

Interdisciplinary Study with Quantum Beams

- Application to materials analysis

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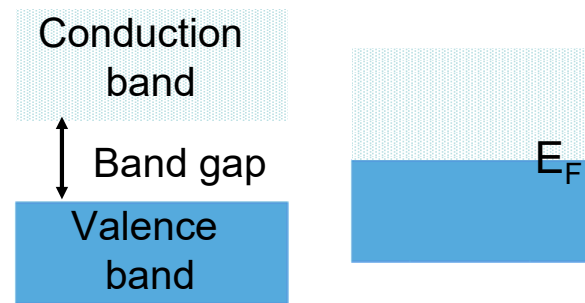
1. Materials analysis : what do we want to know?
2. Quantum beams
 - e^- , e^+ , Photon, Ion, Neutron, Muon
 - Interaction of the beams with materials
 - Information obtained by the beams
3. Application: H in Pd
 - Energetics, dynamics, and kinetics
 - Quantum effects and chemical reaction
4. Summary

Development of novel materials

Materials properties

Electrical
Magnetic
Optical
Chemical
.....

Electronic band structure



Semiconductor

Metal

✓ What we want to know:

Structure and element composition

atomic, mesoscopic, macroscopic scales

→ prediction, design, and production of highly functional materials

Energy-related materials: battery

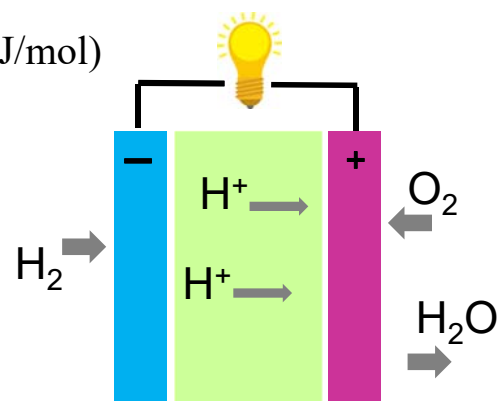
• Fuel cell $\text{H}_2 + \frac{1}{2}\text{O}_2 \rightarrow \text{H}_2\text{O} + Q(286 \text{ kJ/mol})$

Reaction at electrode surfaces

Proton transport

Cf. Photocatalytic H_2 production

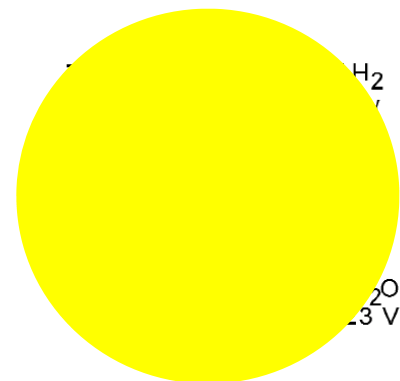
H storage in materials



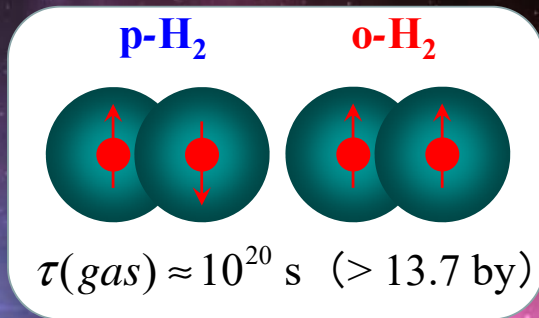
• Li-ion battery

Li ion transport → electricity

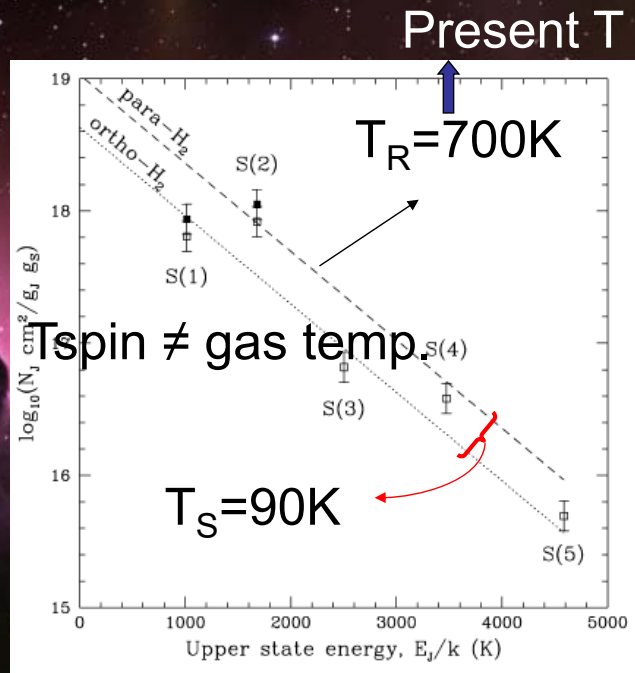
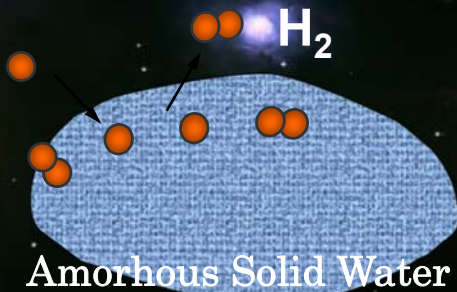
➔ Surfaces & interfaces are important
Light elements often play key roles



Molecular hydrogen in space



- ✓ $\text{H} + \text{H} \rightarrow \text{H}_2$
- ✓ Spin conversion



D.A. Neufeld et al., ApJ 506 (1998) L75,
KF, Prog. Surf. Sci. 88 (2013) 279.

Quantum beam

Interaction with materials (aggregates of atoms)

| A | Interaction | |
|------------|-----------------|---------|
| e^-, e^+ | Coulomb force | } W & P |
| Photon | Charge density | |
| Neutron | Nuclear force | |
| Ion | Coulomb+Nuclear | } P |

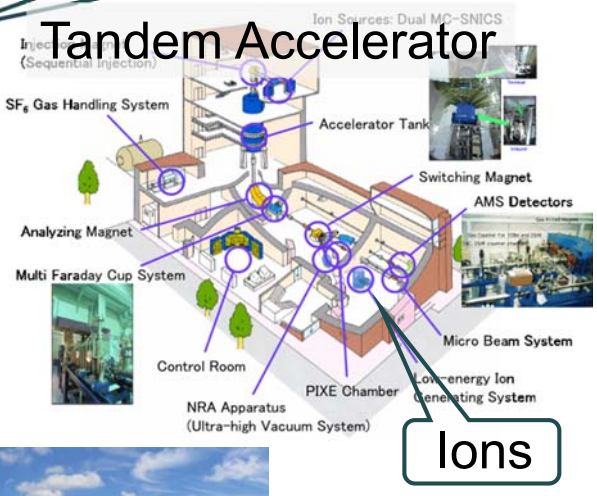
Wave/ particle
duality

$$\lambda = \frac{h}{p}$$

Quantum beam facilities in Japan



23 Neutron beam lines
3 Muon beam lines



UV - X ray, e⁺



Elastic scattering of wave

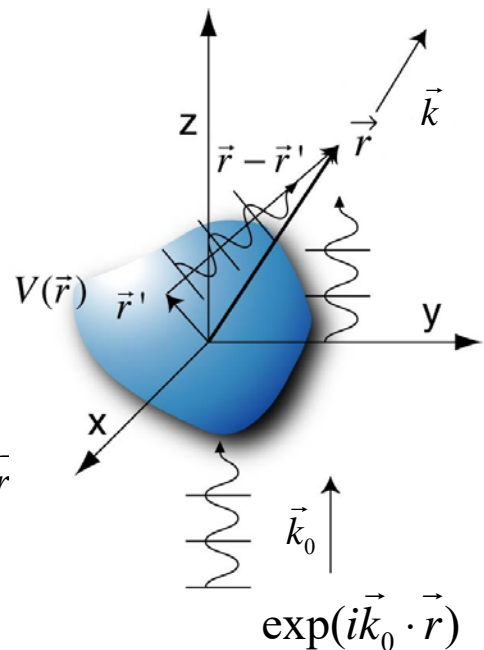
$$\left[-\frac{\hbar^2}{2m} \Delta + V(\vec{r}) \right] \psi(\vec{r}) = E\psi(\vec{r})$$

$$\frac{2m}{\hbar^2} V(\vec{r}) = U(\vec{r}), \quad E = \frac{\hbar^2}{2m} k^2$$

$$[\Delta + k^2] \psi(\vec{r}) = U(\vec{r})\psi(\vec{r})$$

Solution in the integral form

$$\psi(\vec{r}) = \underbrace{\psi_0(\vec{r})}_{\text{Incident wave}} + \underbrace{\int G(\vec{r}, \vec{r}') U(\vec{r}') \psi(\vec{r}') d\vec{r}'}_{\text{Scattering wave}}$$



Cf.) Green function for a free particle

$$G(\vec{r}, \vec{r}') = -\frac{1}{4\pi} \frac{\exp(ik|\vec{r} - \vec{r}'|)}{|\vec{r} - \vec{r}'|} \quad [\Delta + k^2] G(\vec{r}, \vec{r}') = \delta(\vec{r} - \vec{r}')$$

Elastic scattering of wave: diffraction

Weak scattering (Born approx.)

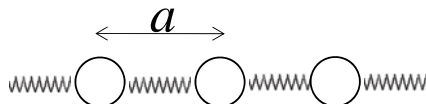
$$\psi^s(\vec{r}) \approx \frac{\exp(ikr)}{4\pi r} \int \exp(i(\vec{k}_0 - \vec{k}) \cdot \vec{r}') U(\vec{r}') d\vec{r}'$$

Spherical wave

Scatt. Amplitude (in the k direction)

= Fourier transform of potential

Crystal: periodic potential



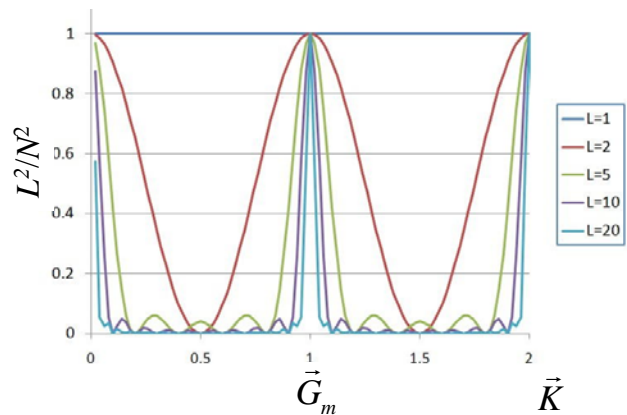
$$U(\vec{r}) = \sum_{\vec{R}_n} U_{unit}(\vec{r} - \vec{R}_n)$$

$$\vec{R}_n = n_1 \vec{a}_1$$

$$I(\vec{K}) \propto \delta(\vec{K} - \vec{G}_m)$$

$$\vec{G}_m = \frac{2\pi m}{a}$$

Reciprocal vector



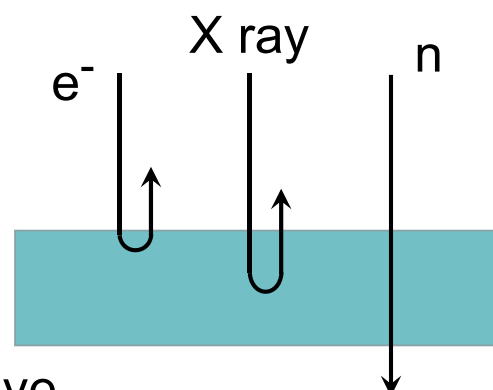
Diffraction → atomic structure

| | E (eV) | λ (Å) | σ (cm ²) | Z dep | |
|------------|-------------|----------------|-----------------------------|---------|----------------|
| e^-, e^+ | ~ 100 | 1.22 | $\sim 10^{-16}$ | Z^2 | Coulomb |
| Photon | $\sim 10^3$ | 1.24 | $\sim 10^{-21}$ | Z^2 | Charge density |
| Neutron | ~ 0.1 | 0.90 | $\sim 10^{-24}$ | - | Nuclear |
| Ion | $\sim 10^6$ | $\sim 10^{-4}$ | | | |

✓ $\lambda \sim$ Interatomic distance

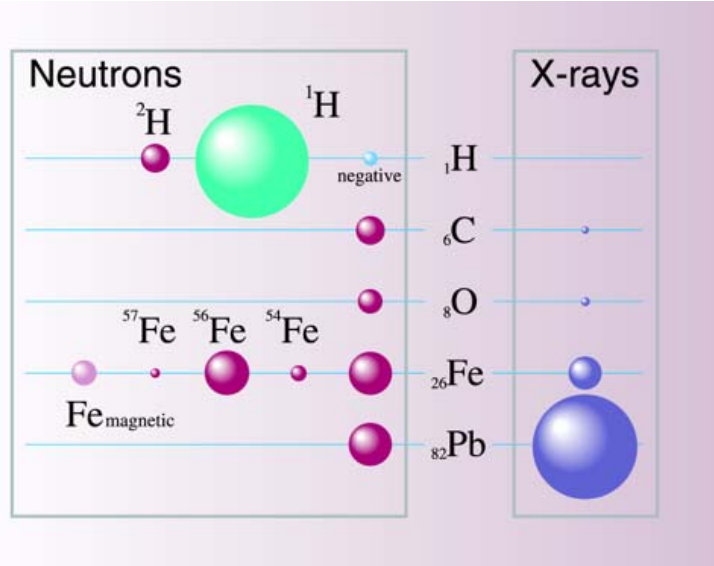
✓ Penetration length
= probing depth

e^- : surface-sensitive



Element specificity

Cross section (\sim atomic size)



