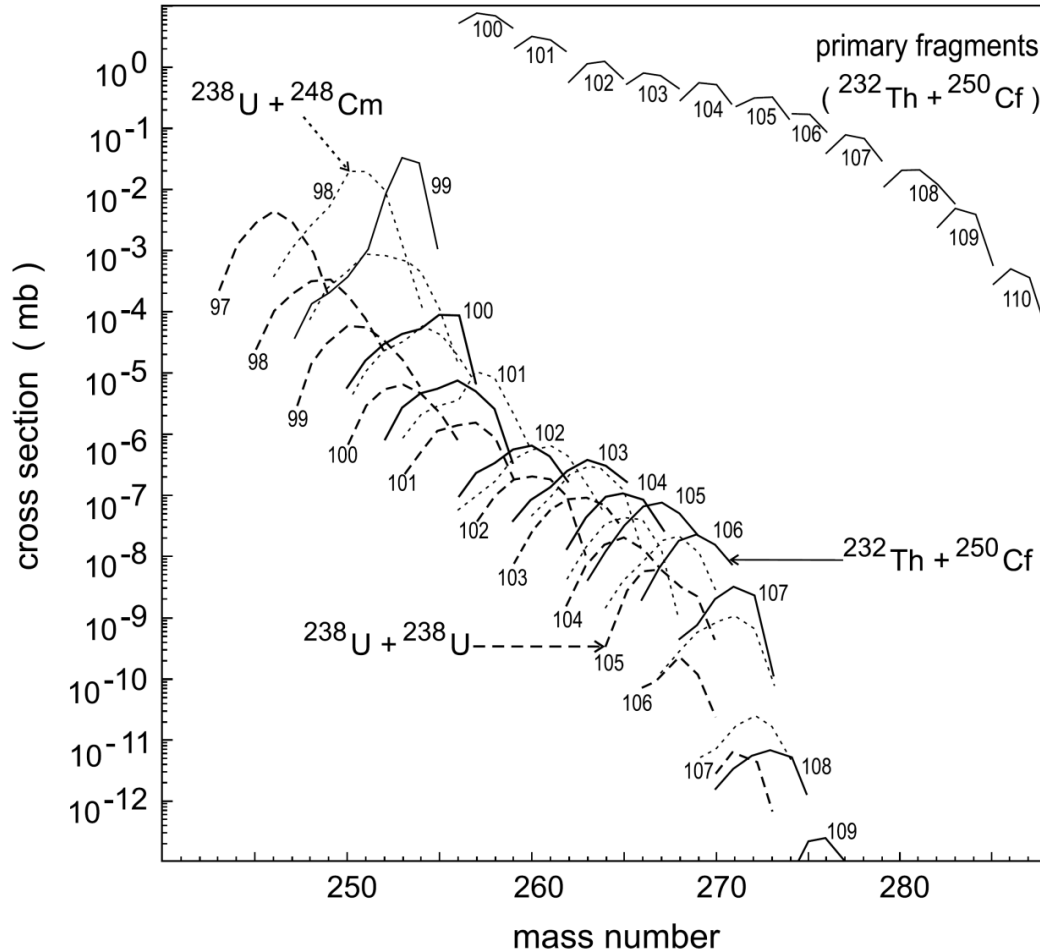
Zhao_Niksic_Vretenar_SGZ
2018_arXiv1809.06114