Z^2

Sensitive to heavy atom

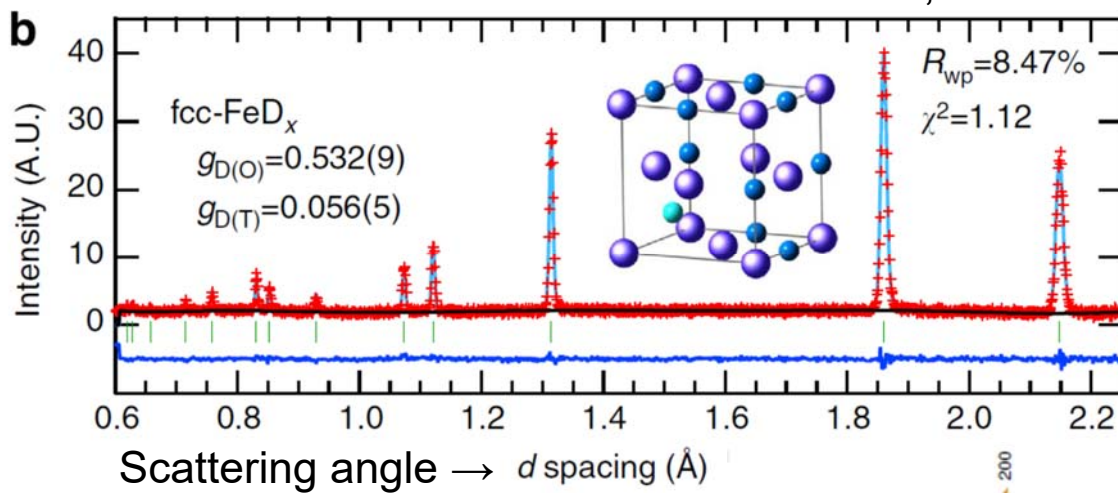
Independent of Z

→ suitable for observation of light atoms

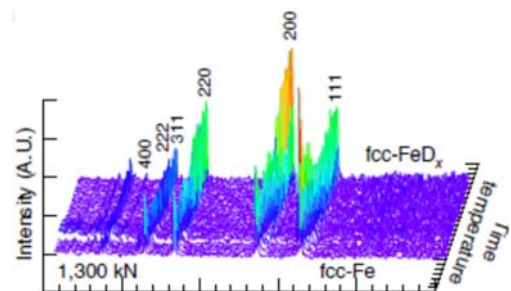
Neutron diffraction

Structure analysis of FeD_x

800 K, 7.4 GPa



Time evolution: operando

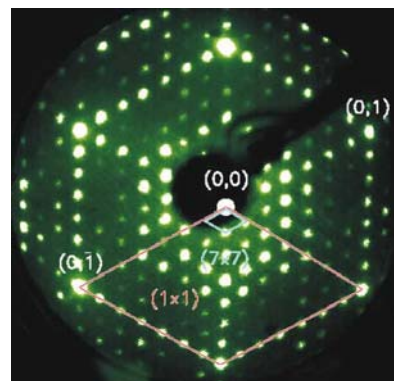


Electron diffraction

Surface structure

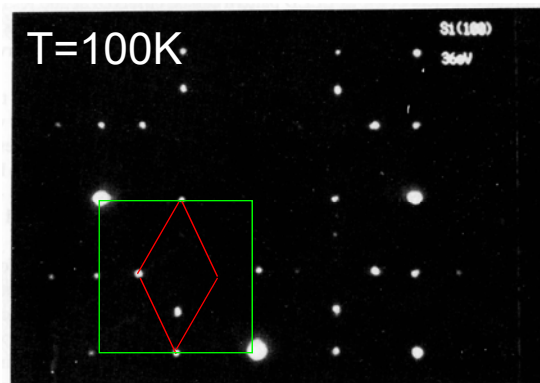
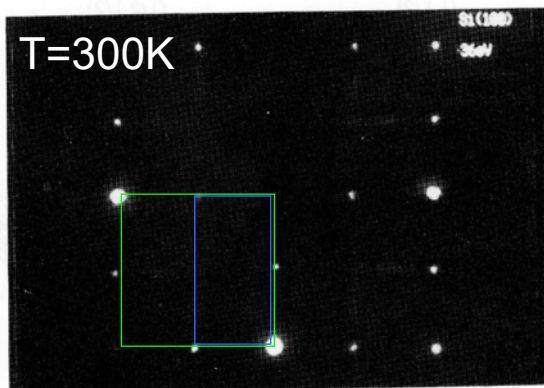
different from the bulk

Si(111)7x7



Phase transition

Si(001)



Phys. Rev. B 49(1994)4810.

Inelastic Scattering

$$A(E=E_0) \rightarrow \text{[Target]} \rightarrow A(E=E_0 - \Delta E)$$

Energy loss

Photon

Raman scatt.

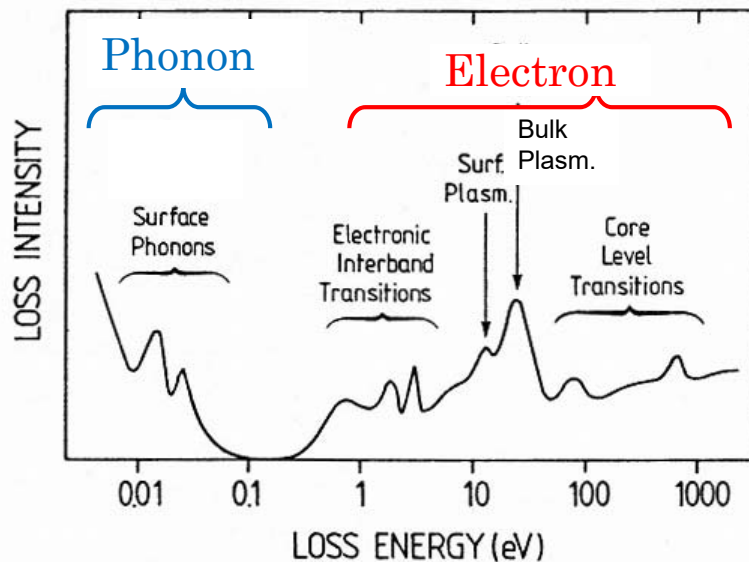
Compton scatt.

Photoemission

e^- , n

Electronic & phonon structures

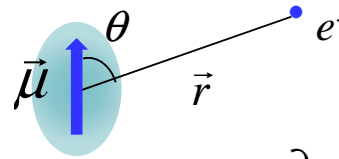
Energy scale



Electron Inelastic scattering due to phonon

Electron-dipole interaction

$$V \cong \frac{\partial}{\partial \vec{r}} \left(\frac{e}{4\pi\epsilon_0 r} \right) \cdot \vec{\mu} = \frac{e}{4\pi\epsilon_0} \frac{\vec{\mu} \cdot \vec{r}}{r^3}$$



Matrix element of V

$$\langle \psi_f | V | \psi_i \rangle \propto \langle \phi_v^f(R) | \hat{\mu}(R) | \phi_v^i(R) \rangle$$

$$\vec{E} = \frac{\partial}{\partial \vec{r}} \left(\frac{e}{4\pi\epsilon_0 r} \right)$$

\vec{R} : atom position

$$\hat{\mu}(R) = \hat{\mu}_0 + \frac{d\hat{\mu}}{dR} R + \quad \text{Taylor expansion}$$

$$\frac{d\hat{\mu}}{dR} \langle \phi_v^f(R) | R | \phi_v^i(R) \rangle \rightarrow \Delta v = \pm 1$$

(harmonic potential)

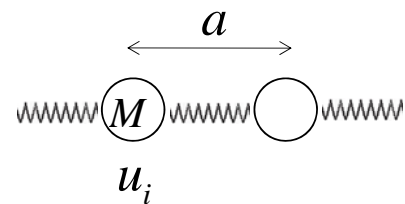
dynamic dipole

Cf. point charge approx. $\hat{\mu}(R) = qR$

Phonon: One-dimensional chain

$$V = \frac{1}{2} \sum_{i=-\infty}^{\infty} k(u_{i-1} - u_i)^2$$

$$M\ddot{u}_i = k(u_{i-1} - u_i) - k(u_i - u_{i+1})$$



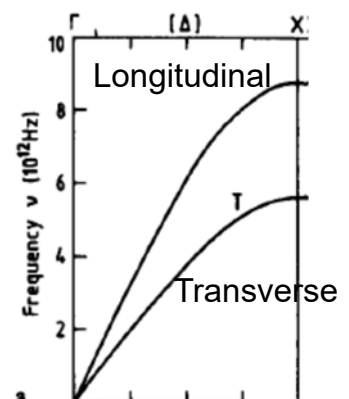
In the form of $u_i = u_i^0 \exp(i\omega t)$

$$\omega^2 M u_i^0 = k(u_{i-1}^0 - u_i^0) - k(u_i^0 - u_{i+1}^0)$$

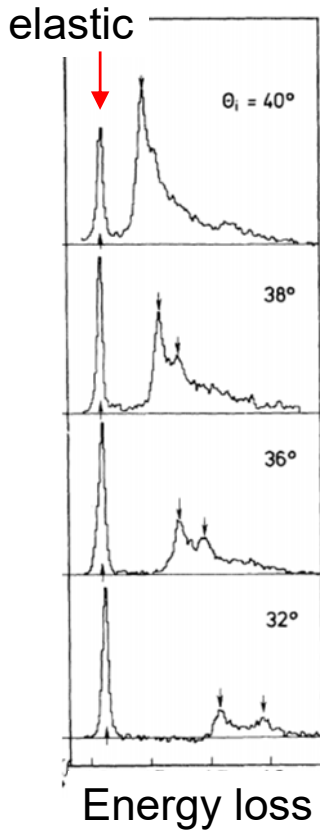
$u_n^0 = u^0 \exp(iq(na))$ Bloch theorem

$$-\omega^2 M u^0 = k(\exp(-iqa) - 2 + \exp(iqa))u^0$$

$$\omega = \sqrt{\frac{4k}{M}} \left| \sin \frac{qa}{2} \right|$$



Phonon dispersion: energy vs. momentum

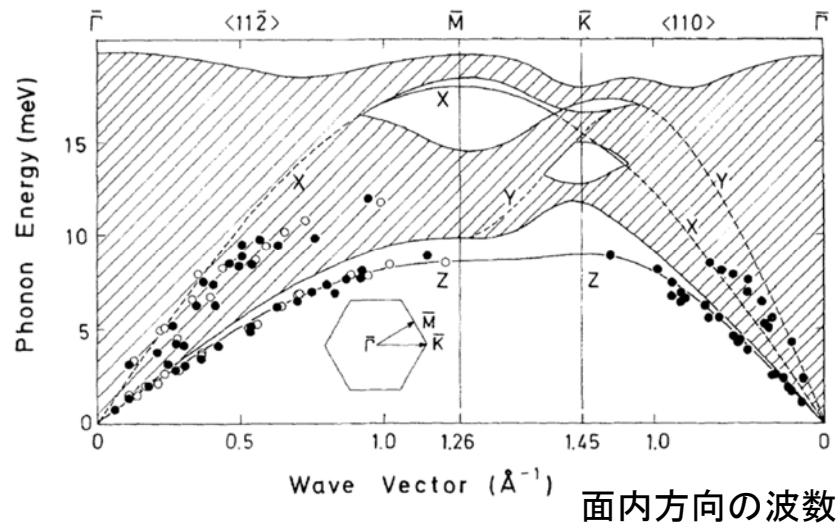


Ag(111)