Lectures 3 & 4

- Challenges in synthesizing SHN

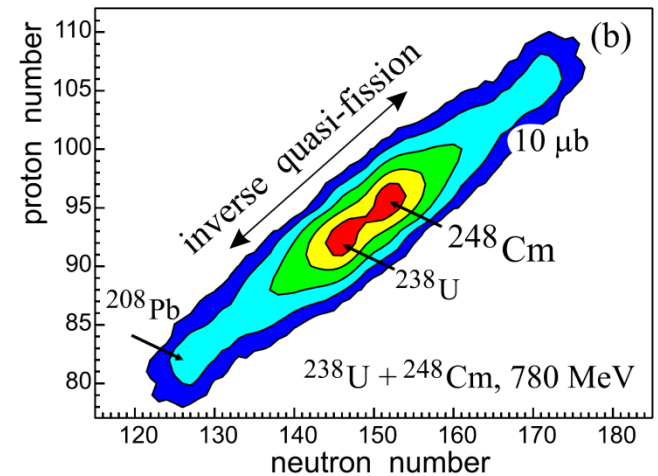
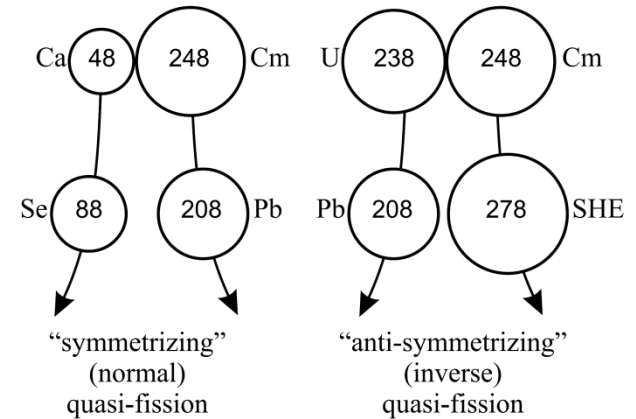
- Synthesis mechanism of SHN
 - Large uncertainties in predicted Xsections
 - Heavy ion fusion reactions
 - Capture
 - Fusion
 - Survival against fission
 - **Multi-nucleon transfer reactions**
 - Reactions w/ radioactive ion beams