Scattering angle \rightarrow momentum

$X : \langle 11\bar{2} \rangle$

$Y : \langle 110 \rangle$

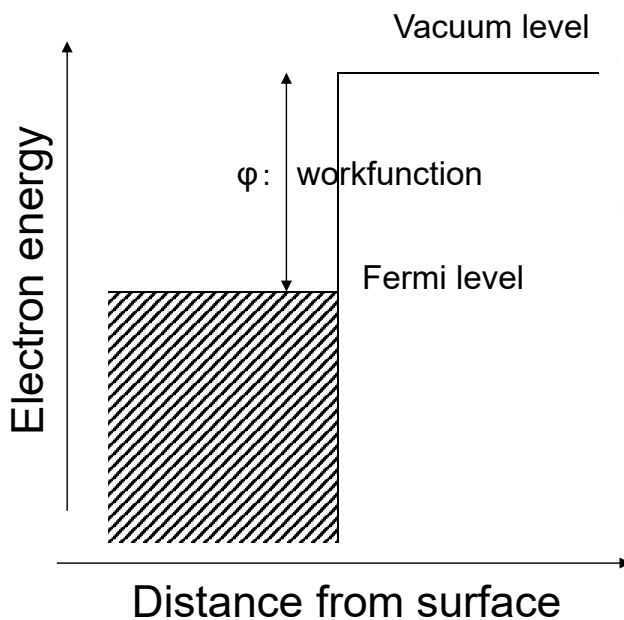


PhysRevLett.51.578

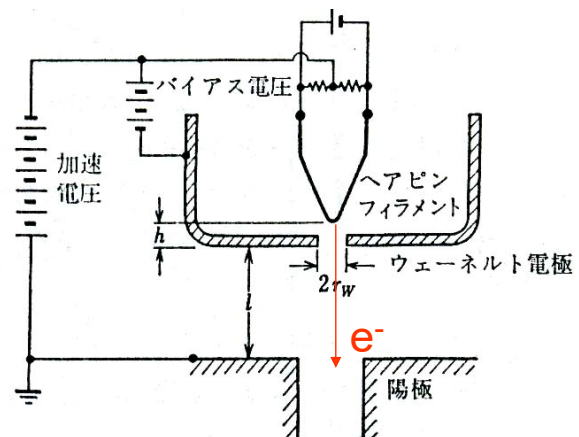
Electron emission

Thermionic electron emission

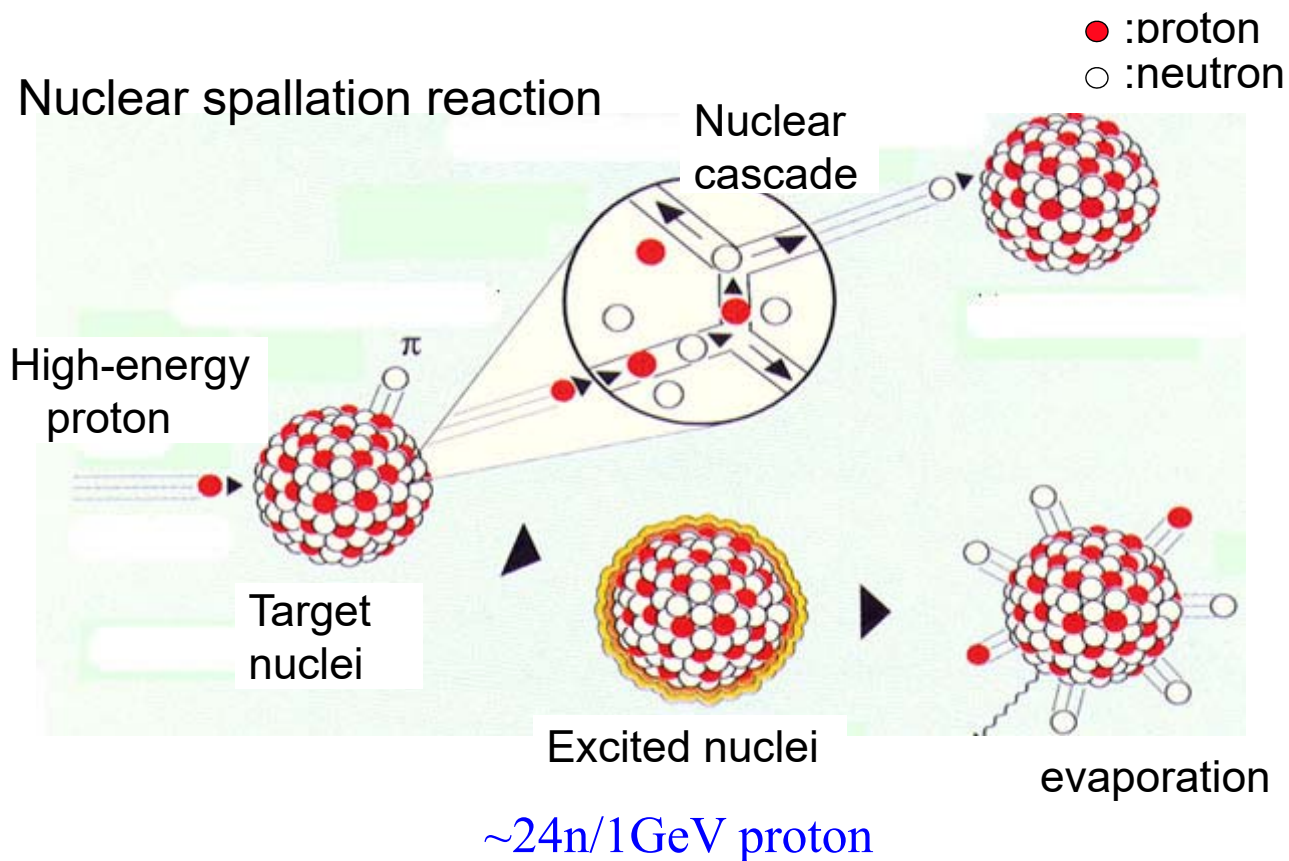
$$J = A_e T^2 \exp(-\phi / k_B T)$$



Electron gun

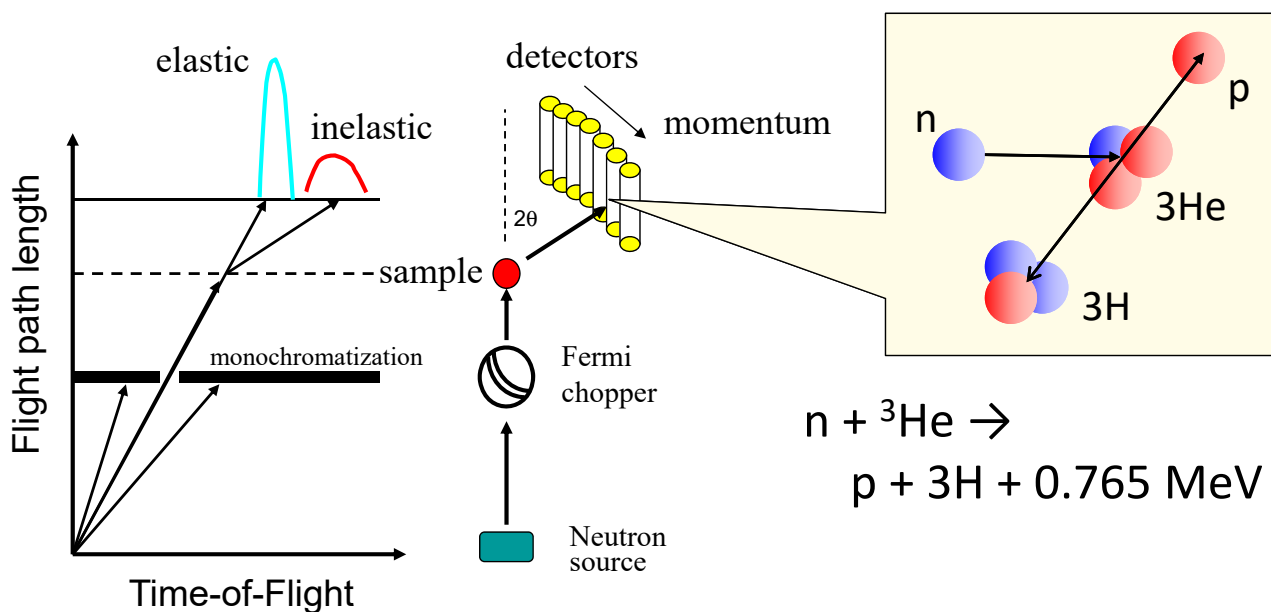


Production of neutron beam



Neutron inelastic measurement

- Neutron detector



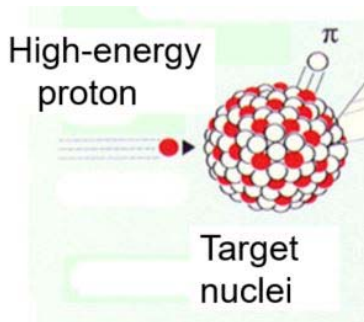
Muon (lepton): μ^+ , μ^-

- mass: $\sim 1/9$ of proton
- Spin: $1/2$
- Gyromagnetic ratio: $\gamma_\mu = 3.2 \gamma_p$
- Lifetime: $2.2 \mu\text{s}$

Magnetic moment

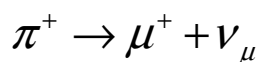
$$\vec{m} = \gamma_\mu \hbar \vec{s}$$

Production of Muon beam



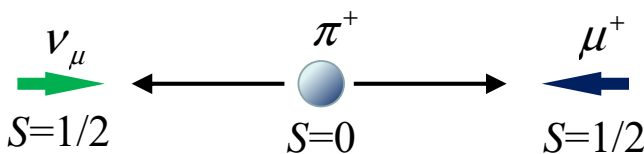
production of π

① μ^+ production

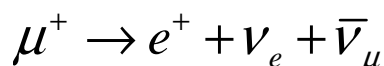


Spin: antiparallel to velocity

$\rightarrow \mu$ -spin: polarized

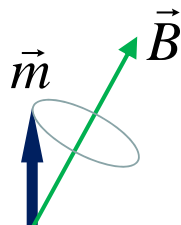


② μ^+ decay



in the spin direction

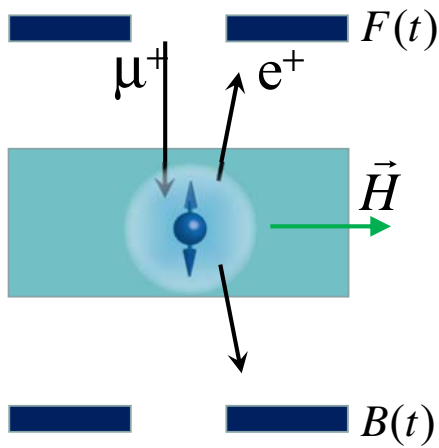
③ Precession under a magnetic field: \vec{B}



Precession frequency

$$\omega = \gamma_\mu B$$

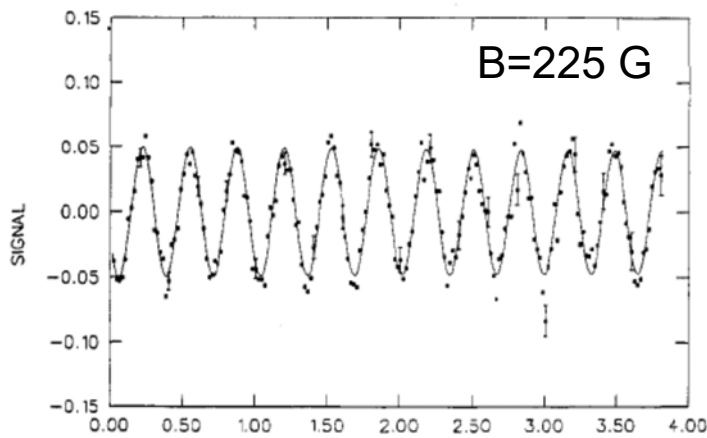
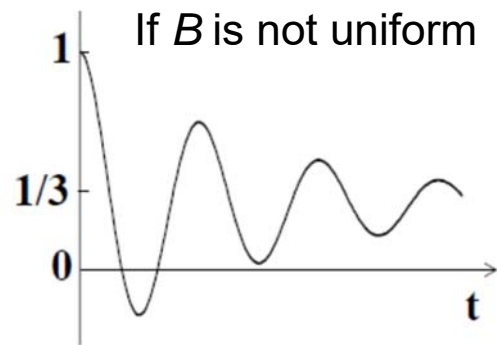
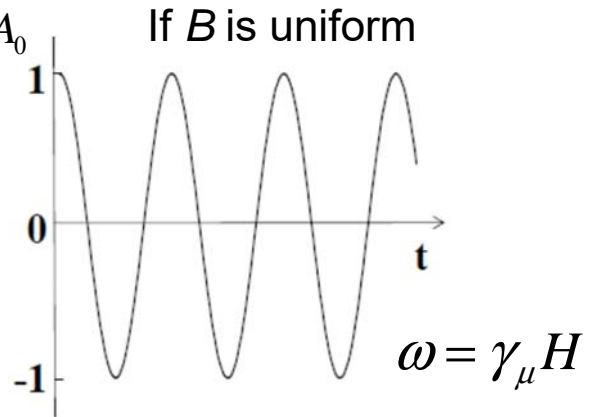
μ Spin Rotation/ relaxation



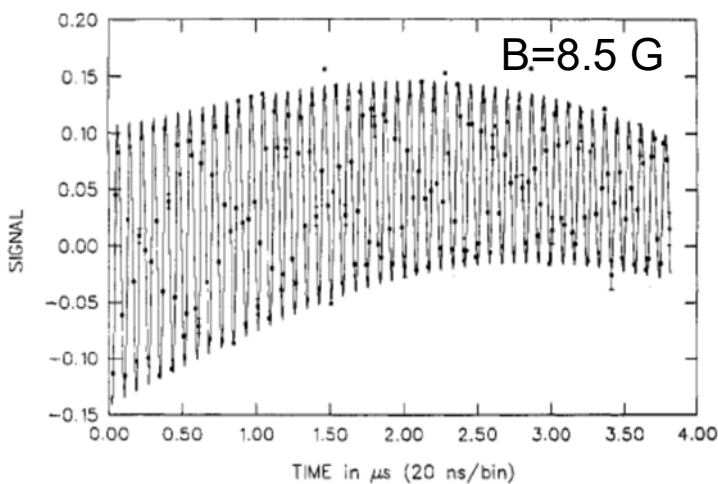
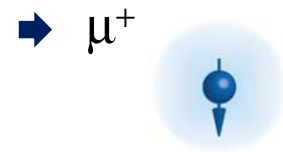
Asymmetry

$$A(t) = \frac{F - B}{F + B}$$

$A(t) / A_0$

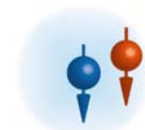


$$\gamma \sim 8.5 \times 10^4 \text{ rad/s/G}$$



Much faster oscillation

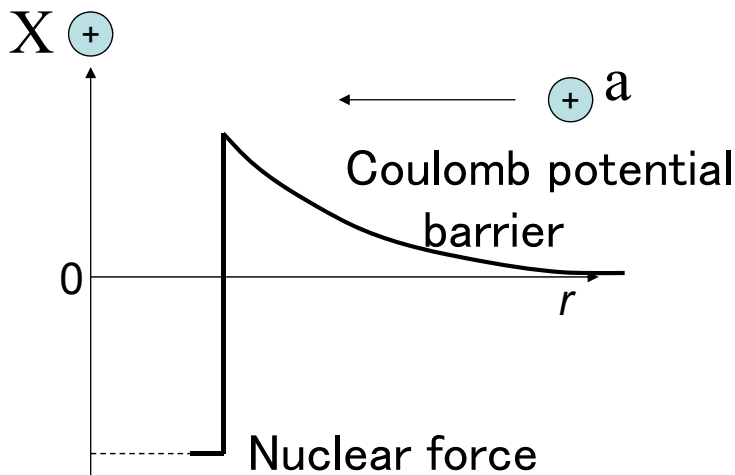
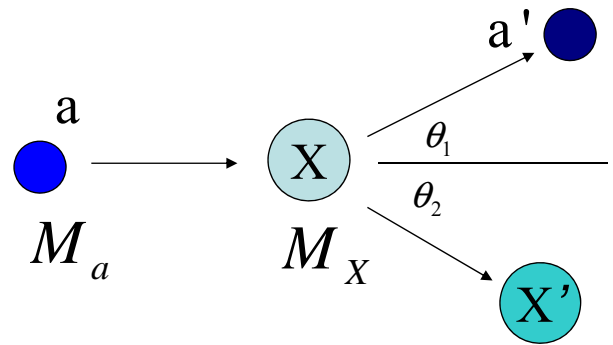
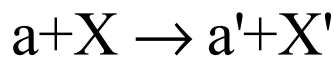
$$\gamma \sim 8.7 \times 10^6 \text{ rad/s/G}$$



$M(\mu^+e^-)$

Interaction of high-energy ions

Two-body scattering



E_a : low

$a' = a$: scattering

E_a : high

reaction

(light element)

Scattering of high-energy ions

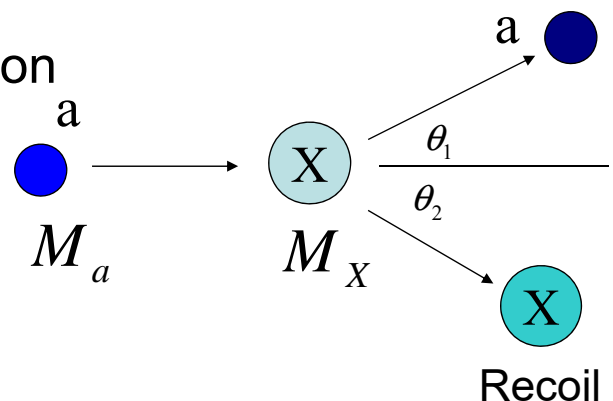
Two-body scattering

via Coulomb interaction

Energy after scattering

$$E' = kE_0$$

(depends on angle & mass)



$$k\text{-factor: } k \equiv \left(\frac{v'}{v}\right)^2 = 1 - \frac{4M_a M_X \cos^2 \theta_2}{(M_a + M_X)^2}$$