Multi-nucleon transfer reactions: Langevin dynamics



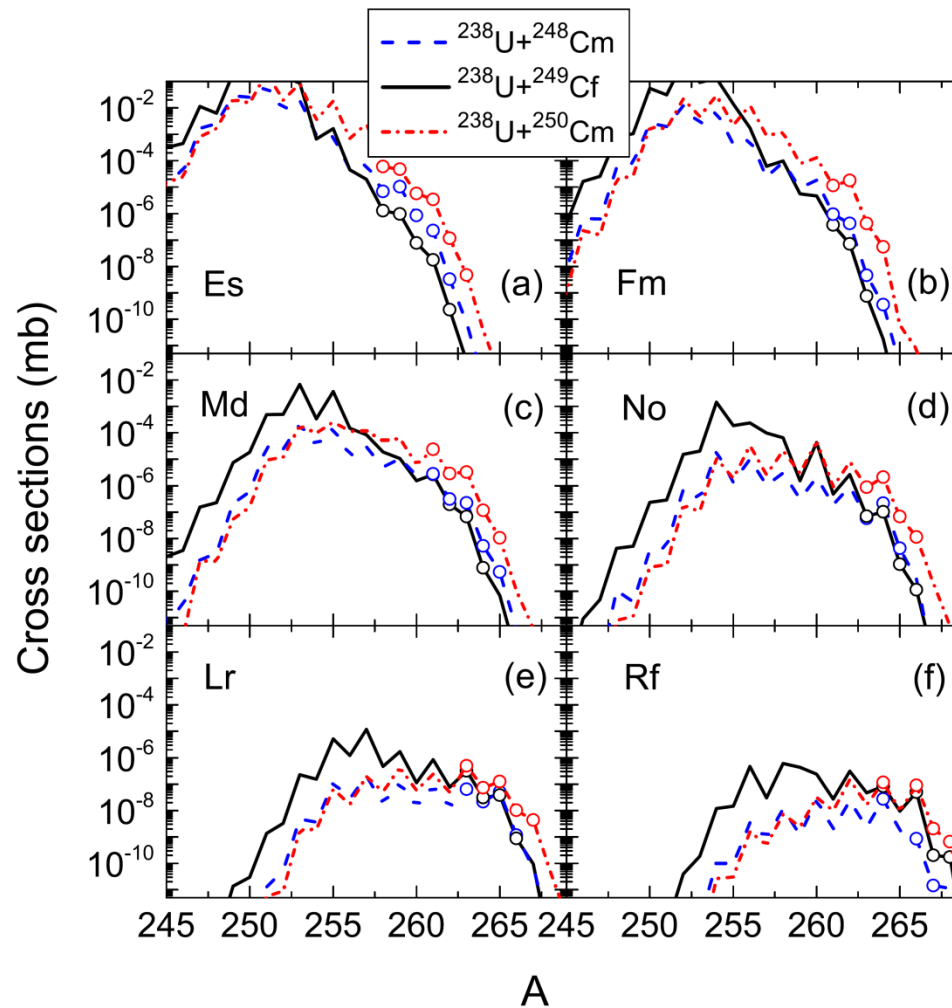
Zagrebaev+2006_PRC73-031602R

Inverse quasi-fission



Zagrebaev&Greiner2015_NPA944-257

Multi-nucleon transfer reactions: DNS model



Zhu+2016_PRC94-054606

Fusion of two charged liquid drops

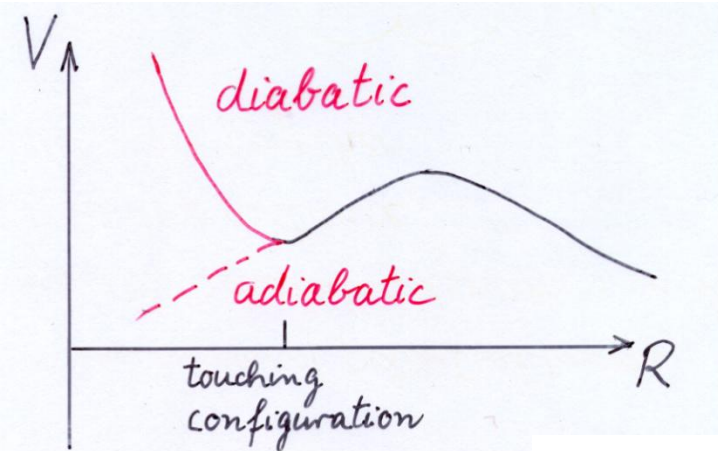
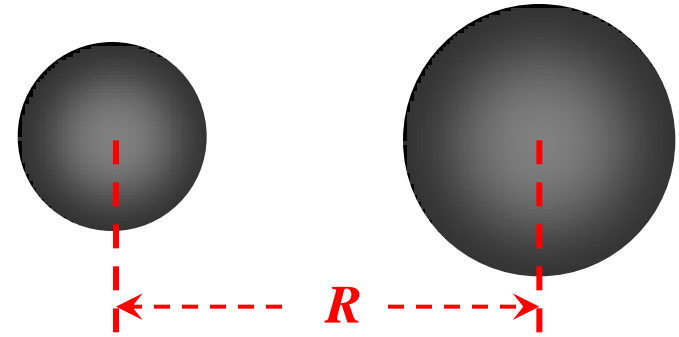
□ Langevin dynamics

$$\left\{ \begin{array}{l} \frac{du(t)}{dt} = - \int_{-\infty}^t \gamma(t-t')u(t')dt' + \frac{1}{\mu}\delta F(t) - \frac{1}{\mu} \frac{dV(R)}{dR} \\ u(t) = \frac{dR(t)}{dt} \end{array} \right.$$

- $R(t)$: Rel. distance
- $u(t)$: Rel. velocity
- $V(R)$: Interaction potl.
- $dF(t)$: Random force
- $g(t-t')$: Friction force

□ Dinuclear system (DNS) model

- Projectile & target keep staying in the potl. Pocket and individuality
- Transfer of nucleons betw. Projectile & target may lead to fusion