Energy analysis \rightarrow mass of X = element

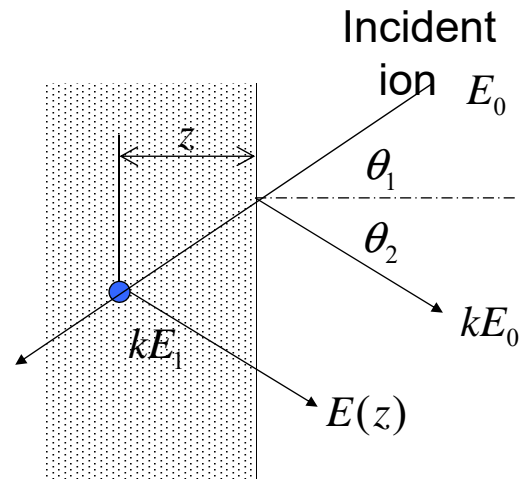
環境 (電子状態) に依らない

Energy loss of ions

During passing through a solid

Ions gradually lose energy

Stopping power $S = \frac{dE}{dz}$



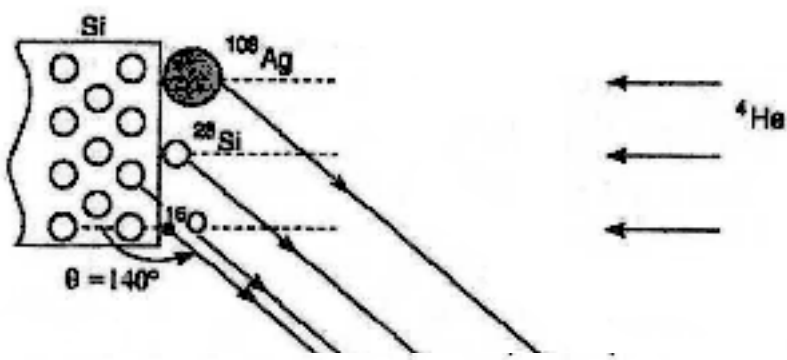
Energy analysis

→ penetration length = depth

$$E(z) = kE_0 - z \left(\frac{kS_{in}}{\cos \theta_1} + \frac{S_{out}}{\cos \theta_2} \right)$$

↑ Incidence path ↑ Exit path

Rutherford Back Scattering (RBS)

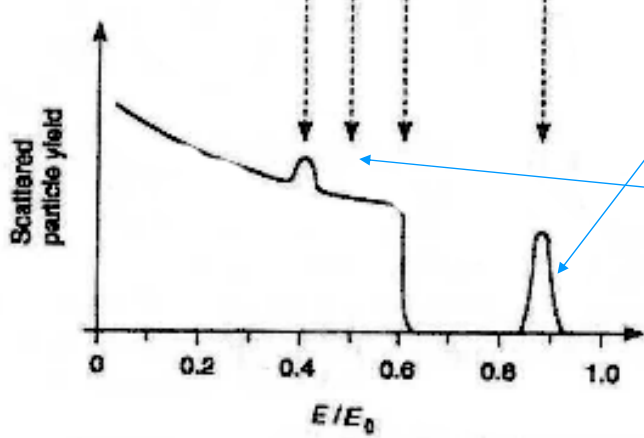


① Scattering from a heavy element

↓
Large k -factor
→ large energy

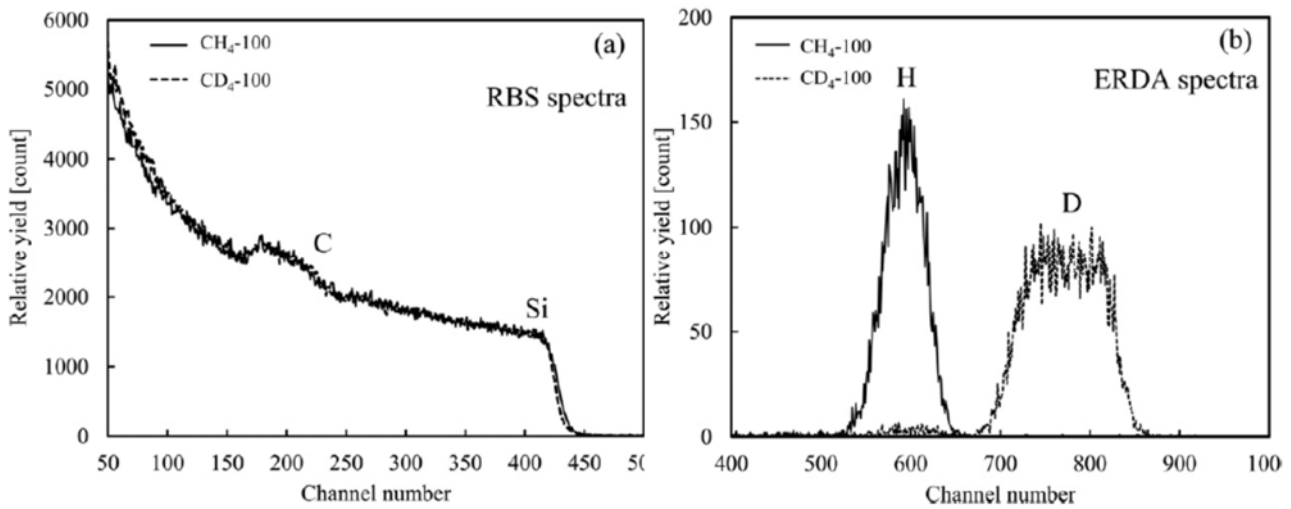
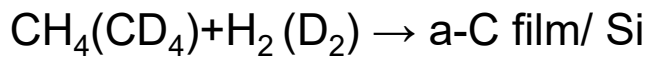
② Scattering from a deep place

↓
Low energy



Suited to heavy-element analysis

Elastic Recoil Detection (ERDA)



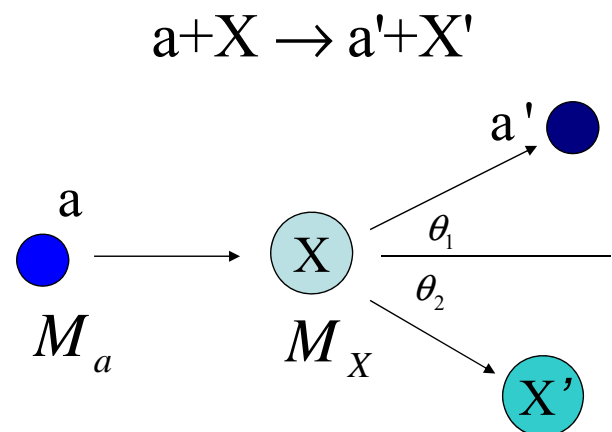
K. Ozeki *et al*, Appl. Surf. Sci. 265 (2013) 750.

Nuclear Reaction Analysis (NRA)

Nuclear reaction

$a \neq a', X \neq X'$

- element-specific
- quantitative



Coulomb barrier due to nuclei

Suited to light elements

- ✓ Hydrogen
- ✓ Lithium

Recent progress

1 . Improvement of resolution & sensitivity

Depth resolution ~ nm

Multi-element analysis: e.g. distinction of N and O

Background reduction: improvement of sensitivity

2 . Operando analysis

Materials characterization

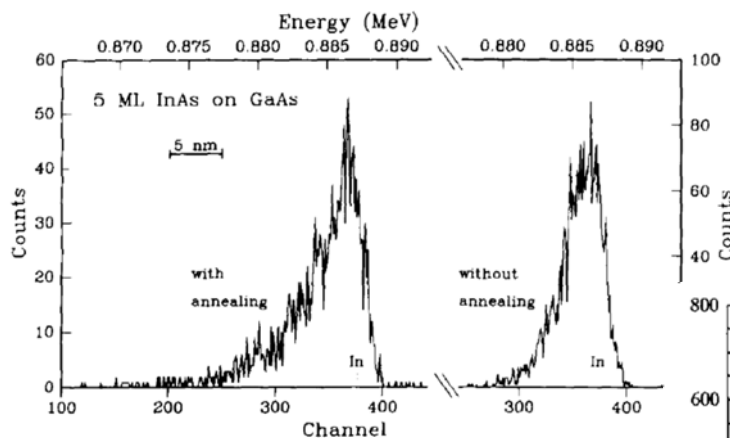
under device operation/ reaction

e.g. gas atmosphere, liquid interface

✓ Quantum beams: high penetration in materials

High-resolution RBS

High-resolution energy analysis: nanometer scale analysis



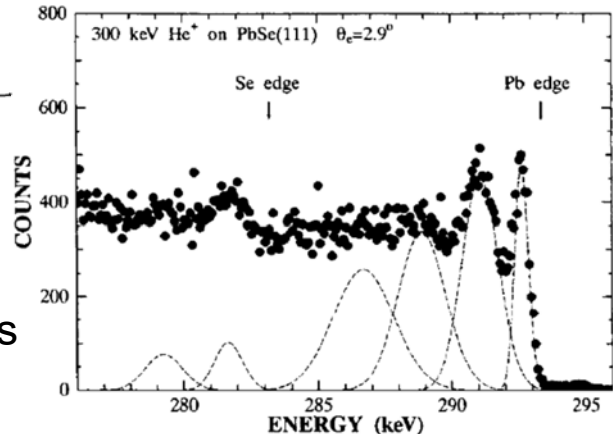
NIMB 64, 817 (1992)

Energy
distinction of depth and mass

NIMB 113, 270 (1996)

SSD → magnetic analyzer

Depth resolution ~ 1/10

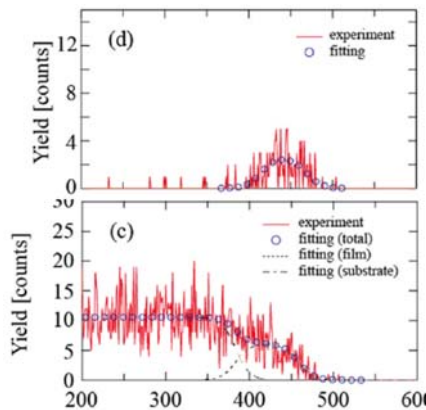
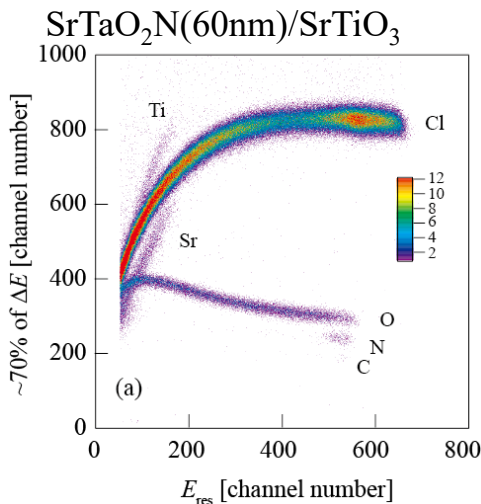
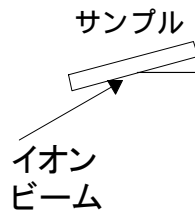


$\Delta E - E$ telescope ERDA: depth and mass

Energy loss in gas



Element identification



Nの深さ分布

Oの深さ分布

I. Harayama, D. Sekiba *et al.*: NIMB 384 (2016) 61.

Application to Archeology & Art

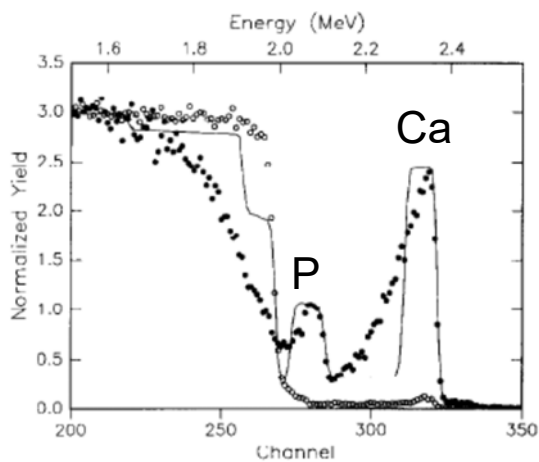
Louvre Museum: AGLAE accelerator

NIMB 64 (1992) 488.

Usewear characterisation of prehistoric flints with IBA

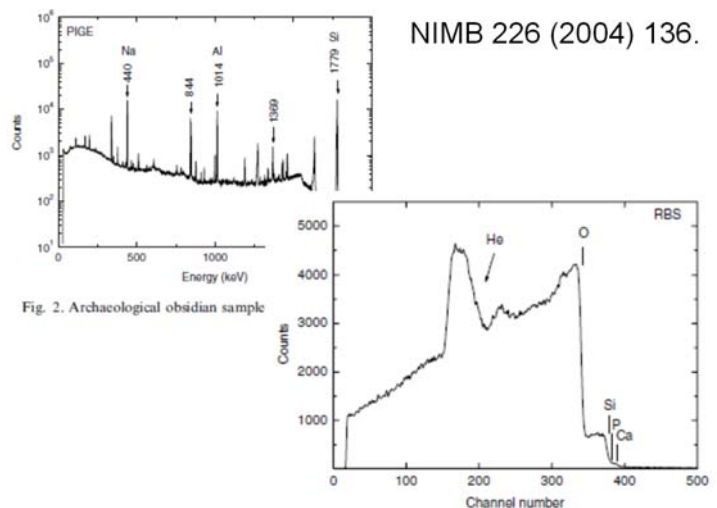
M. Christensen, Ph. Walter and M. Menu

Laboratoire de Recherche des Musées de France, Palais du Louvre, 75041 Paris Cedex 1, France

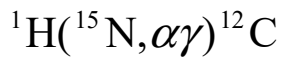


Rosendorf, Germany

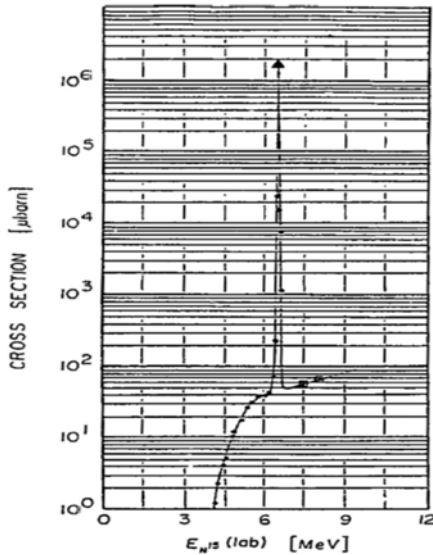
NIMB 226 (2004) 136.



H detection

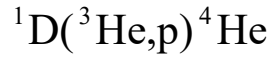


narrow
width=1.8 keV

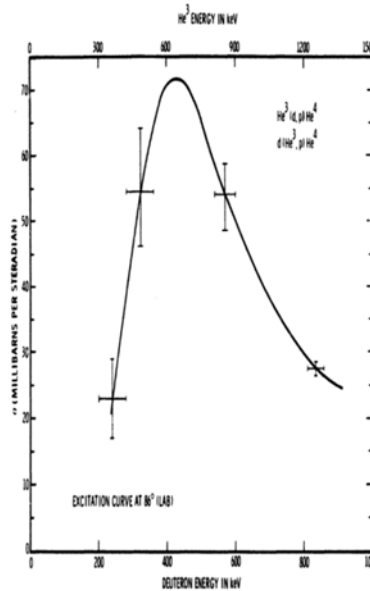


NIMB_66_65 (1992)

D detection

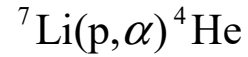


broad
width~500 keV

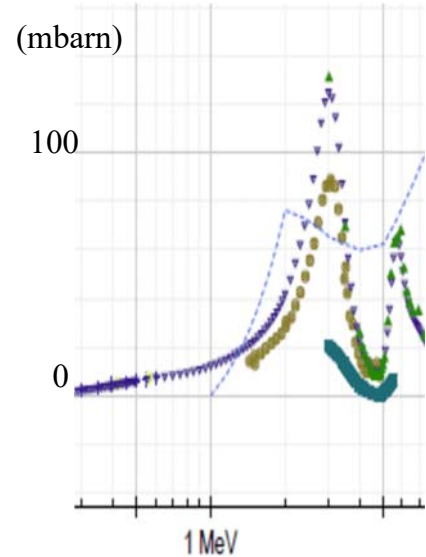


JNuclMater_53_257 (1974)

Li detection

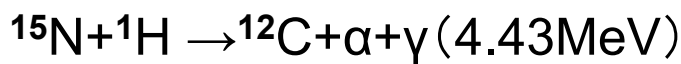


broad
width~1 MeV



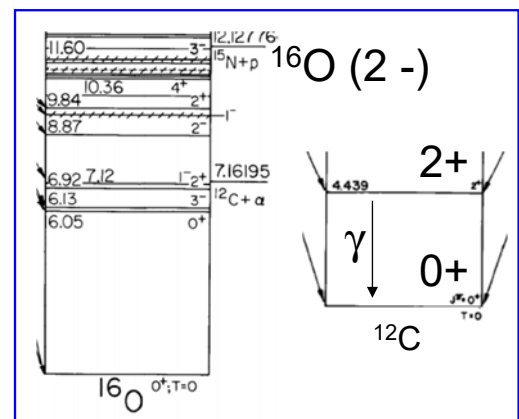
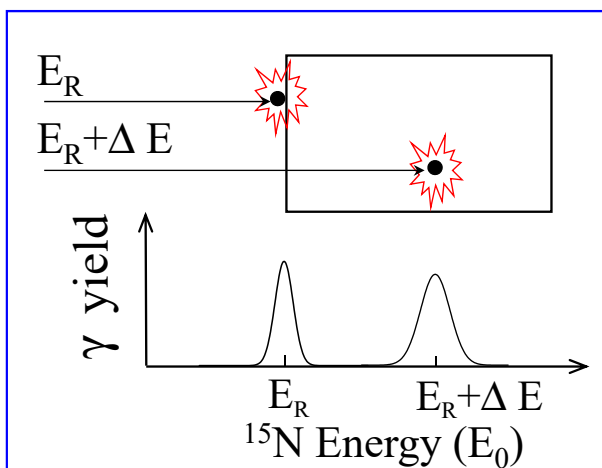
NIMB_66_107 (1992)

H detection with NRA



Resonance at $E_R = 6.385$ MeV

width= 1.8 keV



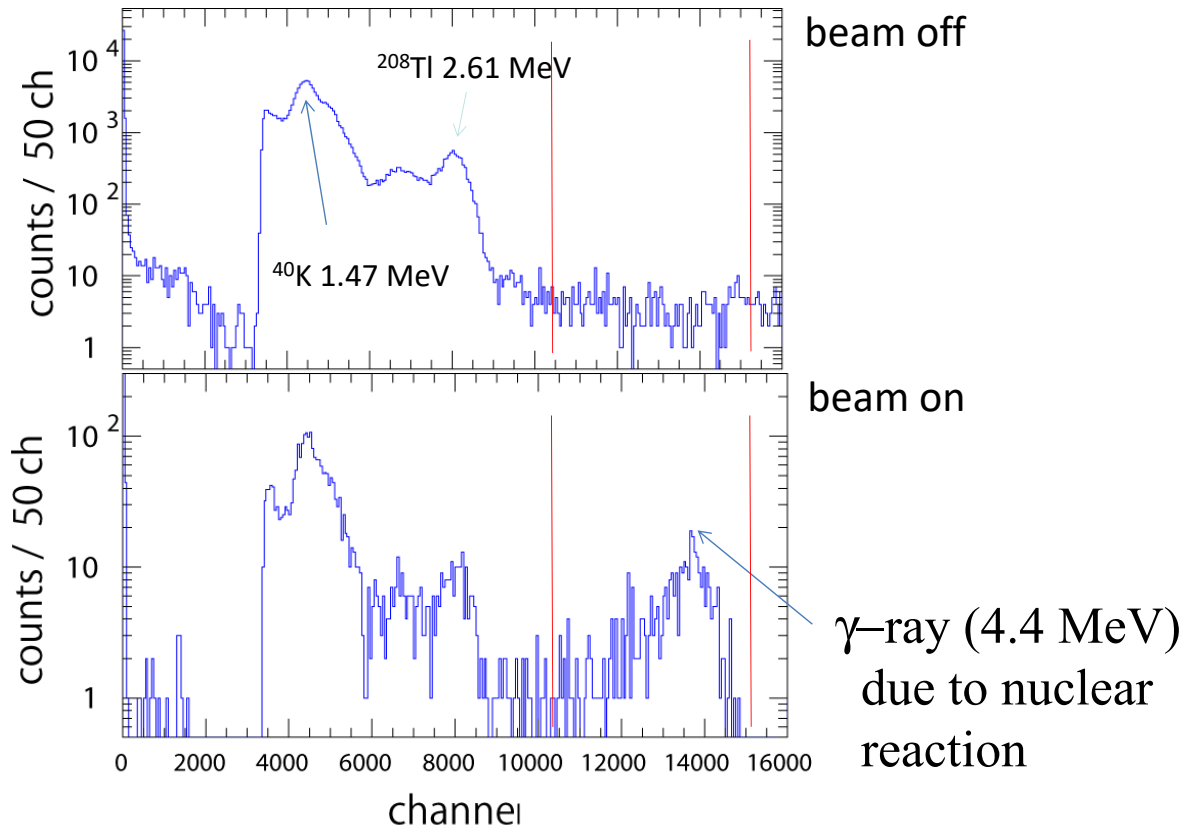
γ -yield vs. ${}^{15}\text{N}$ energy

H depth profile

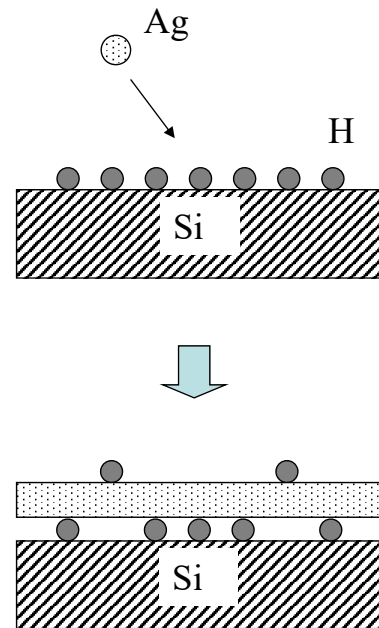
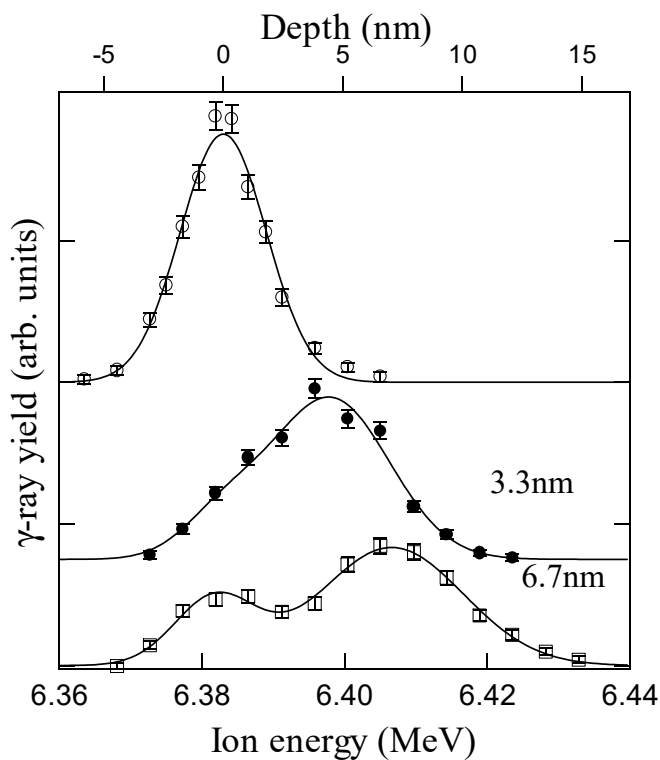
depth resolution

$\sim \text{width}/(dE/dz)$

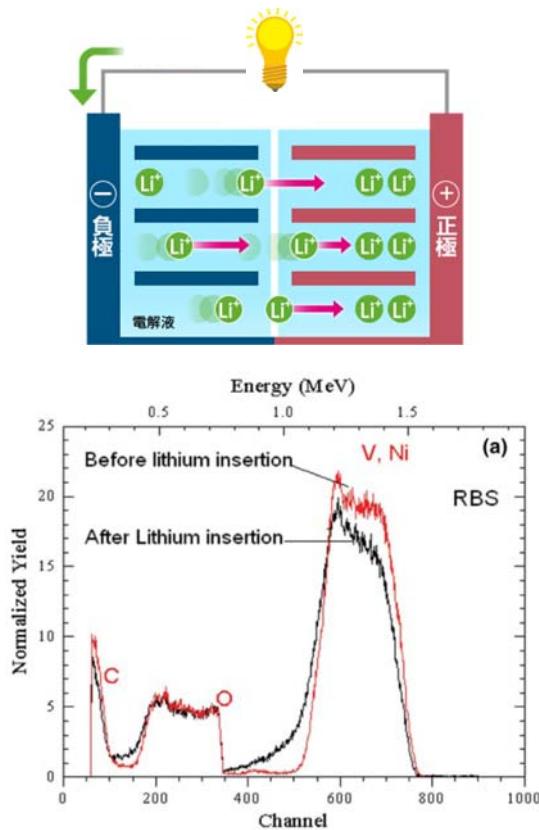
γ -ray spectrum



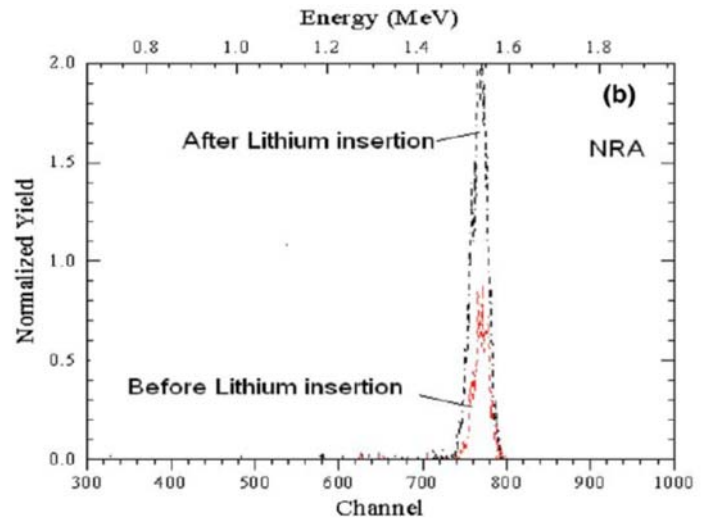
Depth distinction of H



Li ion battery: transport of Li



LiNiVO₄ (300 nm)/C cathode



⁴He (2 MeV) RBS

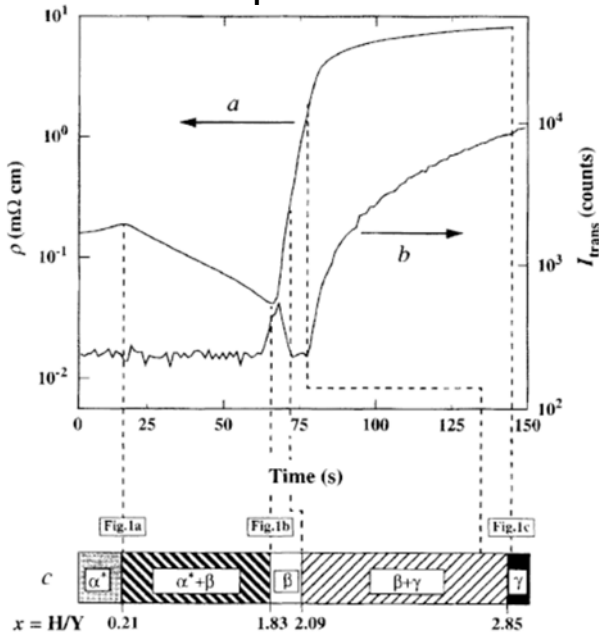
M.V. Reddy *et al.*, NIMB 246 (2006) 397

Hydrogen in solids: dual nature

- * Electronic effects on host properties
 - H-induced metal-insulator transition
 - Relation with electric-device performance
- * Proton motion
 - proton diffusion
 - hydrogenation reaction: photocatalysis
 - electronic to chemical energy conversion

H-induced metal insulator transition

Thin Y films (500 nm):
exposure to H



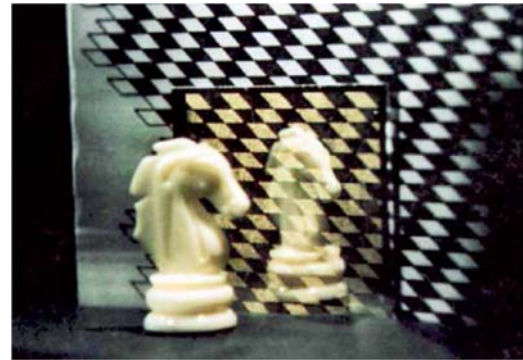
Metal → Insulator

(J. N. Huiberts et al., Nature 380, 231 (1996))