W. Scheid

Nucleon(s) transfer betw. two charged liquid drops

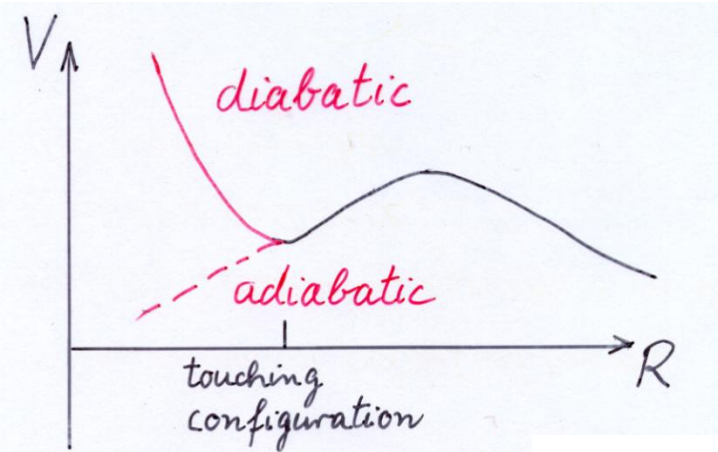
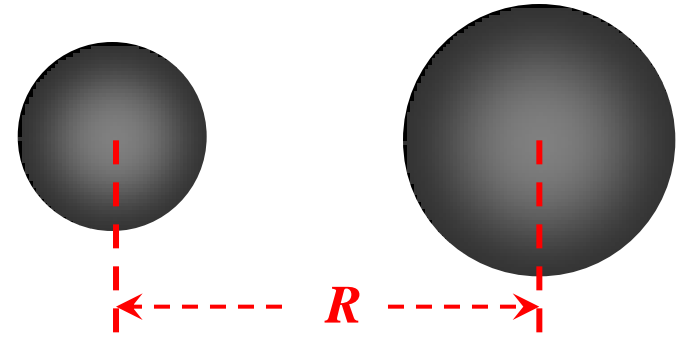
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