Mirror

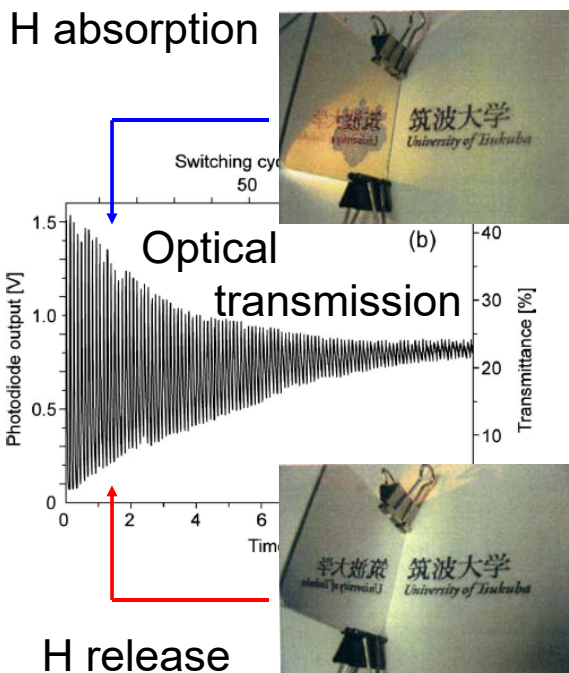


Transparent

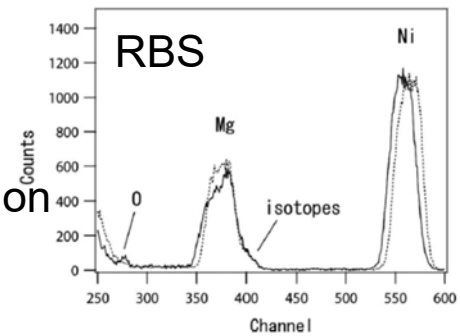


Degradation mechanism: MgNi hydride

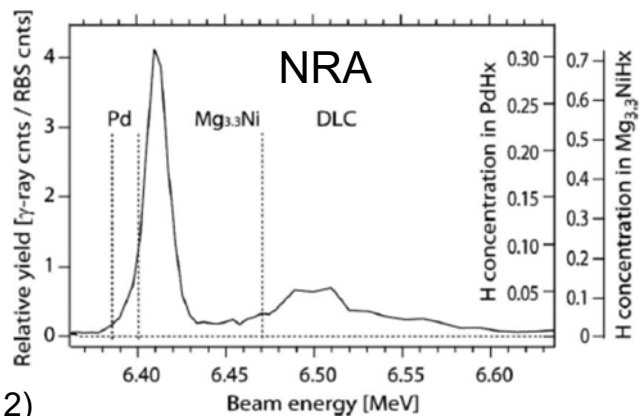
H absorption



Mg segregation



MgH₂, MgO formation



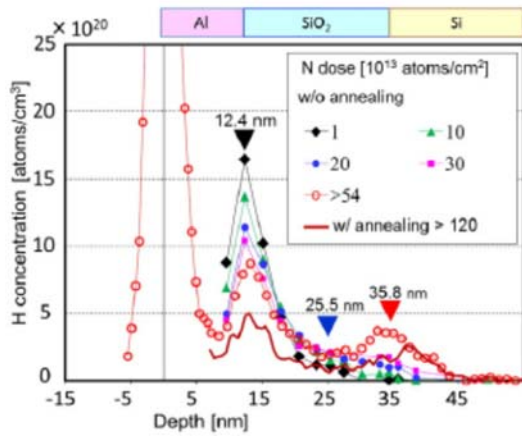
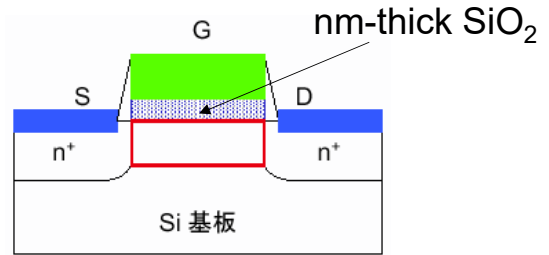
(D. Sekiba et al., JAP 106, 114912)

H in Si devices

Si devices

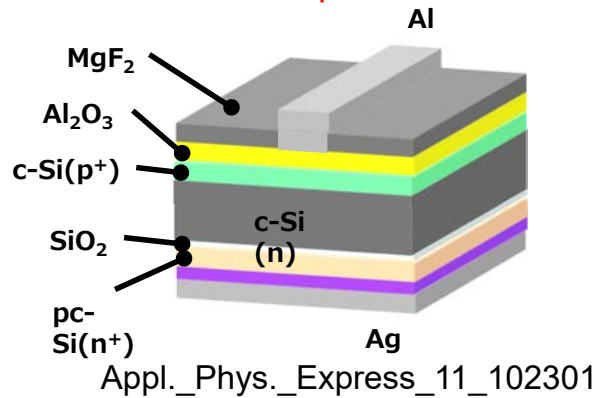
- MOS transistor
- Photo/ particle detector
- Solar cell

MOS transistor



MicroeleRel_70_12

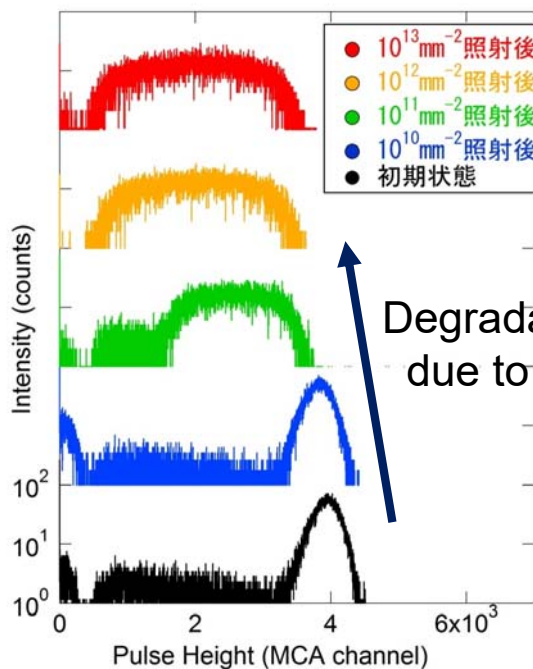
TOPCon solar cells $\eta = 25.3\%$



Appl._Phys._Express_11_102301

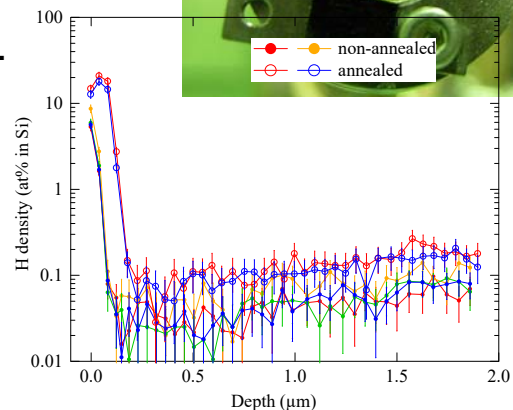
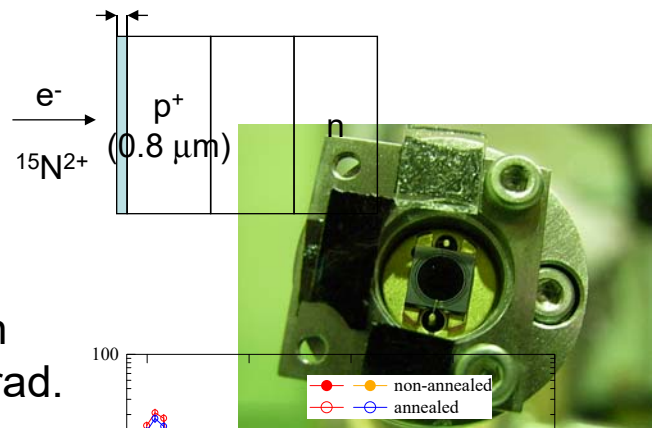
Degradation mechanism of Si detector

Pulse height analysis



IEEE EDL 33, 1162 (2012)

SiO₂(65 nm)



Hydrogen absorption in Pd

H₂ gas: large volume

Store H in metal, and extract when necessary

- How much H is in Pd?
- How stable is H in Pd?
- How fast is H absorption into Pd?
- Does absorbed H participate in reactions?

Hydrogen absorption:

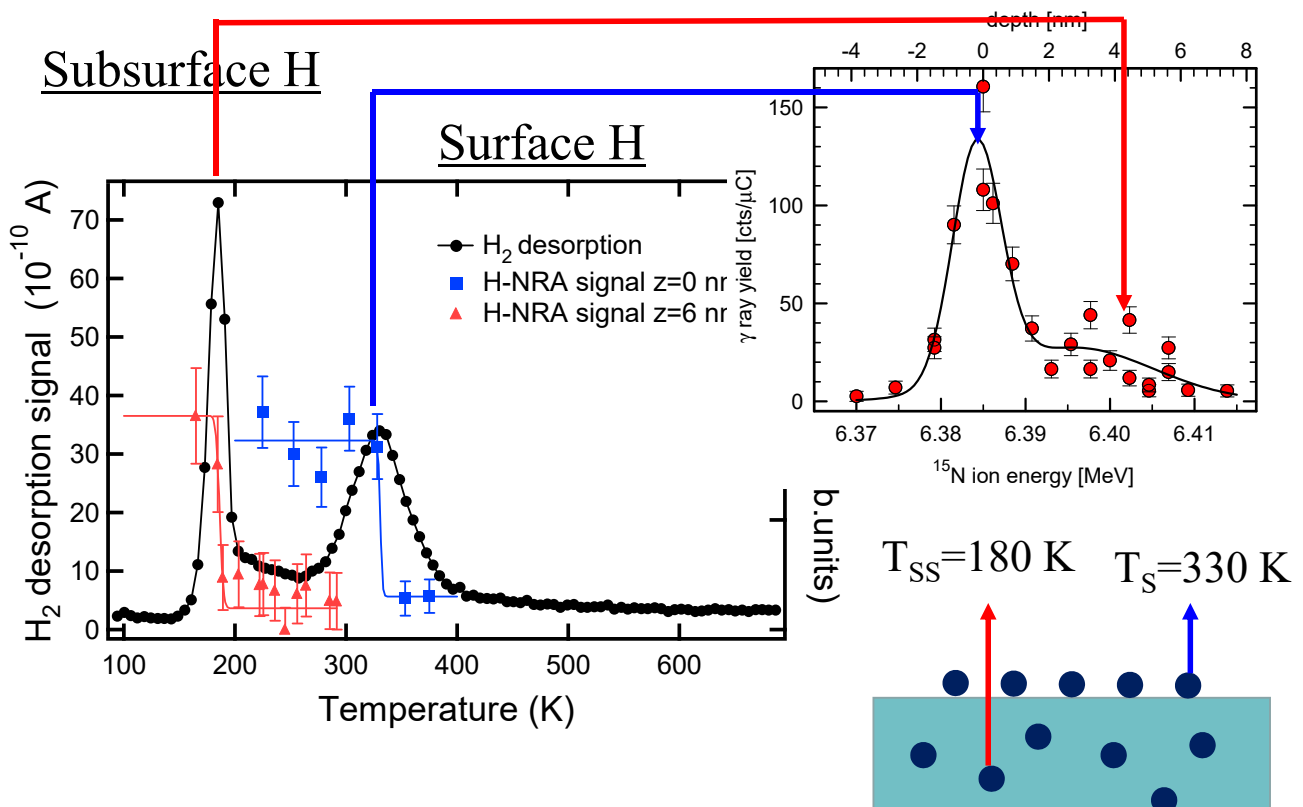
transport of H via surfaces

- Absorption properties
- Surface properties

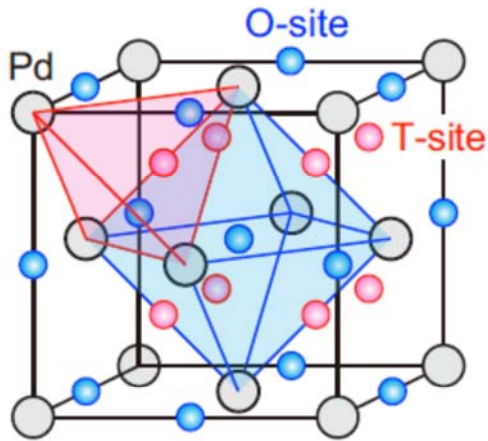


- Nano particles
- Surface modification

Absorption & release of H in Pd

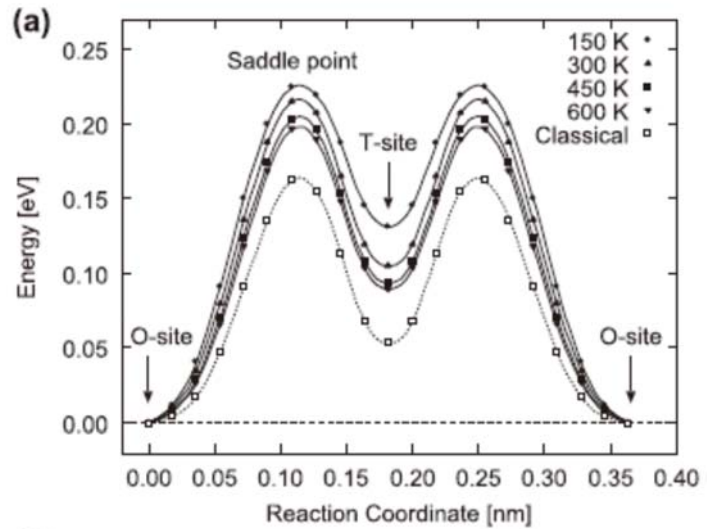


H in interstitials of Pd



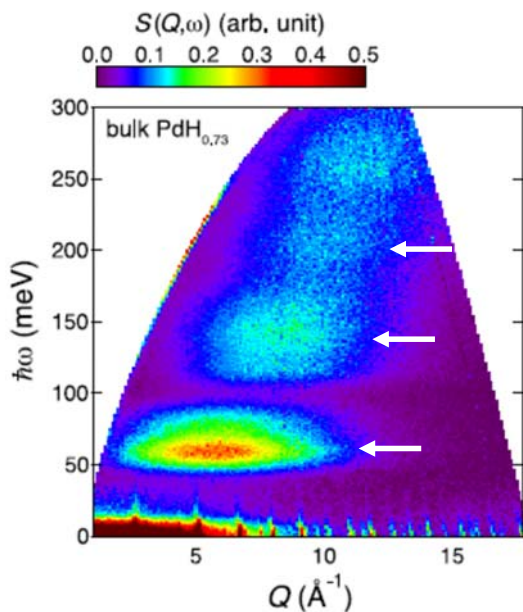
- O-site: most stable
- T-site: metastable

Calculated potential



Looks harmonic

Neutron inelastic scattering: H in Pd



Almost equally spaced peaks

Vibration of H

Z. Phys B 55, 283 (1984); PhysRevB.96.054304

3D harmonic oscillator

$$V(x, y, z) = C(x^2 + y^2 + z^2)$$

$$E = \sum_{i=x,y,z} \hbar\omega(n_i + \frac{1}{2}) \quad |n_x n_y n_z\rangle$$

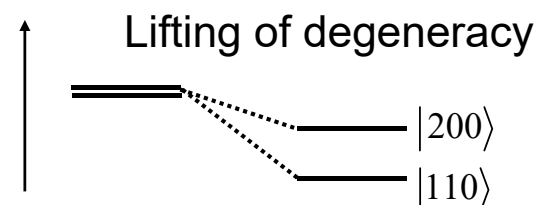
anharmonic term

$$V' = C'(x^2 y^2 + y^2 z^2 + z^2 x^2)$$

$$x = \sqrt{\frac{\hbar}{2m\omega}} (a_x^\dagger + a_x)$$

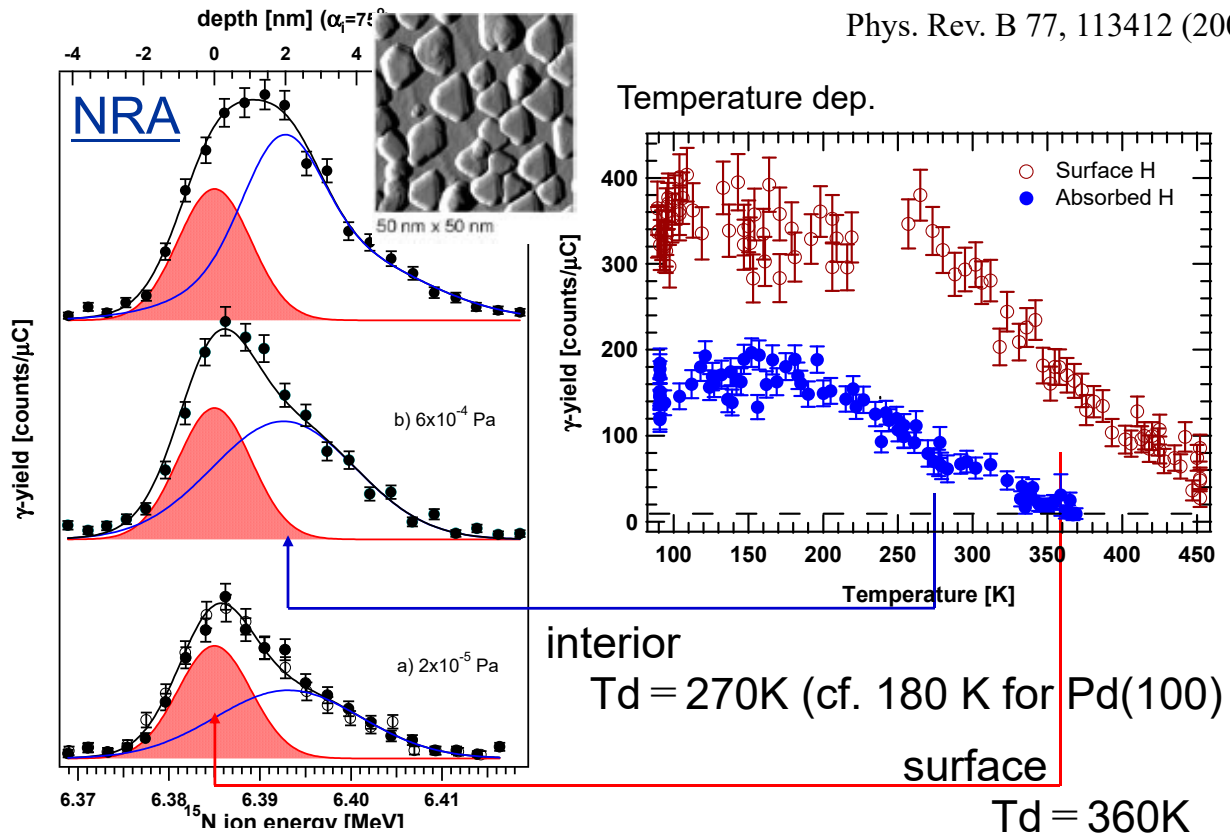
perturbation

$$\langle 200 | V' | 200 \rangle \neq \langle 110 | V' | 110 \rangle$$

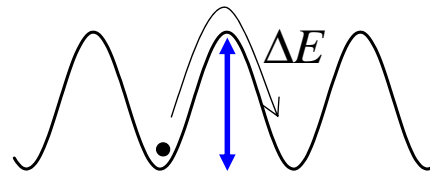
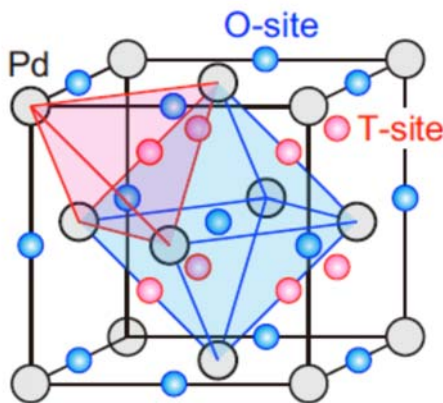


H absorption in Pd nano particles

Phys. Rev. B 77, 113412 (2008)



Diffusion of a particle



H undergoes hopping
from one site to neighboring sites

$$D = \Gamma \frac{\ell^2}{2d_s}$$

Γ : Jump rate

ℓ^2 : Mean sq. jump distance

d_s : dimension

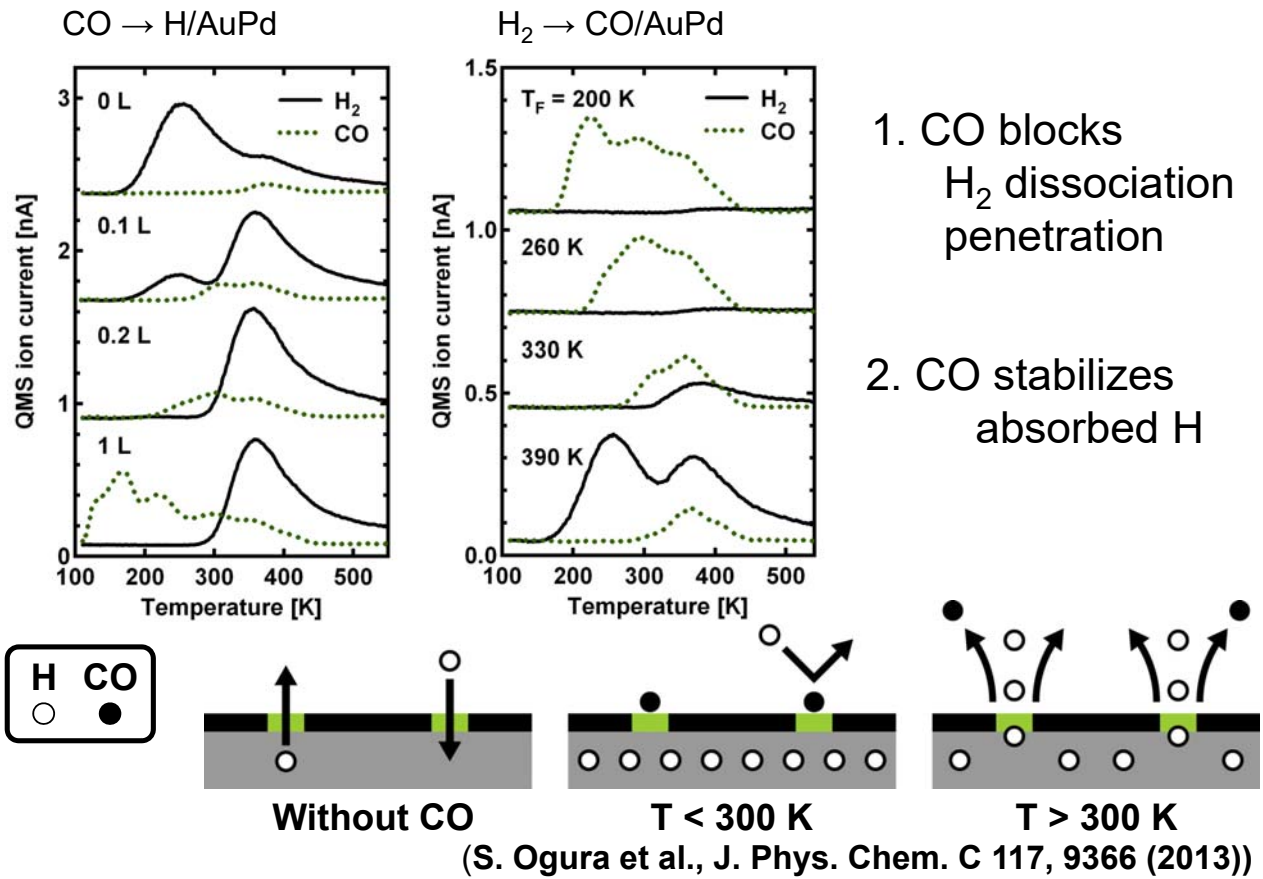
Diffusion equation

$$\frac{\partial c(\vec{r}, t)}{\partial t} = D \nabla^2 c(\vec{r}, t)$$

Thermal diffusion

$$D = D_0 \exp(-\Delta E/kT)$$

Molecular cap by CO



Stability & kinetics

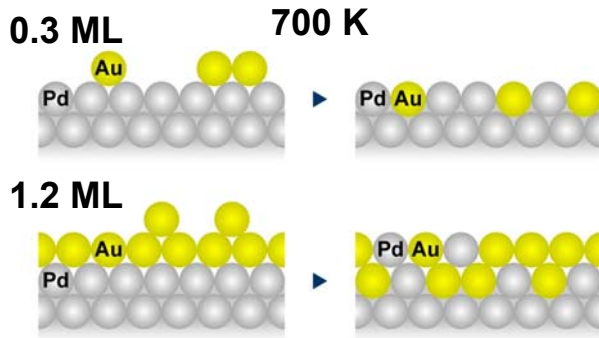
| | desorption T | | abs. prob. |
|---------|--------------|----------|-------------------------|
| | surface | subsurf. | |
| Pd | 320 K | 180 K | $\sim 7 \times 10^{-4}$ |
| AuPd | — | 250 K | $\sim 3 \times 10^{-2}$ |
| Nano-Pd | 360 K | 270 K | |
| CO-cap | | 350 K | |

Compared with Pd

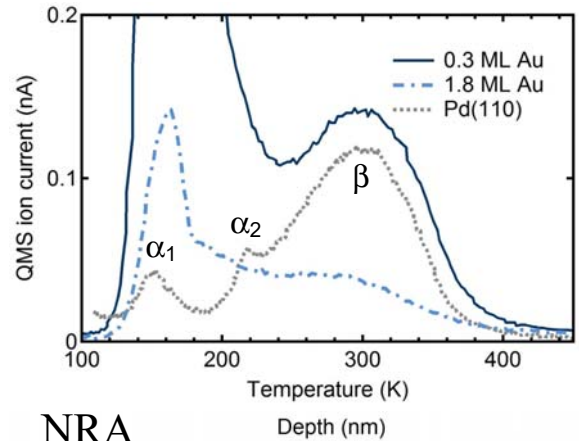
AuPd: lattice expansion

Pd nano particle: lattice deformation

Surface modification by Au



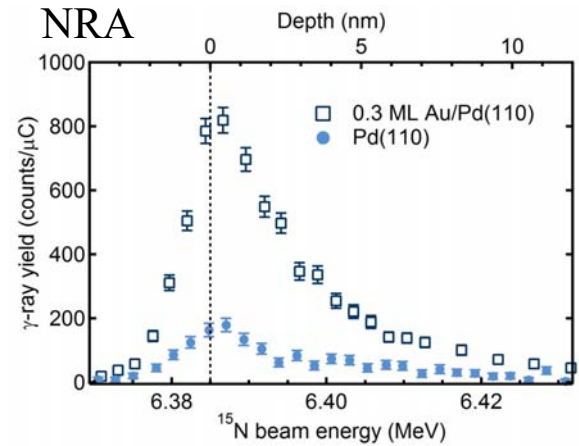
TDS



α_1 : absorbed H
increase
 β : surface-adsorbed H
decrease

K. Namba et al., PNAS 115, 7896 (2018)

NRA



Enhancement of H absorption by Au

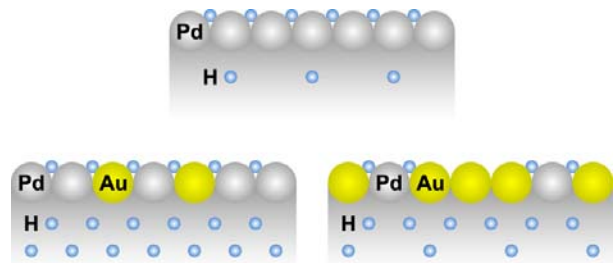
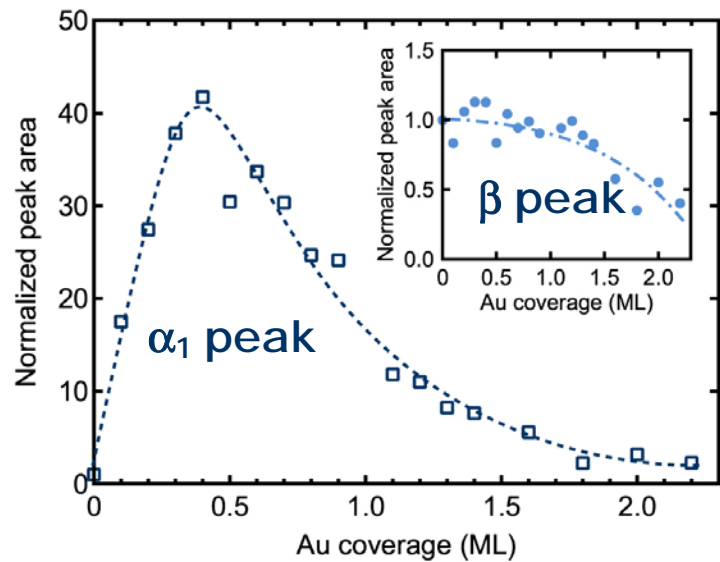
Absorption probability
Max at $\theta_{Au} = 40\%$

Decrease of Surf. H
with increasing Au

“Conventional idea”

Au: chemically inert

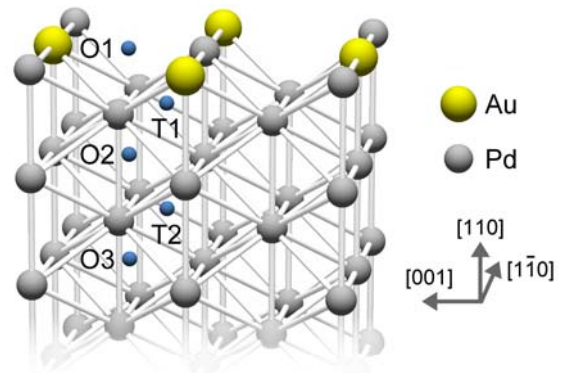
→ useless for H absorption



Role of Au for H absorption

First-principles total-energy cal.