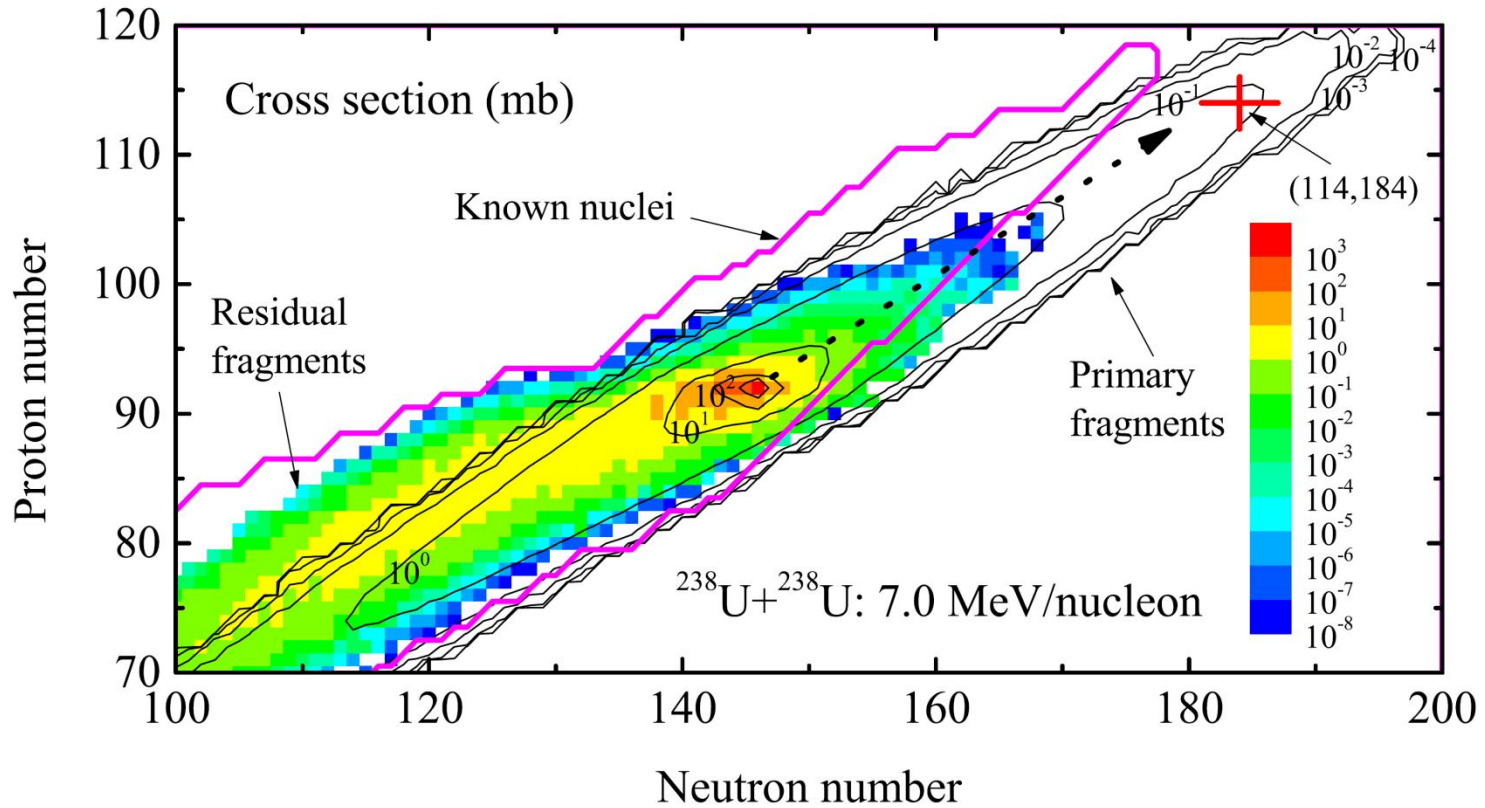
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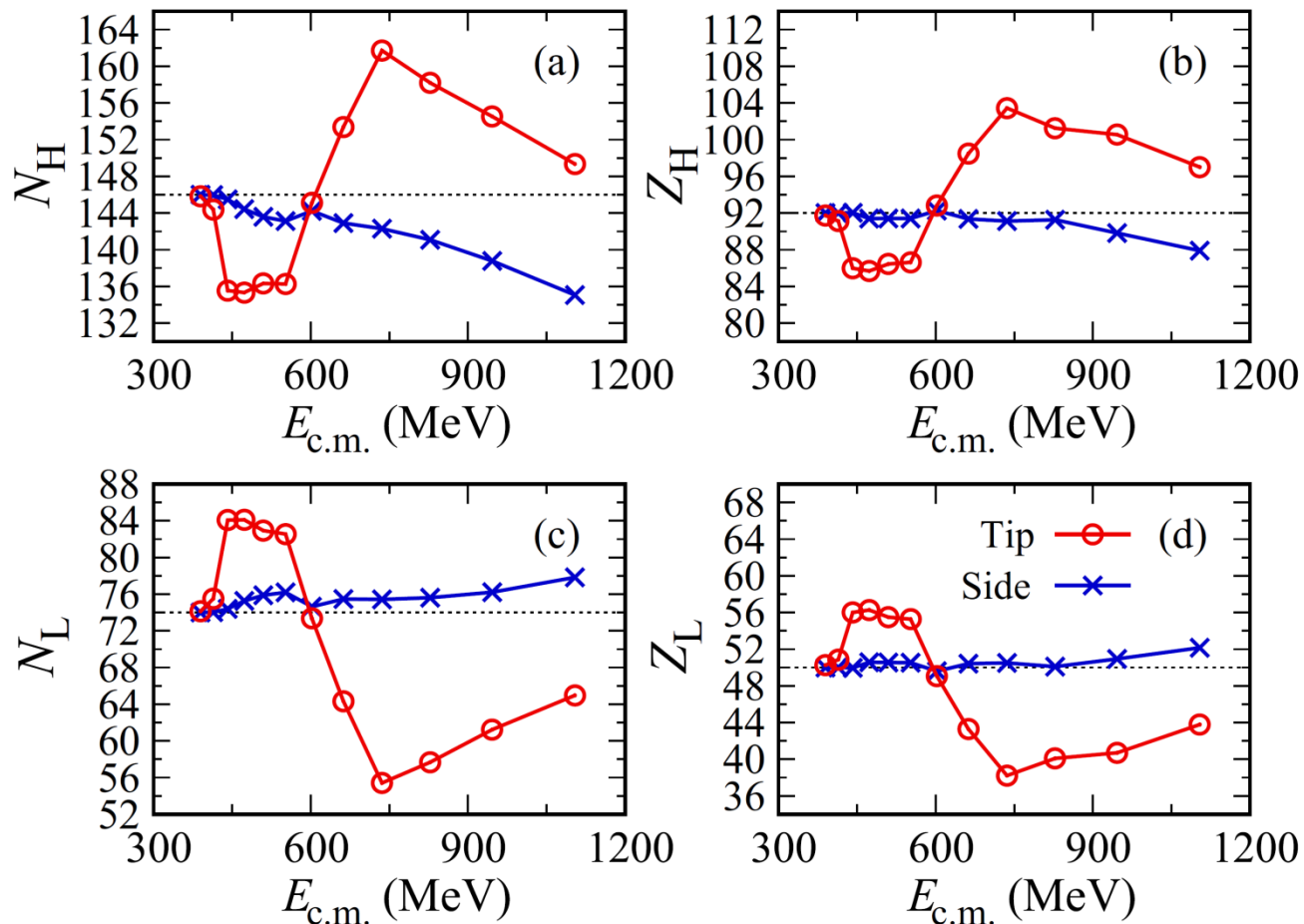
W. Scheid