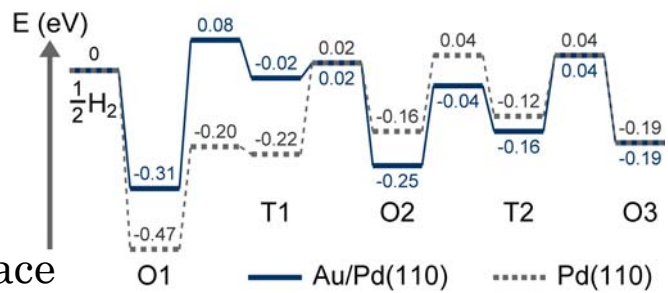
- (1) Adsorption on surface
- (2) Diffusion into subsurface
- (3) Diffusion to bulk



Where is the bottleneck?

Surface destabilization
→ enhancement of (2)

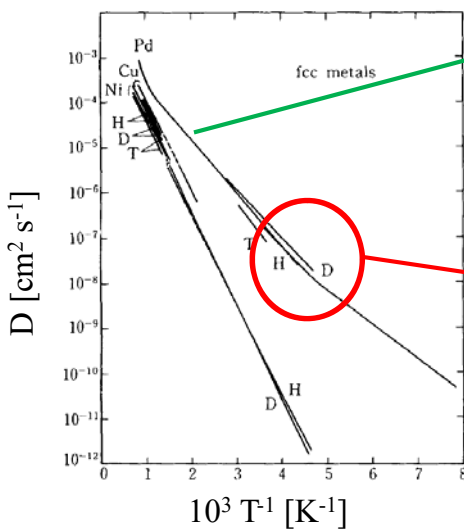
Too much Au
→ no adsorption on surface



K. Namba et al., PNAS 115, 7896 (2018)

Quantum effects in diffusion

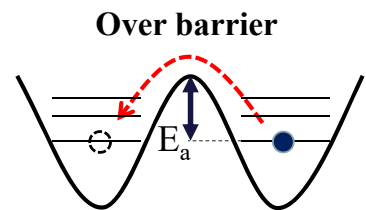
Diffusion coefficient
Temp. dep.



High T

① Thermal diffusion

$$D \propto \exp(-E_a/kT)$$



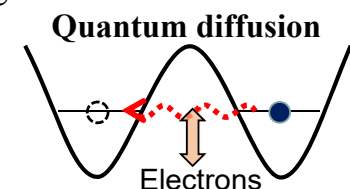
② D is faster than H

Inverse isotope effect

Low T

③ Quantum tunneling?

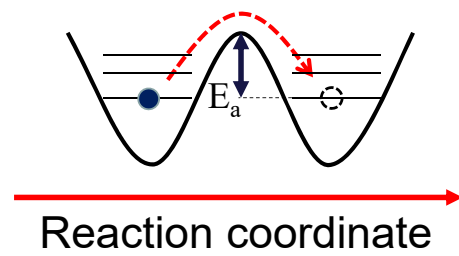
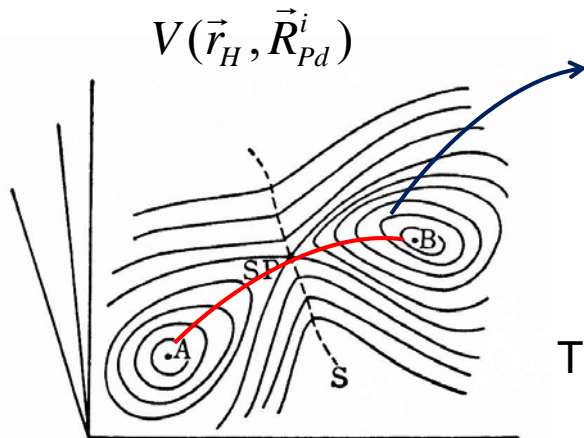
$$D \sim \text{const.}$$



Y. Fukai, The Metal-Hydrogen System, Springer Series in Materials Science (2005).

Origin of inverse isotope effect

Multi-dimensional potential

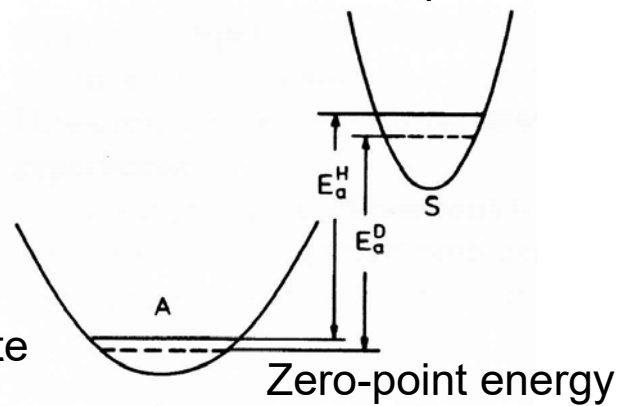


Transition state = saddle point

If confinement at SP is strong,

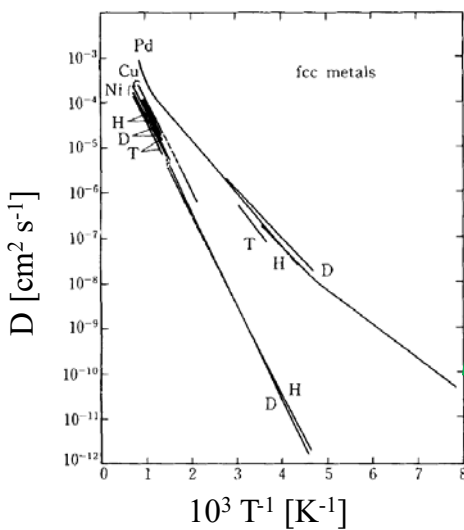
$$E_a^H > E_a^D$$

Initial state



Quantum effects in diffusion

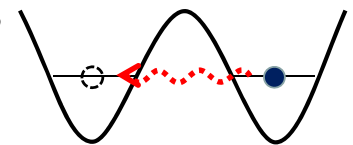
Diffusion coefficient
Temp. dep.



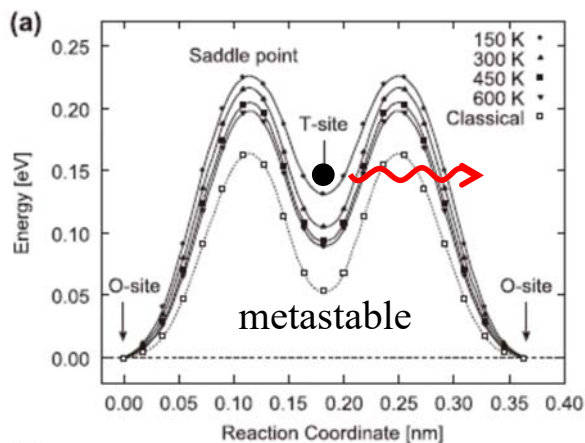
Low T: no data

③ Quantum tunneling?

$$D \sim \text{const.}$$

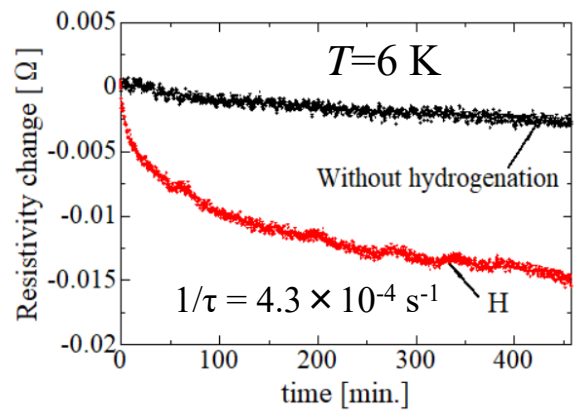
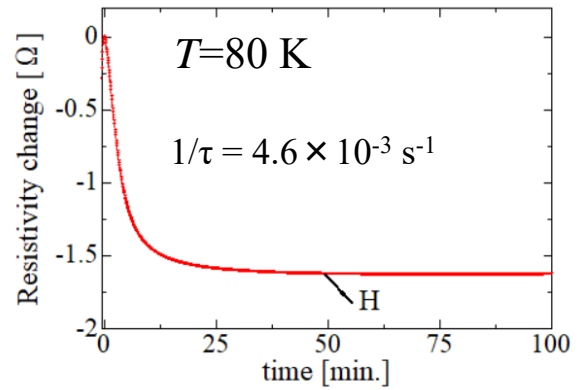
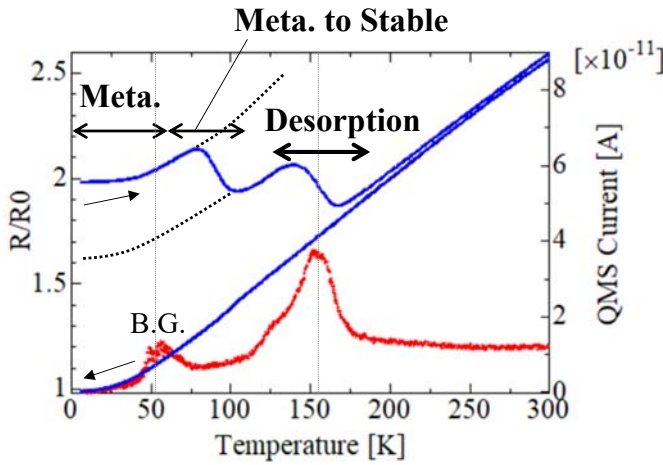


Our recent strategy: T to O diffusion



Y. Fukai, The Metal-Hydrogen System, Springer Series in Materials Science (2005).

Probing H diffusion by resistance



H ion implantation
→ T-site occupation

Resistance change
by H diffusion from T to O

T. Ozawa et al., Vac. Surf. Sci. in press.

Temperature dependence of H/D hopping rate

$$\exp(-E_a/kT)$$

① Thermal diffusion above 80 K

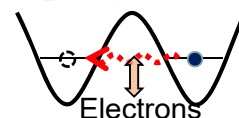
② Normal isotope effect

③ small T dependence

Isotope dependence

Proton tunneling

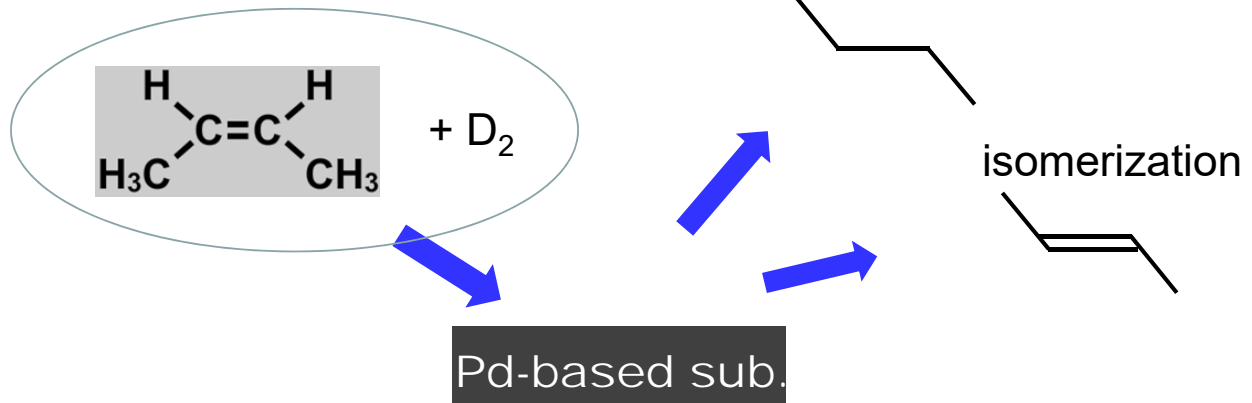
Coupling with electrons



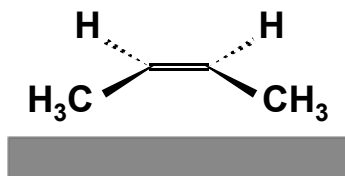
Cf. J. Kondo, Physica 125B (1984) 279.

Reaction: hydrogenation & isomerization

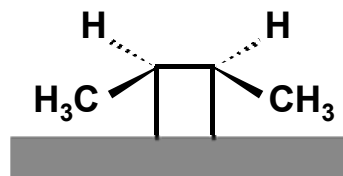
Model reaction



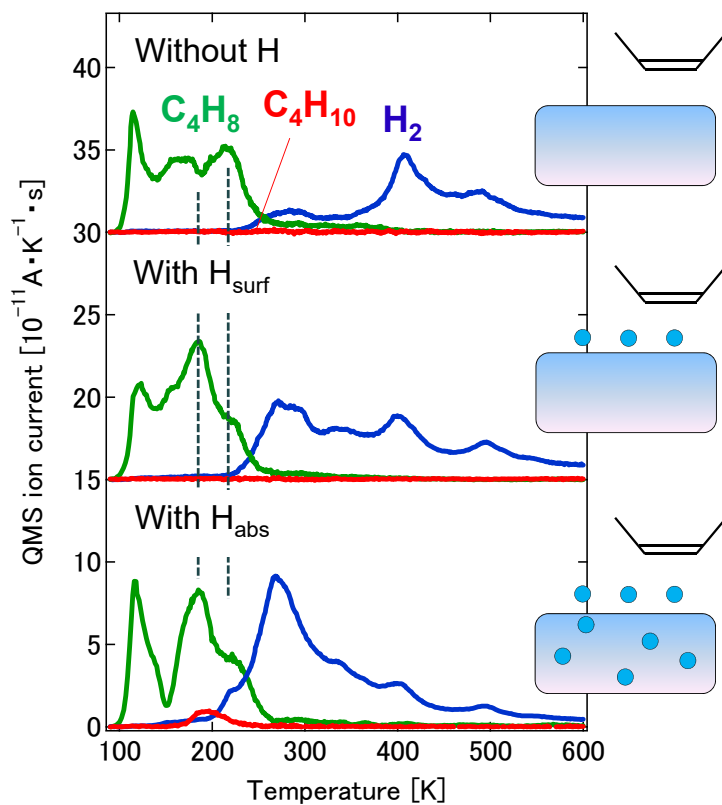
π bond



di- σ



Reaction TPD: $C_4H_8 + H$ on Pd(110)



H_2 desorption:

C_4H_8 decomposition

LT desorption of C_4H_8

modification of C_4H_8

C_4H_{10} production

Summary

Use of Quantum beams (e^- , Ion, Neutron, Muon)
for materials analysis

Elastic & Inelastic scattering by materials
information on atomic & electronic structures
→ origin of material functions

Application: analysis of H

Electronic effect: metal and semiconductor

Proton motion: Energetics, dynamics, kinetics

Quantum effects and chemical reaction