Multi-nucleon transfer reactions: ImQMD model



Zhao+2016_PRC94-024601

Multi-nucleon transfer reactions: TDHF theory



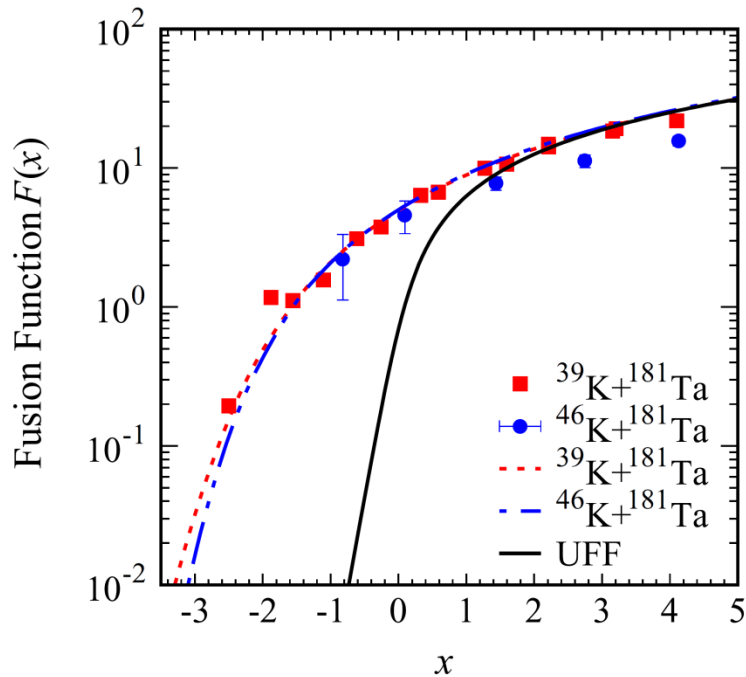
Lectures 3 & 4

- Challenges in synthesizing SHN

- Synthesis mechanism of SHN
 - Large uncertainties in predicted Xsections
 - Heavy ion fusion reactions
 - Capture
 - Fusion
 - Survival against fission
 - Multi-nucleon transfer reactions
 - Reactions w/ radioactive ion beams

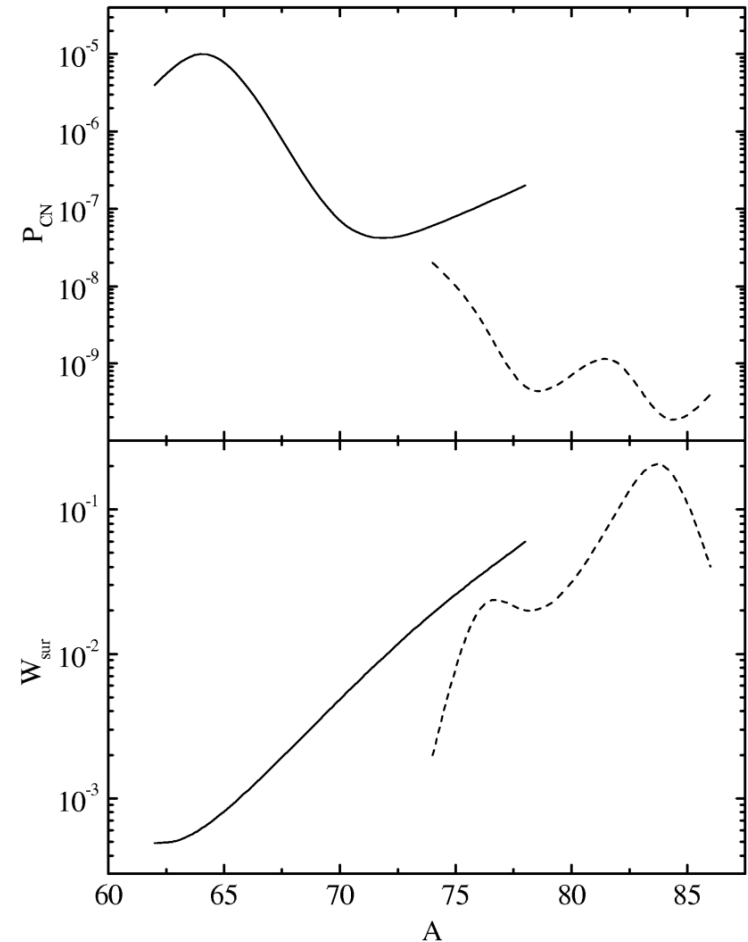
Synthesis of SHN w/ radioactive ion beams

$$\sigma_{\text{ER}}(E_{\text{cm}}) = \sum_J \sigma_{\text{cap}}(E_{\text{cm}}, J) P_{\text{CN}}(E_{\text{cm}}, J) W_{\text{sur}}(E_{\text{cm}}, J)$$



Wang+2018_PRC98-014615

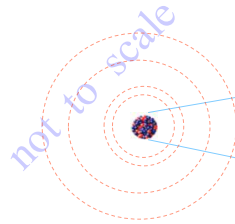
A thorough investigation:
Wu+2018_PRC97-064609



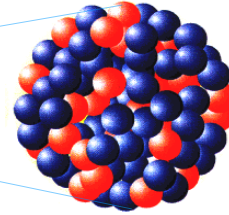
Adamian_Antonenko&Scheid2000_NPA678-24

Summary

10^{-9} m



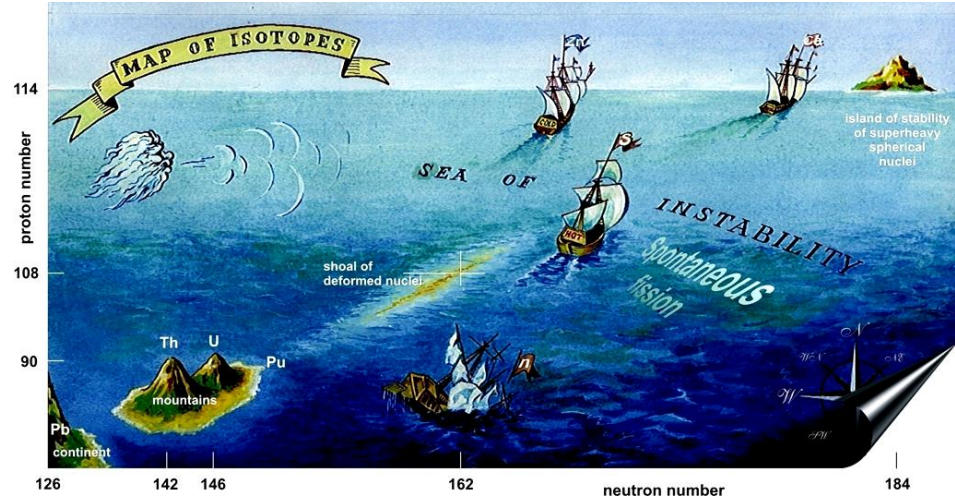
10^{-14} m



带*为放射性元素
其中黑色为天然放射性元素
红色为人造元素

IA	主族金属																IIIA	IVA	VA	VIA	VIIA	0
1	副族金属																5	6	7	8	9	10
2	副族金属																13	14	15	16	17	18
3	副族金属																31	32	33	34	35	36
4	副族金属																49	50	51	52	53	54
5	副族金属																63	64	65	66	67	68
6	副族金属																79	80	81	82	83	84
7	副族金属																89	90	91	92	93	94

镧系	57	58	59	60	61	62	63	64	65	66	67	68	69	70	71
	La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu
锕系	89	90	91	92	93	94	95	96	97	98	99	100	101	102	103
	Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr



- ❑ Where is the end of PT of elements?
- ❑ Are there stable high-atomic-number elements?
- ❑ What are similarities & differences in chemistry of SHEs & their lighter homologs
- ❑ ...

- ❑ Where is the island of stability of SHN?
- ❑ Are there stable or long-lived SHN?
- ❑ How to reach the island of stability?
- ❑ Are there exotic shapes in SHN?
- ❑ Are there isomers in SHN longer-lived than their ground states?
- ❑ ...